

16.360 Lecture 1

Units and dimensions

- Six fundamental International System of Units

Dimensions	Unit	symbol
Length	meter	m
Mass	kilogram	kg
Time	second	s
Electric Current	Ampere	A
Temperature	Kelvin	K
Amount of substance	mole	mol

- any other dimension can be derived from the fundamental dimensions, e.g.:

$$F = ma = m \frac{dv}{dt} = \text{kgm} / \text{s}^2$$

$$E = \frac{F}{q} = \frac{F}{\int Idt} = \text{kgm} / \text{As}^3$$

16.360 Lecture 1

Electromagnetic spectrum

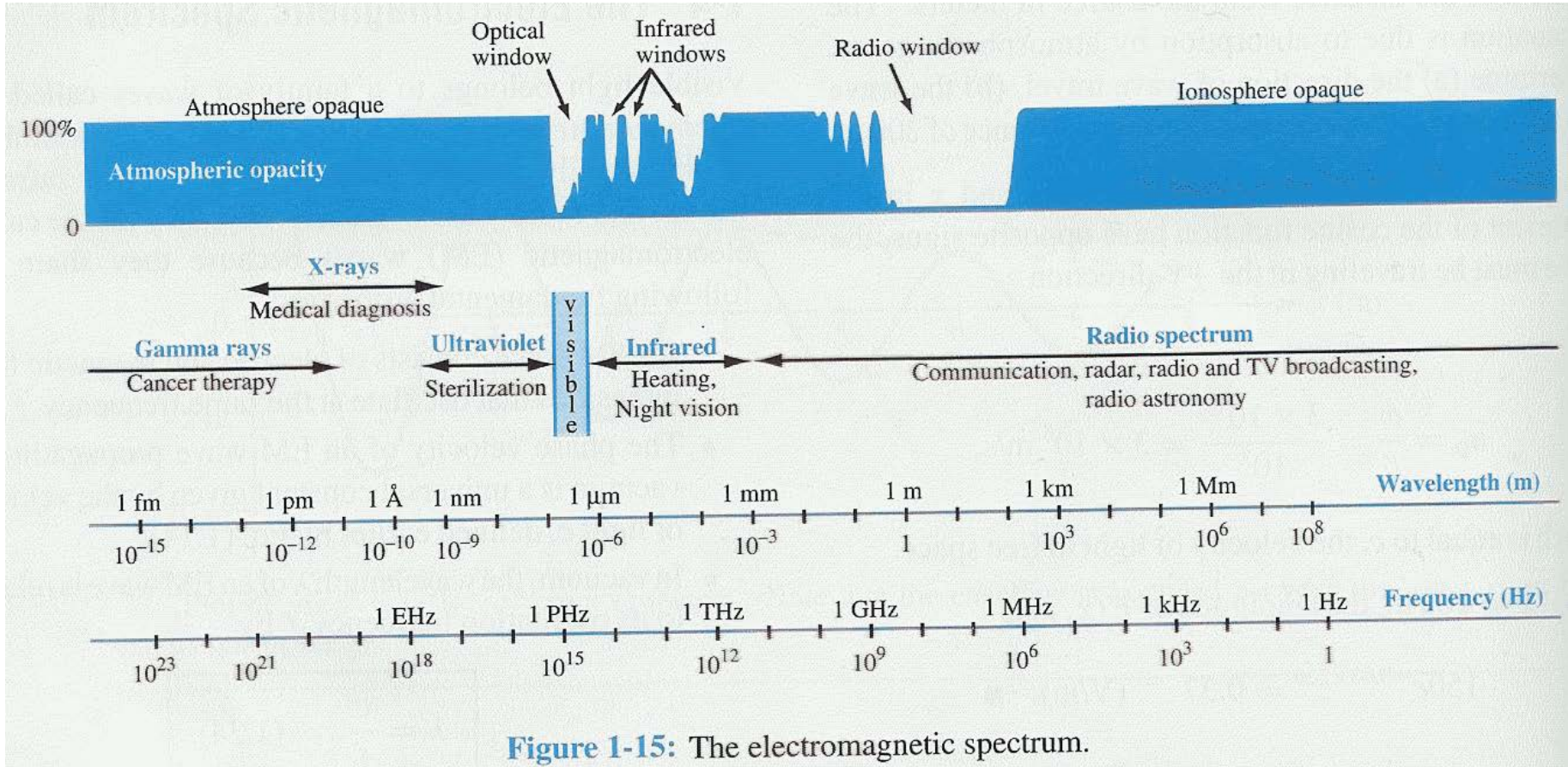
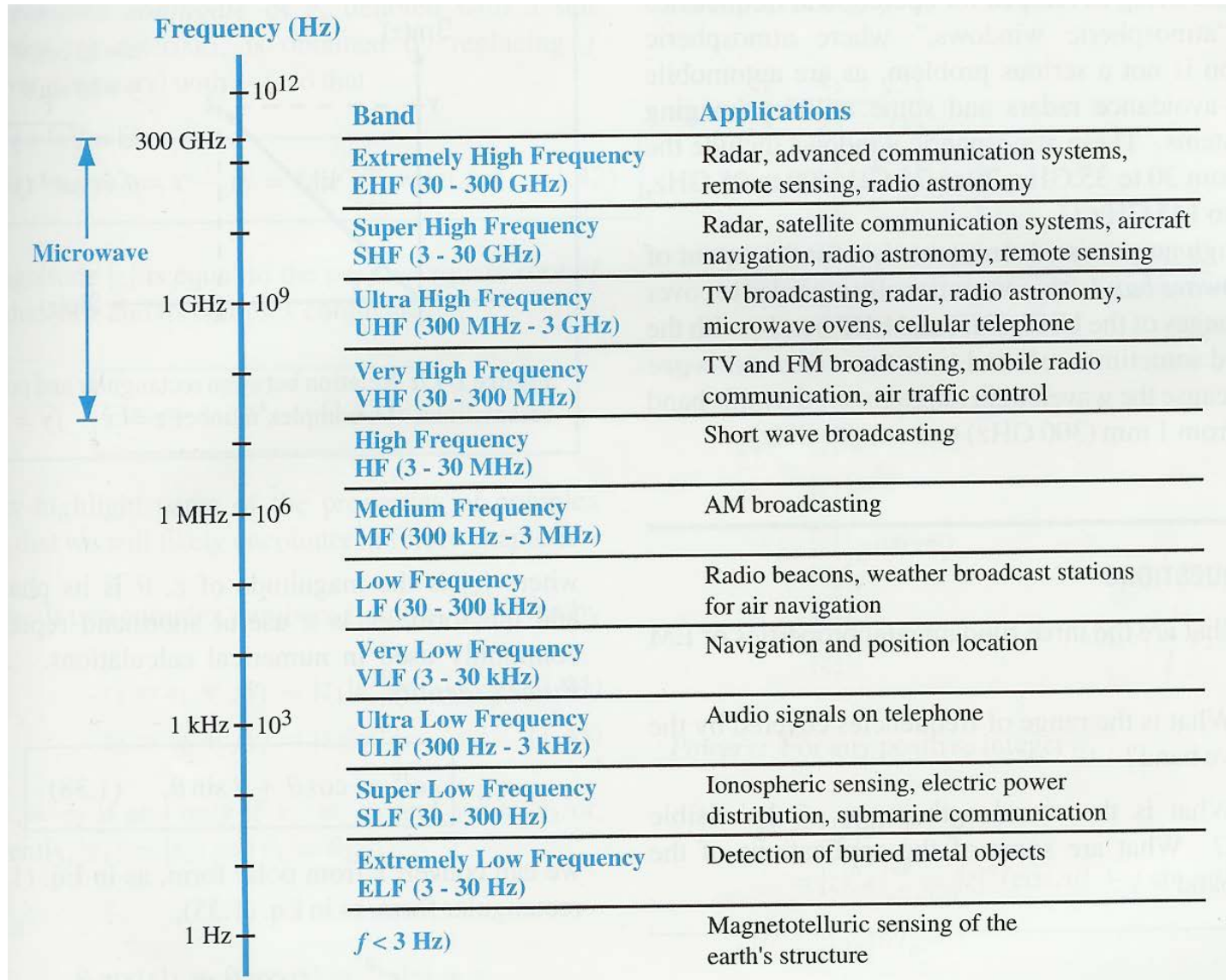


Figure 1-15: The electromagnetic spectrum.

16.360 Lecture 1

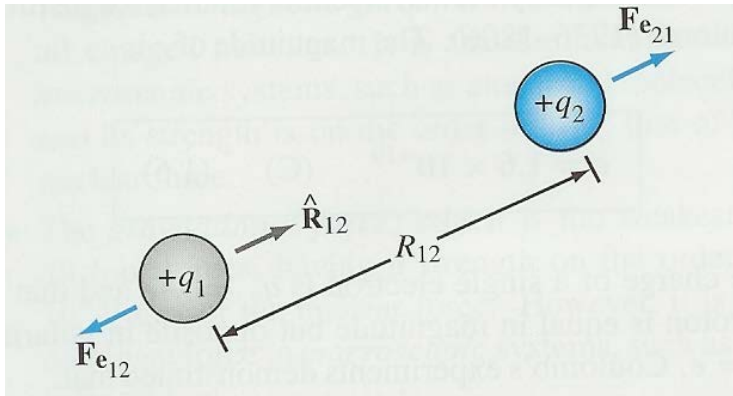
Electromagnetic bands and applications



16.360 Lecture 2

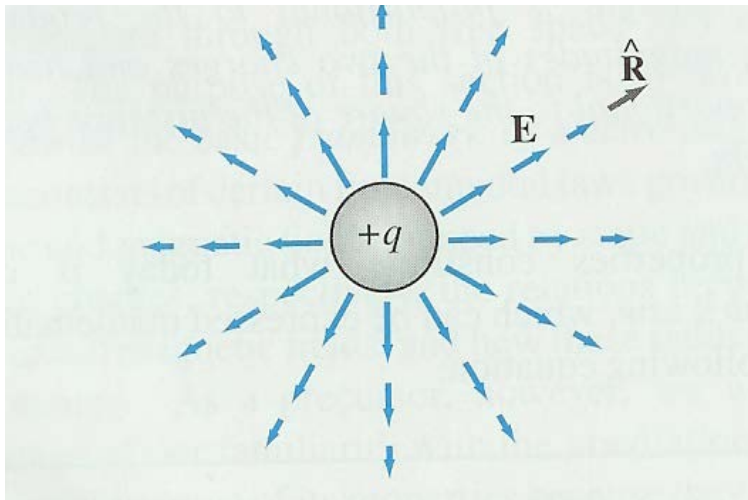
Electric field

- Electric forces on point charges, Columb's law



$$\vec{F}_{12} = \hat{R}_{12} \frac{q_1 q_2}{4\pi\epsilon_0 R_{12}^2},$$

$$\vec{F}_{21} = \hat{R}_{21} \frac{q_1 q_2}{4\pi\epsilon_0 R_{21}^2},$$

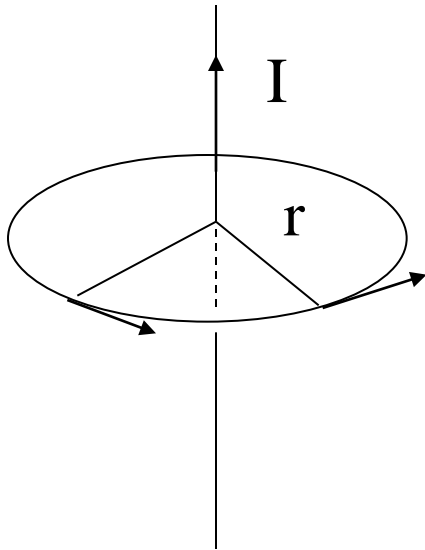


$$\vec{F}_{21} = q_1 \frac{q}{4\pi\epsilon_0 R^2} = q_1 \vec{E},$$

$$\vec{E} = \hat{R} \frac{q}{4\pi\epsilon_0 R^2},$$

16.360 Lecture 2

Magnetic field by constant current



$$\mathbf{B} = \frac{\mu I}{2\pi r} \hat{\phi}$$

$$\mu = \mu_r \mu_0,$$

μ_r : relative magnetic permeability

$\mu_r = 1$ for most materials

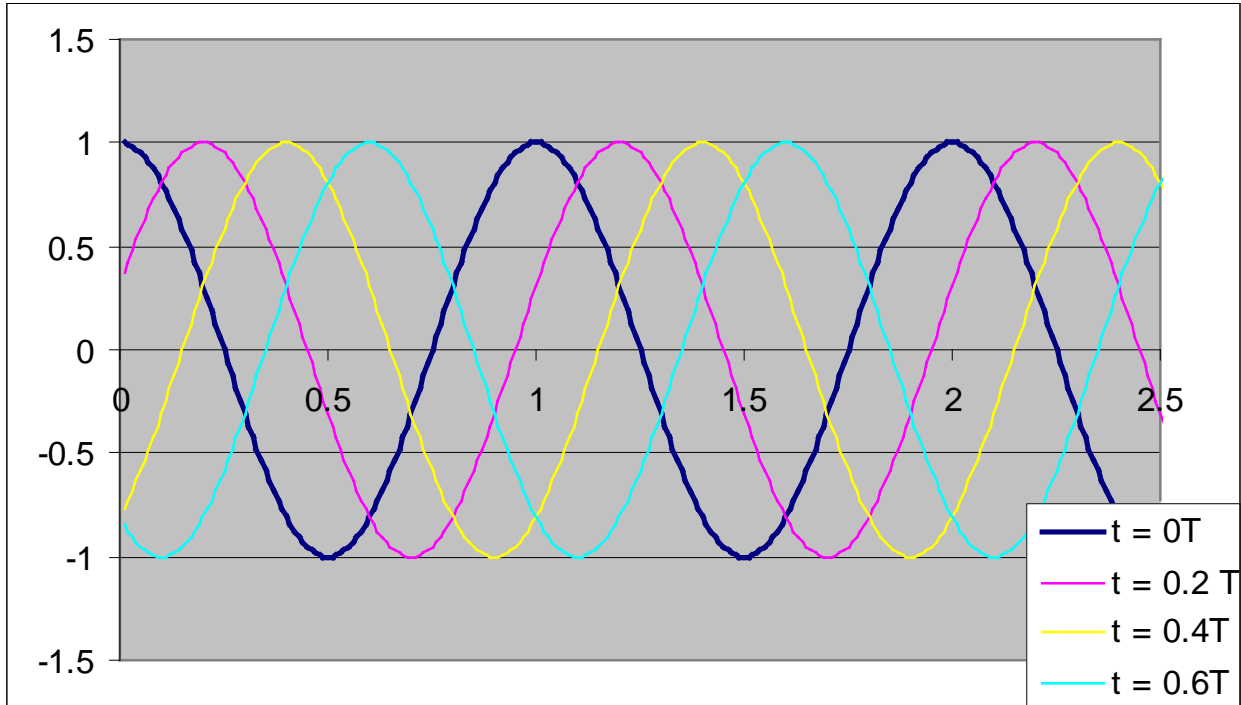
$$\begin{aligned} \mathbf{H} &= \frac{\mathbf{B}}{\mu} \\ &= \frac{I}{2\pi r} \hat{\phi} \end{aligned}$$

16.360 Lecture 3

Traveling wave

$$y(x,t) = A \cos(2\pi t/T - 2\pi x/\lambda),$$

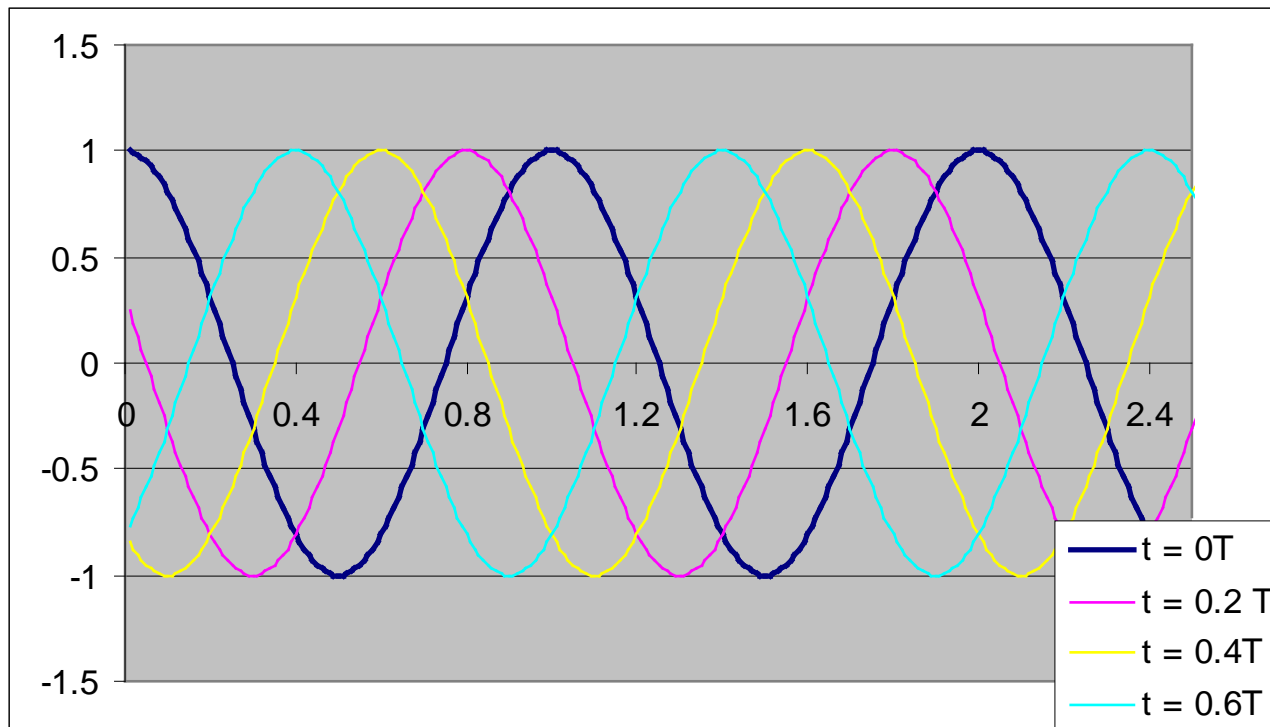
$$\phi(x,t) = 2\pi t/T - 2\pi x/\lambda, \quad y(x,t) = A \cos\phi(x,t),$$



16.360 Lecture 3

Traveling wave

$$y(x,t) = A \cos(2\pi t/T + 2\pi x/\lambda),$$

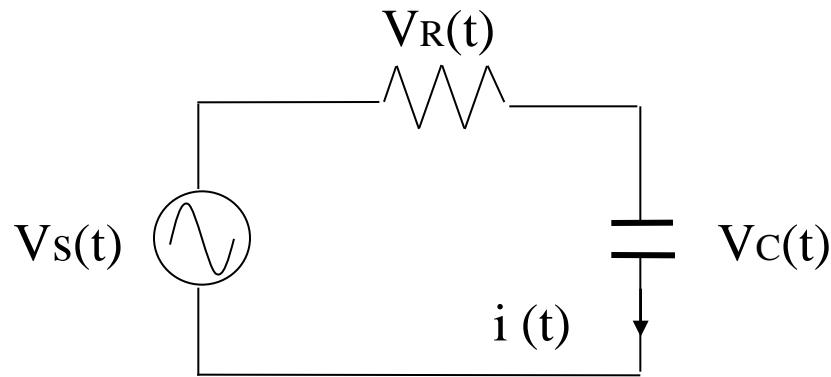


$$\text{Velocity} = \frac{0.6\lambda}{0.6T} \\ = \lambda/T$$

Phase velocity:

$$V_p = \frac{dx}{dt} \\ = -\lambda/T$$

- Phasor



$$V_S(t) = V_0 \sin(\omega t + \phi_0),$$

$$V_R(t) = i(t)R,$$

$$V_C(t) = \int i(t) dt / C,$$

$$V_S(t) = V_R(t) + V_C(t),$$

$$\longrightarrow V_0 \sin(\omega t + \phi_0) = \int i(t) dt / C + i(t)R, \quad \text{Integral equation,}$$

Using phasor to solve integral and differential equations

- Phasor

$$Z(t) = \text{Re}(\tilde{Z} e^{j\omega t})$$

\tilde{Z} is time independent function of $Z(t)$, i.e. phasor

$$\begin{aligned} V_s(t) &= V_0 \sin(\omega t + \phi_0) \\ &= \text{Re}(V_0 e^{j(\phi_0 - \pi/2)} e^{j\omega t}) \\ &= \text{Re}(\tilde{V} e^{j\omega t}), \end{aligned}$$

$$\tilde{V} = V_0 e^{j(\phi_0 - \pi/2)},$$

- Phasor

$$i(t) = \text{Re}(\tilde{I} e^{j\omega t})$$

$$\begin{aligned} \int i(t) dt &= \int \text{Re}(\tilde{I} e^{j\omega t}) dt \\ &= \text{Re}(\tilde{I} \frac{1}{j\omega} e^{j\omega t}), \end{aligned}$$

time domain equation:

$$V_0 \sin(\omega t + \phi_0) = \int i(t) dt / C + i(t)R, \quad \longrightarrow$$

phasor domain equation:

$$\text{Re}(\tilde{V} e^{j\omega t}) = \text{Re}(\tilde{I} \frac{1}{j\omega} e^{j\omega t}) / C + \text{Re}(\tilde{I} e^{j\omega t}),$$

- Phasor domain

$$\tilde{V} = \tilde{I} \frac{1}{j\omega C} + \tilde{I} R$$

$$\tilde{I} = \frac{\tilde{V}}{R + 1/(j\omega C)}$$

$$= \frac{V_0 e^{j(\phi_0 - \pi/2)}}{R + 1/(j\omega C)},$$

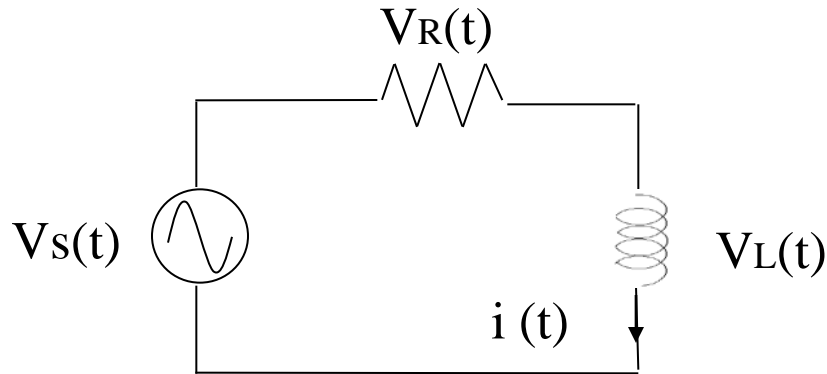


Back to time domain:

$$i(t) = \text{Re}(\tilde{I} e^{j\omega t}) = \text{Re}\left(\frac{V_0 e^{j(\phi_0 - \pi/2)}}{R + 1/(j\omega C)} e^{j\omega t}\right)$$

16.360 Lecture 3

- An Example :



$$V_S(t) = V_0 \sin(\omega t + \phi_0),$$

$$V_R(t) = i(t)R,$$

$$V_L(t) = L \frac{di(t)}{dt},$$

$$V_S(t) = V_R(t) + V_L(t),$$

→ $V_0 \sin(\omega t + \phi_0) = L \frac{di(t)}{dt} + i(t)R,$ differential equation,

Using phasor to solve the differential equation.

- Phasor

$$i(t) = \text{Re}(\tilde{I} e^{j\omega t})$$

$$di(t)/dt = \text{Re}(d\tilde{I} e^{j\omega t})/dt$$

$$= \text{Re}(\tilde{I} j\omega e^{j\omega t}),$$

time domain equation:

$$V_0 \sin(\omega t + \phi_0) = L di(t)/dt + i(t)R, \quad \longrightarrow$$

phasor domain equation:

$$\text{Re}(\tilde{V} e^{j\omega t}) = \text{Re}(\tilde{I} j\omega e^{j\omega t})L + \text{Re}(\tilde{I} e^{j\omega t}),$$

16.360 Lecture 3

- Phasor domain

$$\tilde{V} = \tilde{I} j\omega L + \tilde{I} R,$$

$$\begin{aligned}\tilde{I} &= \frac{\tilde{V}}{R + (j\omega L)} \\ &= \frac{V_0 e^{j(\phi_0 - \pi/2)}}{R + j\omega L},\end{aligned}$$



Back to time domain:

$$i(t) = \text{Re}(\tilde{I} e^{j\omega t}) = \text{Re}\left(\frac{V_0 e^{j(\phi_0 - \pi/2)}}{R + j\omega L} e^{j\omega t}\right)$$

16.360 Lecture 3

- Steps of transferring integral or differential equations to linear equations using phasor.
 1. Express time-dependent variables as phasor.
 2. Rewrite integral or differential equations in phasor domain.
 3. Solve phasor domain equations
 4. Change phasors variable to their time domain value

16.360 Lecture 3

- Electromagnetic spectrum.

Recall relation: $\lambda f = v$.

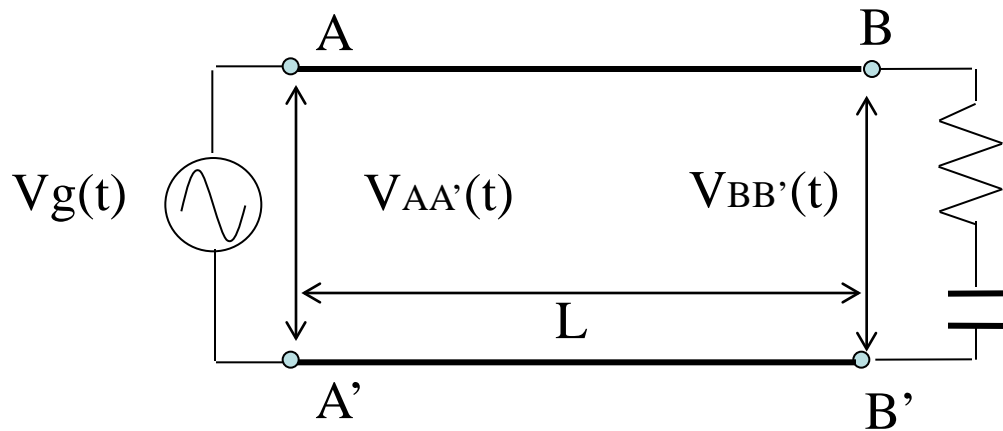
- Some important wavelength ranges:
 1. Fiber optical communication: $\lambda = 1.3 - 1.5\mu\text{m}$.
 2. Free space communication: $\sim 700\text{nm} - 980\text{nm}$.
 3. TV broadcasting and cellular phone: $300\text{MHz} - 3\text{GHz}$.
 4. Radar and remote sensing: $30\text{GHz} - 300\text{GHz}$

16.360 Lecture 3

- Transmission lines
 1. Transmission line parameters, equations
 2. Wave propagations
 3. Lossless line, standing wave and reflection coefficient
 4. Input impedance
 5. Special cases of lossless line
 6. Power flow
 7. Smith chart
 8. Impedance matching
 9. Transients on transmission lines

16.360 Lecture 3

1. Transmission line parameters, equations



$$V_{AA'}(t) = V_g(t) = V_0 \cos(\omega t),$$

Low frequency circuits:

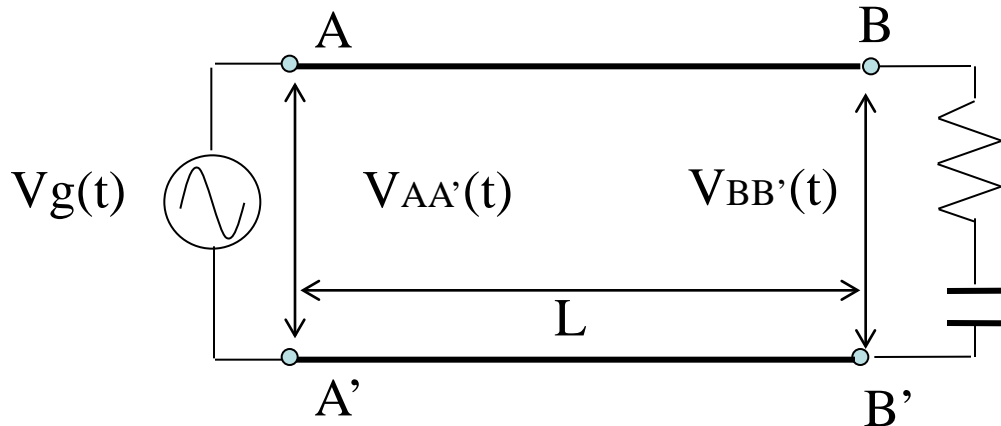
$$V_{BB'}(t) = V_{AA'}(t)$$

← Approximate result

$$\begin{aligned} V_{BB'}(t) &= V_{AA'}(t-t_d) = V_{AA'}(t-L/c) \\ &= V_0 \cos(\omega(t-L/c)), \end{aligned}$$

16.360 Lecture 3

1. Transmission line parameters, equations



Recall: $f\lambda=c$, and $\omega = 2\pi f$

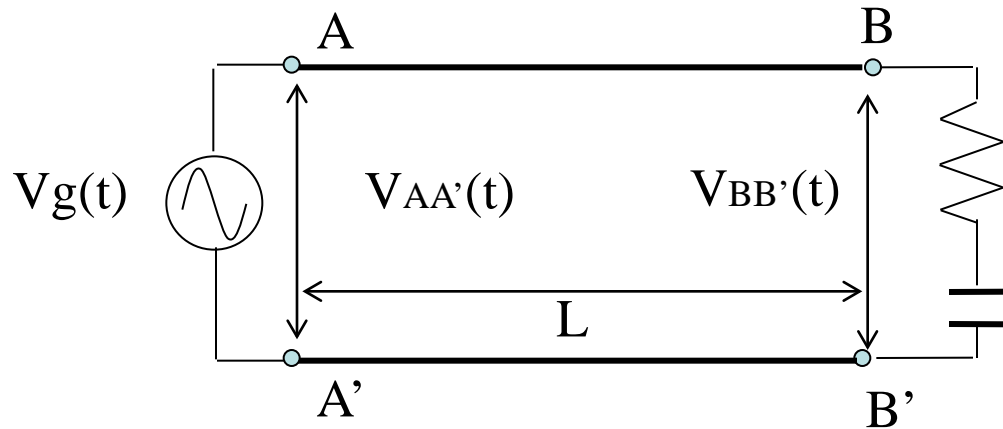
$$\begin{aligned} V_{BB'}(t) &= V_{AA'}(t-t_d) = V_{AA'}(t-L/c) \\ &= V_0 \cos(\omega(t-L/c)) \\ &= V_0 \cos(\omega t - 2\pi L/\lambda), \end{aligned}$$

If $\lambda \gg L$, $V_{BB'}(t) \approx V_0 \cos(\omega t) = V_{AA'}(t)$,

If $\lambda \leq L$, $V_{BB'}(t) \neq V_{AA'}(t)$, the circuit theory has to be replaced.

16.360 Lecture 4

- Transmission line parameters



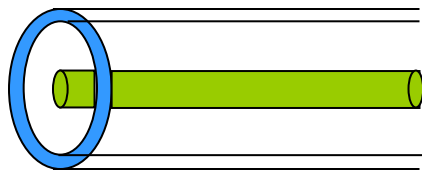
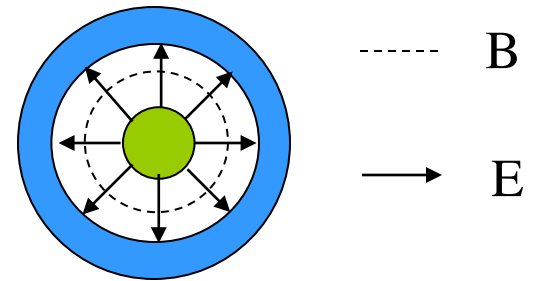
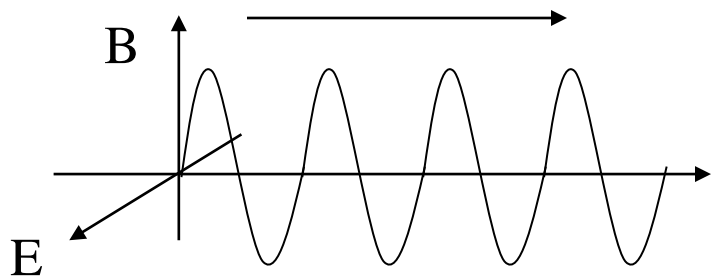
- time delay

$$V_{BB'}(t) = V_{AA'}(t-t_d) = V_{AA'}(t-L/v_p),$$

- Reflection: the voltage has to be treated as a wave, some bounce back
- power loss: due to reflection and some other loss mechanism,
- Dispersion: in material, V_p could be different for different wavelengths

16.360 Lecture 4

- Types of transmission lines
 - Transverse electromagnetic (TEM) transmission lines



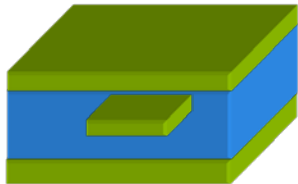
a) Coaxial line



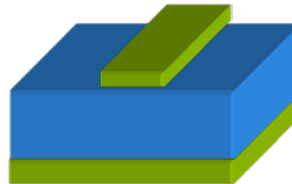
b) Two-wire line



c) Parallel-plate line



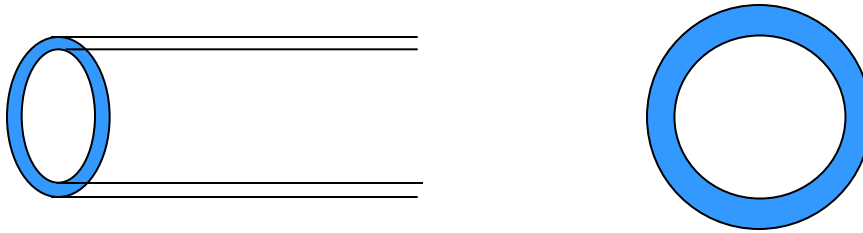
d) Strip line



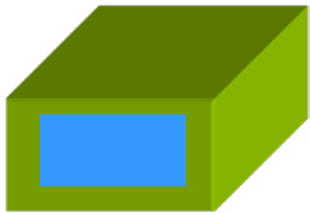
e) Microstrip line

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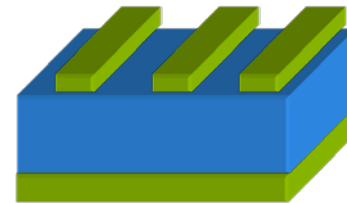
- Types of transmission lines
 - Higher-order transmission lines



a) Optical fiber



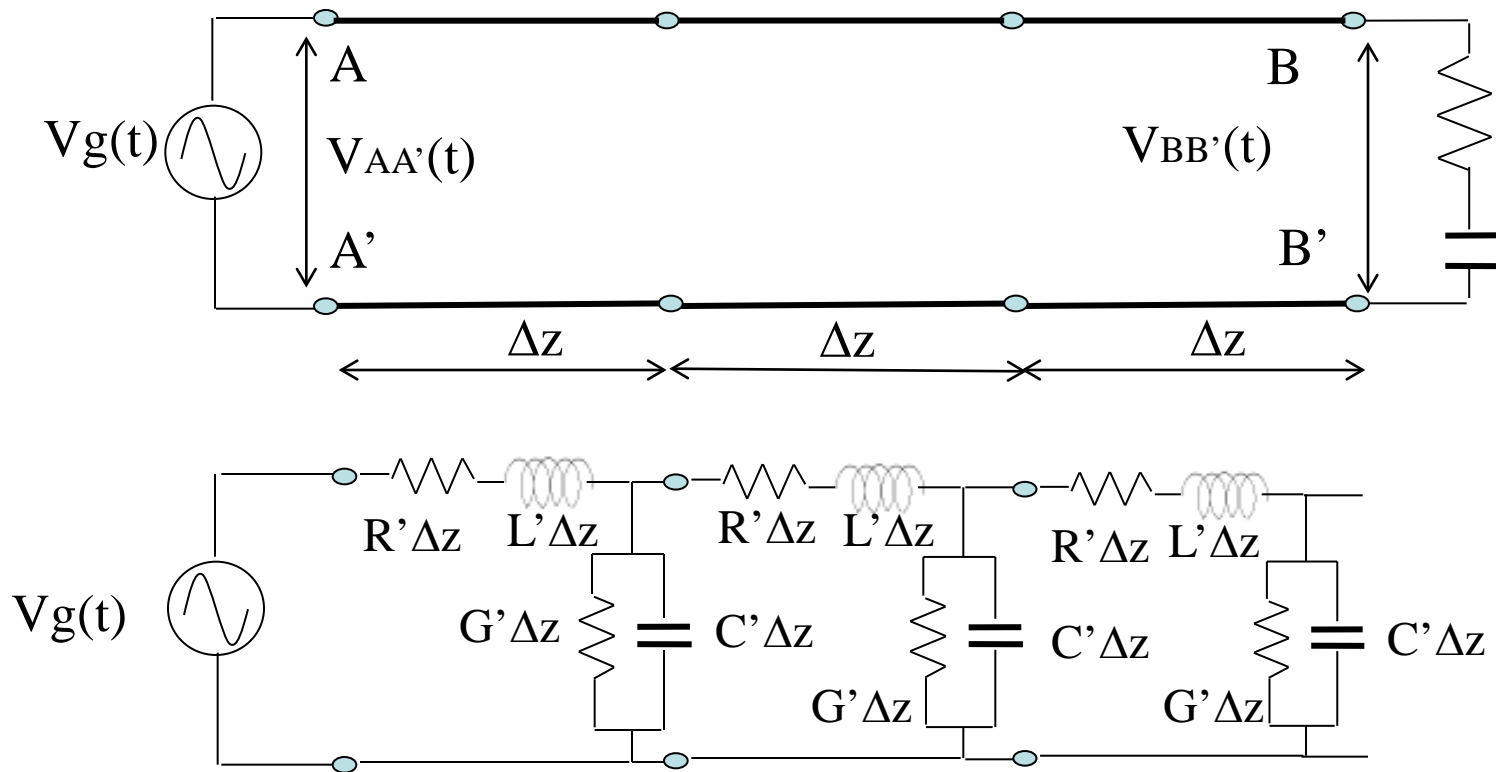
b) Rectangular waveguide



c) Coplanar waveguide

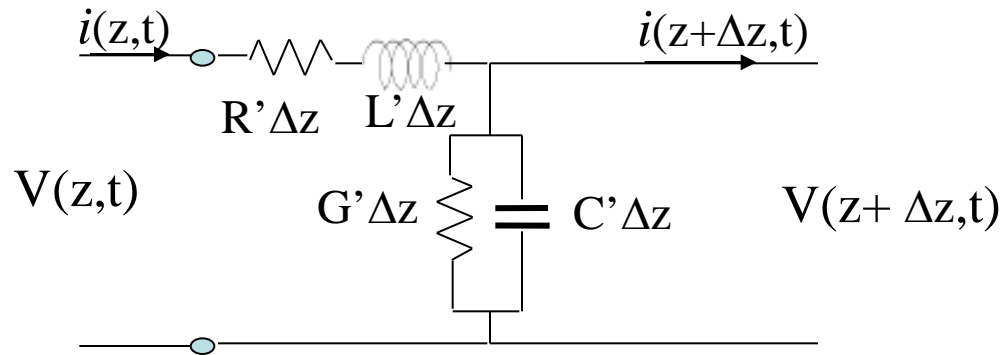
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- Lumped-element Model
 - Represent transmission lines as parallel-wire configuration



16.360 Lecture 4

- Transmission line equations
- Represent transmission lines as parallel-wire configuration

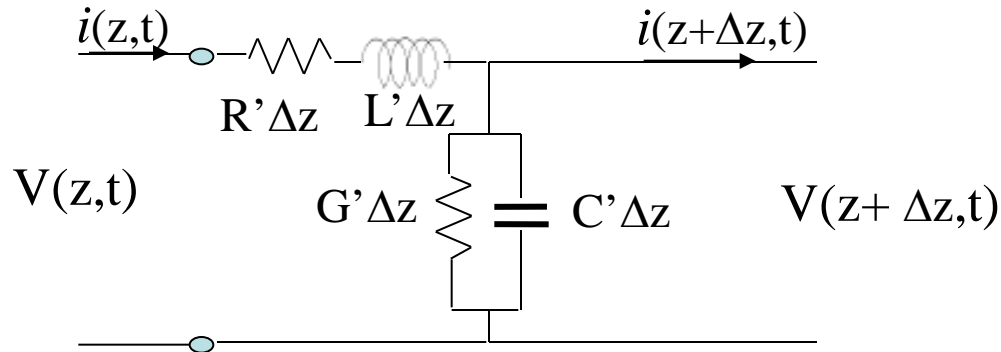


$$V(z,t) = R'\Delta z i(z,t) + L'\Delta z \frac{\partial i(z,t)}{\partial t} + V(z+\Delta z,t), \quad (1)$$

$$i(z,t) = G'\Delta z V(z+\Delta z,t) + C'\Delta z \frac{\partial V(z+\Delta z,t)}{\partial t} + i(z+\Delta z,t), \quad (2)$$

16.360 Lecture 4

- Transmission line equations



$$V(z,t) = R' \Delta z i(z,t) + L' \Delta z \partial i(z,t) / \partial t + V(z + \Delta z, t), \quad (1)$$

→ $-V(z + \Delta z, t) + V(z, t) = R' \Delta z i(z, t) + L' \Delta z \partial i(z, t) / \partial t$

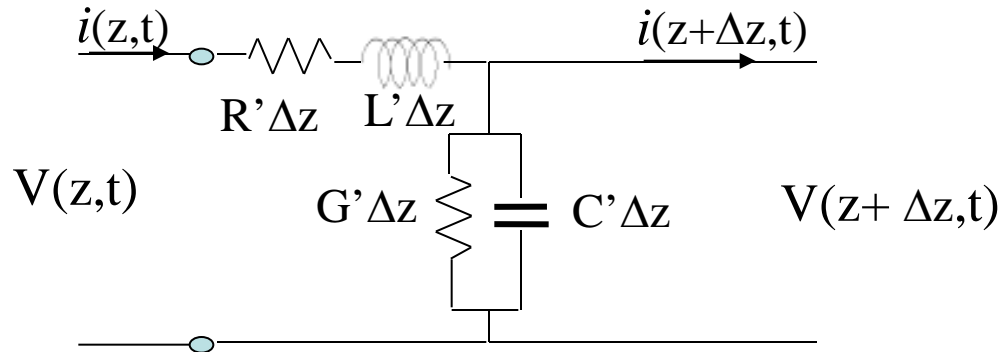
→ $-\partial V(z, t) / \partial z = R' i(z, t) + L' \partial i(z, t) / \partial t, \quad (3)$

Rewrite $V(z, t)$ and $i(z, t)$ as phasors, for sinusoidal $V(z, t)$ and $i(z, t)$:

$$V(z, t) = \text{Re}(\tilde{V}(z) e^{j\omega t}), \quad i(z, t) = \text{Re}(\tilde{i}(z) e^{j\omega t}),$$

16.360 Lecture 4

- Transmission line equations



Recall:

$$di(t)/dt = \text{Re}(d\tilde{i} e^{j\omega t})/dt = \text{Re}(\tilde{i} j\omega e^{j\omega t}),$$

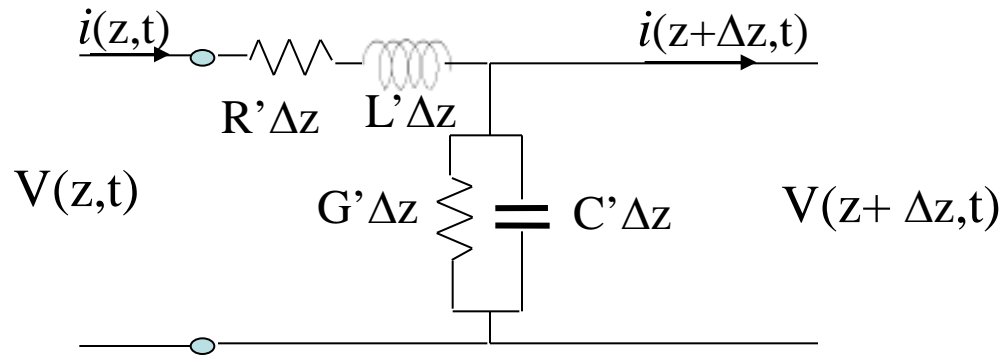
$$-\partial V(z,t)/\partial z = R' i(z,t) + L' \partial i(z,t)/\partial t, \quad (3)$$

→

$$-d\tilde{V}(z)/dz = R' \tilde{i}(z) + j\omega L' \tilde{i}(z), \quad (4)$$

16.360 Lecture 4

- Transmission line equations
- Represent transmission lines as parallel-wire configuration

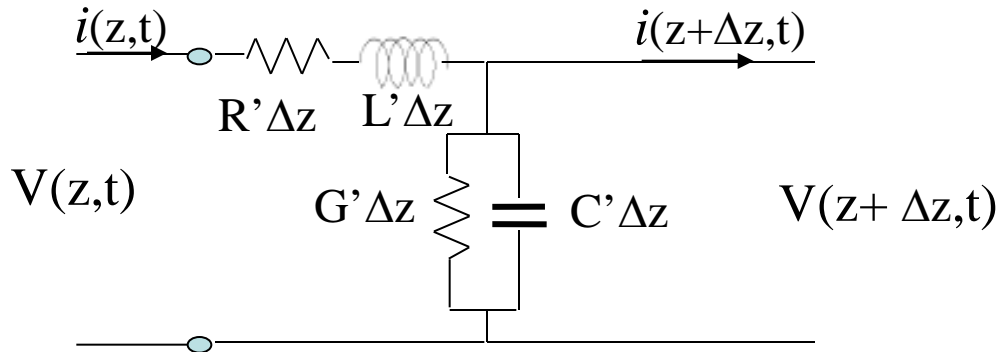


$$V(z,t) = R'\Delta z i(z,t) + L'\Delta z \frac{\partial i(z,t)}{\partial t} + V(z+\Delta z,t), \quad (1)$$

$$i(z,t) = G'\Delta z V(z+\Delta z,t) + C'\Delta z \frac{\partial V(z+\Delta z,t)}{\partial t} + i(z+\Delta z,t), \quad (2)$$

16.360 Lecture 4

- Transmission line equations



$$i(z,t) = G'\Delta z V(z+\Delta z,t) + C'\Delta z \frac{\partial V(z+\Delta z,t)}{\partial t} + i(z+\Delta z,t), \quad (2)$$

$$\Rightarrow -i(z+\Delta z,t) + i(z,t) = G'\Delta z V(z+\Delta z,t) + C'\Delta z \frac{\partial V(z+\Delta z,t)}{\partial t}$$

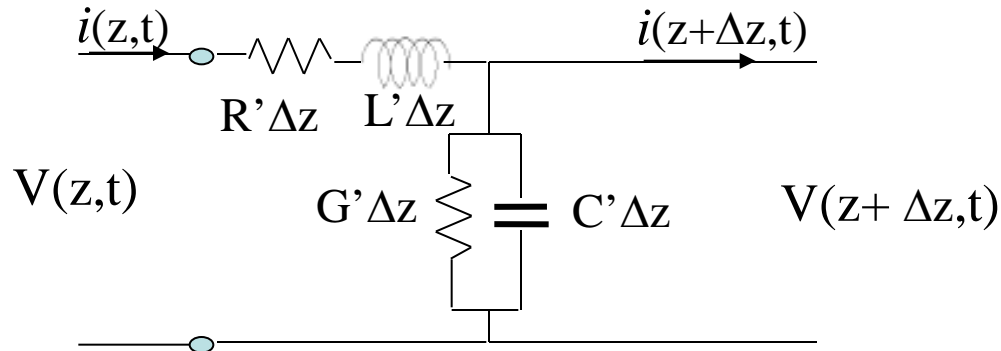
$$\Rightarrow -\frac{\partial i(z,t)}{\partial z} = G' V(z,t) + C' \frac{\partial V(z,t)}{\partial t}, \quad (5)$$

Rewrite $V(z,t)$ and $i(z,t)$ as phasors, for sinusoidal $V(z,t)$ and $i(z,t)$:

$$V(z,t) = \text{Re}(\tilde{V}(z) e^{j\omega t}), \quad i(z,t) = \text{Re}(\tilde{i}(z) e^{j\omega t}),$$

16.360 Lecture 4

- Transmission line equations



Recall:

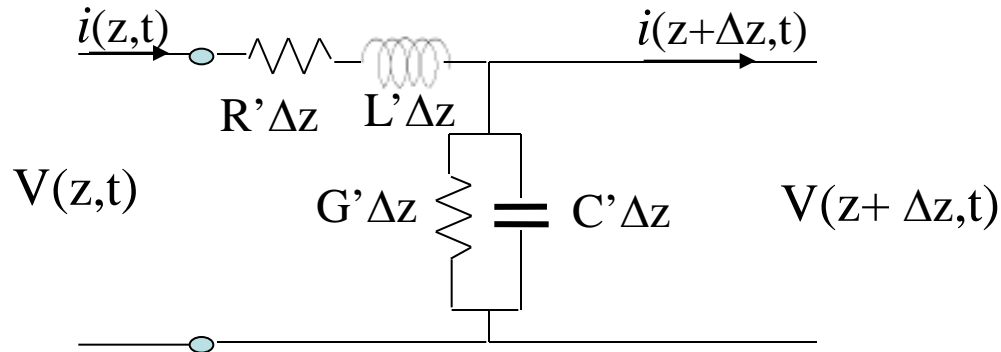
$$dV(t)/dt = \text{Re}(d \tilde{V} e^{j\omega t})/dt = \text{Re}(\tilde{V} j\omega e^{j\omega t}),$$

$$-\partial i(z,t)/\partial z = G' V(z,t) + C' \partial V(z,t)/\partial t, \quad (6)$$

→
$$-d \tilde{i}(z)/dz = G' \tilde{V}(z) + j\omega C' \tilde{V}(z), \quad (7)$$

16.360 Lecture 4

- Telegrapher's equation in phasor domain



$$\left\{ \begin{array}{l} -d\tilde{V}(z)/dz = R' \tilde{i}(z) + j\omega L' \tilde{i}(z), \\ -d\tilde{i}(z)/dz = G' \tilde{V}(z) + j\omega C' \tilde{V}(z), \end{array} \right. \quad (4)$$

Take d/dz on both sides of eq. (4)

$$\longrightarrow -d^2\tilde{V}(z)/dz^2 = R' d\tilde{i}(z)/dz + j\omega L' d\tilde{i}(z)/dz, \quad (8)$$


16.360 Lecture 4

- Telegrapher's equation in phasor domain

$$\left\{ \begin{array}{l} -d\tilde{V}(z)/dz = R' \tilde{i}(z) + j\omega L' \tilde{i}(z), \\ -d\tilde{i}(z)/dz = G' \tilde{V}(z) + j\omega C' \tilde{V}(z), \end{array} \right. \quad (4)$$

$$-d^2\tilde{V}(z)/dz^2 = R' d\tilde{i}(z)/dz + j\omega L' d\tilde{i}(z)/dz, \quad (8)$$

substitute (7) to (8)

 $d^2\tilde{V}(z)/dz^2 = (R' + j\omega L') (G' + j\omega C') \tilde{V}(z),$ or

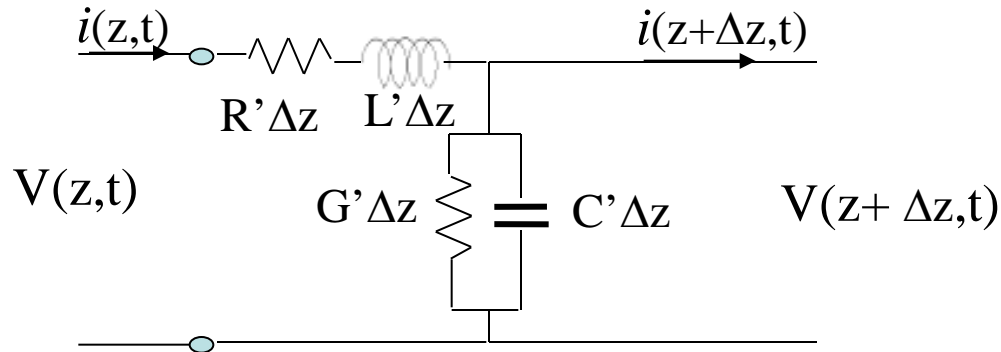
$$\boxed{d^2\tilde{V}(z)/dz^2 - (R' + j\omega L') (G' + j\omega C') \tilde{V}(z) = 0,} \quad (9)$$

$$d^2\tilde{V}(z)/dz^2 - \gamma^2 \tilde{V}(z) = 0, \quad (10)$$

$$\boxed{\gamma^2 = (R' + j\omega L') (G' + j\omega C'),} \quad (11)$$

16.360 Lecture 4

- Telegrapher's equation in phasor domain



$$\left\{ \begin{array}{l} -d\tilde{V}(z)/dz = R' \tilde{i}(z) + j\omega L' \tilde{i}(z), \\ -d\tilde{i}(z)/dz = G' \tilde{V}(z) + j\omega C' \tilde{V}(z), \end{array} \right. \quad (4)$$

$$\left\{ \begin{array}{l} -d\tilde{V}(z)/dz = R' \tilde{i}(z) + j\omega L' \tilde{i}(z), \\ -d\tilde{i}(z)/dz = G' \tilde{V}(z) + j\omega C' \tilde{V}(z), \end{array} \right. \quad (7)$$

Take d/dz on both sides of eq. (7)

$$\longrightarrow -d^2 \tilde{i}(z)/dz^2 = G' d\tilde{V}(z)/dz + j\omega C' d\tilde{V}(z)/dz, \quad (12)$$


16.360 Lecture 4

- Telegrapher's equation in phasor domain

$$\left\{ \begin{array}{l} -d\tilde{V}(z)/dz = R' \tilde{i}(z) + j\omega L' \tilde{i}(z), \\ -d\tilde{i}(z)/dz = G' \tilde{V}(z) + j\omega C' \tilde{V}(z), \end{array} \right. \quad (4)$$

$$-d^2 \tilde{i}(z)/dz^2 = G' d\tilde{V}(z)/dz + j\omega C' d\tilde{V}(z)/dz, \quad (12)$$

substitute (4) to (12)

 $d^2 \tilde{i}(z)/dz^2 = (R' + j\omega L') (G' + j\omega C') \tilde{i}(z),$ or

$$d^2 \tilde{i}(z)/dz^2 - (R' + j\omega L') (G' + j\omega C') \tilde{i}(z) = 0, \quad (9)$$

$$d^2 \tilde{i}(z)/dz^2 - \gamma^2 \tilde{i}(z) = 0, \quad (13)$$

$$\gamma^2 = (R' + j\omega L') (G' + j\omega C'), \quad (11)$$

16.360 Lecture 4

- Wave equations

$$\left\{ \begin{array}{l} d^2 \tilde{V}(z)/dz^2 - \gamma^2 \tilde{V}(z) = 0, \\ d^2 \tilde{i}(z)/dz^2 - \gamma^2 \tilde{i}(z) = 0, \end{array} \right. \quad \begin{array}{l} (10) \\ (13) \end{array}$$

$$\gamma = \alpha + j\beta,$$

$$\alpha = \operatorname{Re} \sqrt{(R' + j\omega L')(G' + j\omega C')},$$

$$\beta = \operatorname{Im} \sqrt{(R' + j\omega L')(G' + j\omega C')},$$

16.360 Lecture 4

- Summary
 - Transmission line parameters
 - Types of transmission lines
 - Lumped-element model
 - Transmission line equations
 - Telegrapher's equations
 - Wave equations

16.360 Lecture 4

- Next lecture
 - Wave propagation on a Transmission line
 - Characteristic impedance
 - Standing wave and traveling wave
 - Lossless transmission line
 - Reflection coefficient

16.360 Lecture 5

- Today:
 - Wave propagation on a Transmission line
 - Characteristic impedance
 - Standing wave and traveling wave
 - Lossless transmission line
 - Reflection coefficient

16.360 Lecture 5

- Wave equations

$$\left\{ \begin{array}{l} d^2 \tilde{V}(z)/dz^2 - \gamma^2 \tilde{V}(z) = 0, \\ d^2 \tilde{i}(z)/dz^2 - \gamma^2 \tilde{i}(z) = 0, \end{array} \right. \quad \begin{array}{l} (10) \\ (13) \end{array}$$

$$\gamma = \alpha + j\beta,$$

$$\alpha = \operatorname{Re} \sqrt{(R' + j\omega L')(G' + j\omega C')},$$

$$\beta = \operatorname{Im} \sqrt{(R' + j\omega L')(G' + j\omega C')},$$

Solving the second order differential equation

$$\left\{ \begin{array}{l} \tilde{V}(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z} \end{array} \right. \quad (14)$$

$$\left\{ \begin{array}{l} \tilde{i}(z) = I_0^+ e^{-\gamma z} + I_0^- e^{\gamma z} \end{array} \right. \quad (15)$$

16.360 Lecture 5

- Wave equations

$$\left\{ \begin{array}{l} \tilde{V}(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z} \\ \tilde{i}(z) = I_0^+ e^{-\gamma z} + I_0^- e^{\gamma z} \end{array} \right. \quad (14)$$

$$\left\{ \begin{array}{l} \tilde{V}(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z} \\ \tilde{i}(z) = I_0^+ e^{-\gamma z} + I_0^- e^{\gamma z} \end{array} \right. \quad (15)$$

where:

V_0^+ and V_0^- are determined by boundary conditions.

I_0^+ and I_0^- are related to V_0^+ and V_0^- by characteristic impedance Z_0 .

16.360 Lecture 5

- Characteristic impedance Z_0

$$\begin{cases} \tilde{V}(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z} & (14) \\ \tilde{i}(z) = I_0^+ e^{-\gamma z} + I_0^- e^{\gamma z} & (15) \end{cases}$$

recall:

$$-d\tilde{V}(z)/dz = R' \tilde{i}(z) + j\omega L' \tilde{i}(z), \quad (4)$$

$$\Rightarrow \gamma V_0^+ e^{-\gamma z} - \gamma V_0^- e^{\gamma z} = (R' + j\omega L') \tilde{i}(z), \quad (16)$$

$$\Rightarrow \tilde{i}(z) = \frac{\gamma}{(R' + j\omega L')} (V_0^+ e^{-\gamma z} - V_0^- e^{\gamma z})$$

$$\Rightarrow I_0^+ = \frac{\gamma}{(R' + j\omega L')} V_0^+ \quad (17) \quad I_0^- = \frac{-\gamma}{(R' + j\omega L')} V_0^- \quad (18)$$

16.360 Lecture 5

- Characteristic impedance Z_0

$$I_0^+ = \frac{\gamma}{(R' + j\omega L')} V_0^+ \quad (17)$$

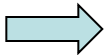
$$I_0^- = \frac{-\gamma}{(R' + j\omega L')} V_0^- \quad (18)$$

Define characteristic impedance Z_0

$$\begin{aligned} Z_0 &\equiv \frac{V_0^+}{I_0^+} = \frac{(R' + j\omega L')}{\gamma} \\ &= \sqrt{\frac{(R' + j\omega L')}{(G' + j\omega C')}} \end{aligned}$$

recall:

$$\gamma = \sqrt{(R' + j\omega L')(G' + j\omega C')}$$



16.360 Lecture 5

- Summary:

$$\left\{ \begin{array}{l} \tilde{V}(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z} \\ \tilde{i}(z) = I_0^+ e^{-\gamma z} + I_0^- e^{\gamma z} \end{array} \right. \quad (14)$$

$$\left\{ \begin{array}{l} \tilde{V}(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z} \\ \tilde{i}(z) = I_0^+ e^{-\gamma z} + I_0^- e^{\gamma z} \end{array} \right. \quad (15)$$

$$\gamma = \sqrt{(R' + j\omega L')(G' + j\omega C')} \quad (19)$$

$$Z_0 \equiv \frac{V_0^+}{I_0^+} = \sqrt{\frac{(R' + j\omega L')}{(G' + j\omega C')}} \quad (20)$$

16.360 Lecture 5

- Example, an air line :

$$R' = 0 \Omega, G' = 0 /\Omega, Z_0 = 50\Omega, \beta = 20 \text{ rad/m}, f = 700 \text{ MHz}$$


$$L' = ? \text{ and } C' = ?$$

solution:

$$Z_0 = \sqrt{\frac{(R' + j\omega L')}{(G' + j\omega C')}} = \sqrt{\frac{j\omega L'}{j\omega C'}} = 50\Omega$$

$$\gamma = \sqrt{(R' + j\omega L')(G' + j\omega C')} = j\omega \sqrt{L'C'}$$

$$\gamma = \alpha + j\beta,$$

 $\beta = \omega \sqrt{L'C'} = 20 \text{ rad/m}$

16.360 Lecture 5

- lossless transmission line :

$$\begin{aligned}\gamma &= \alpha + j\beta, \\ &= \sqrt{(R' + j\omega L')(G' + j\omega C')}\end{aligned}$$

If $R' \ll j\omega L'$ and $G' \ll j\omega C'$,

$$\begin{aligned}\gamma &= \sqrt{(R' + j\omega L')(G' + j\omega C')} \\ &= j\omega \sqrt{L'C'}\end{aligned}$$

$$\alpha = 0$$

$$\beta = \omega \sqrt{L'C'}$$

lossless line

16.360 Lecture 5

- lossless transmission line :

$$Z_0 = \sqrt{\frac{(R' + j\omega L')}{(G' + j\omega C')}} = \sqrt{\frac{j\omega L'}{j\omega C'}}$$

$$Z_0 = \sqrt{\frac{L'}{C'}}$$

lossless line

$$\alpha = 0$$
$$\beta = \omega \sqrt{L'C'}$$

$$\beta = 2\pi/\lambda \quad \longrightarrow \quad \lambda = 2\pi/\beta = \frac{1}{\omega\sqrt{L'C'}}$$
$$v_p = \omega/\beta = \frac{1}{\sqrt{L'C'}}$$

16.360 Lecture 5

- For TEM transmission line :

$$L'C' = \mu\epsilon$$

$$V_p = \frac{1}{\sqrt{L'C'}} = \frac{1}{\sqrt{\mu\epsilon}} = \frac{c}{\sqrt{\mu_r\epsilon_r}}$$

$$\beta = \omega \sqrt{L'C'} = \omega \sqrt{\mu\epsilon}$$

- summary :

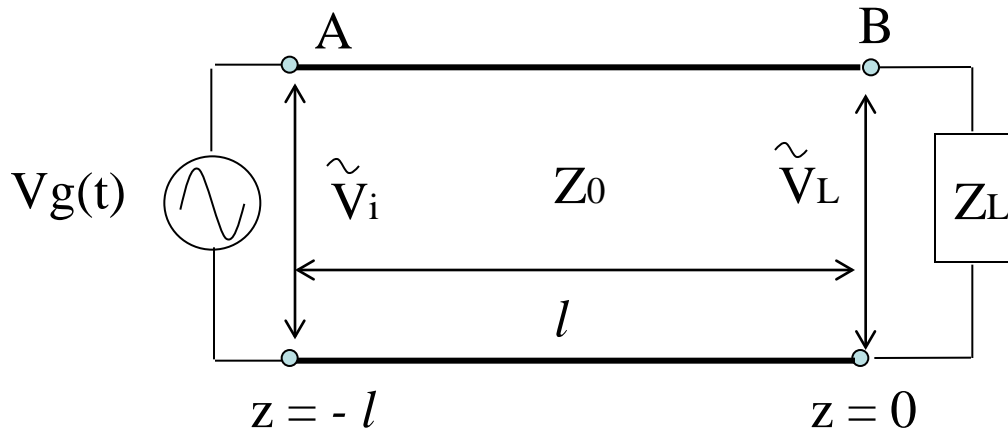
$$\begin{cases} \tilde{V}(z) = V_0^+ e^{-j\beta z} + V_0^- e^{j\beta z} \\ \tilde{i}(z) = I_0^+ e^{-j\beta z} + I_0^- e^{j\beta z} \end{cases}$$

$$Z_0 = \sqrt{\frac{L'}{C'}} \qquad V_p = \frac{1}{\sqrt{L'C'}} = \frac{c}{\sqrt{\mu_r\epsilon_r}}$$

$$\beta = \omega \sqrt{L'C'} = \omega \sqrt{\mu\epsilon}$$

16.360 Lecture 5

- Voltage reflection coefficient :



$$\begin{cases} \tilde{V}(z) = V_0^+ e^{-j\beta z} + V_0^- e^{j\beta z} \\ \tilde{i}(z) = \frac{V_0^+}{Z_0} e^{-j\beta z} - \frac{V_0^-}{Z_0} e^{j\beta z} \end{cases}$$

$$\tilde{V}_L = \tilde{V}(z) \Big|_{z=0} = V_0^+ + V_0^-$$

$$\tilde{i}_L = \tilde{i}(z) \Big|_{z=0} = \frac{V_0^+}{Z_0} - \frac{V_0^-}{Z_0}$$

$$Z_L = \frac{\tilde{V}_L}{\tilde{i}_L} = \frac{V_0^+ + V_0^-}{\frac{V_0^+}{Z_0} - \frac{V_0^-}{Z_0}} \quad \longrightarrow \quad \frac{V_0^+}{V_0^-} = \frac{Z_L - Z_0}{Z_L + Z_0}$$

16.360 Lecture 5

- Voltage reflection coefficient :

$$\Gamma \equiv \frac{V_0^-}{V_0^+} = \frac{Z_L - Z_0}{Z_L + Z_0}$$

- Current reflection coefficient :

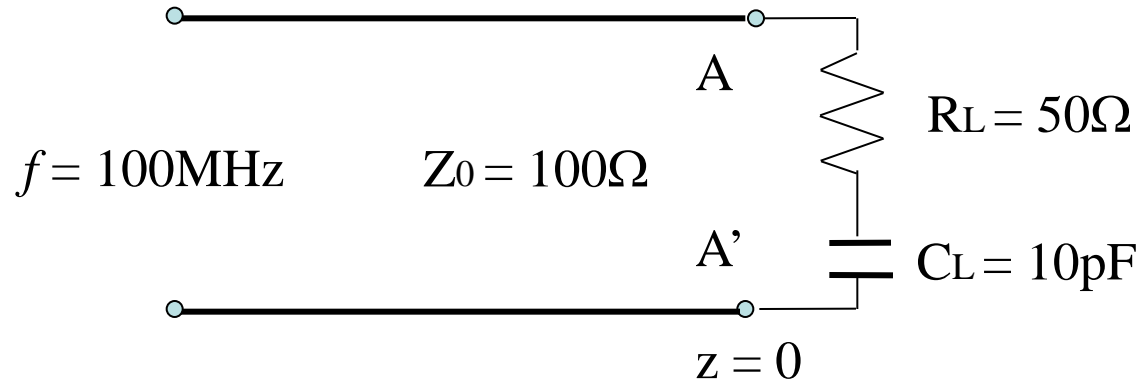
$$\Gamma_i \equiv \frac{i_0^-}{i_0^+} = - \frac{V_0^-}{V_0^+} = - \Gamma$$

- Notes :

1. $|\Gamma| \leq 1$, how to prove it?
2. If $Z_L = Z_0$, $\Gamma = 0$. Impedance match, no reflection from the load Z_L .

16.360 Lecture 5

- An example :



$$Z_L = R_L + j/\omega C_L = 50 - j159$$

$$\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0}$$

16.360 Lecture 6

- Standing wave
- Input impedance

$$\begin{cases} \tilde{V}(z) = V_0^+ e^{-j\beta z} + V_0^- e^{j\beta z} \\ \tilde{i}(z) = \frac{V_0^+}{Z_0} e^{-j\beta z} - \frac{V_0^-}{Z_0} e^{j\beta z} \end{cases} \quad \text{with } \Gamma = \frac{V_0^-}{V_0^+}$$

$$\begin{cases} \tilde{V}(z) = V_0^+ (e^{-j\beta z} + \Gamma e^{j\beta z}) \\ \tilde{i}(z) = \frac{V_0^+}{Z_0} (e^{-j\beta z} - \Gamma e^{j\beta z}) \end{cases}$$

$$|\tilde{V}(z)| = |V_0^+| |e^{-j\beta z} + |\Gamma| e^{j\theta_r} e^{j\beta z}|$$

$$= |V_0^+| [1 + |\Gamma|^2 + 2|\Gamma|\cos(2\beta z + \theta_r)]^{1/2}$$

16.360 Lecture 6

- Standing wave

$$\begin{cases} \tilde{V}(z) = V_0^+ e^{-j\beta z} + V_0^- e^{j\beta z} \\ \tilde{i}(z) = \frac{V_0^+}{Z_0} e^{-j\beta z} - \frac{V_0^-}{Z_0} e^{j\beta z} \end{cases} \quad \text{with } \Gamma = \frac{V_0^-}{V_0^+}$$

$$\begin{cases} \tilde{V}(z) = V_0^+ (e^{-j\beta z} + \Gamma e^{j\beta z}) \\ \tilde{i}(z) = \frac{V_0^+}{Z_0} (e^{-j\beta z} - \Gamma e^{j\beta z}) \end{cases}$$

$$|\tilde{V}(z)| = |V_0^+| |e^{-j\beta z} + |\Gamma| e^{j\theta_r} e^{j\beta z}|$$

$$= |V_0^+| [1 + |\Gamma|^2 + 2|\Gamma|\cos(2\beta z + \theta_r)]^{1/2}$$

16.360 Lecture 6

- Standing wave

$$\begin{cases} \tilde{V}(z) = V_0^+ (e^{-j\beta z} + \Gamma e^{j\beta z}) \\ \tilde{i}(z) = \frac{V_0^+}{Z_0} (e^{-j\beta z} - \Gamma e^{j\beta z}) \end{cases}$$

$$\begin{aligned} |\tilde{i}(z)| &= |V_0^+|/|Z_0| |e^{-j\beta z} - |\Gamma| e^{j\theta_r} e^{j\beta z}| \\ &= |V_0^+|/|Z_0| [1 + |\Gamma|^2 - 2|\Gamma|\cos(2\beta z + \theta_r)]^{1/2} \end{aligned}$$

$$|\tilde{V}(z)| = |V_0^+| [1 + |\Gamma|^2 + 2|\Gamma|\cos(2\beta z + \theta_r)]^{1/2}$$

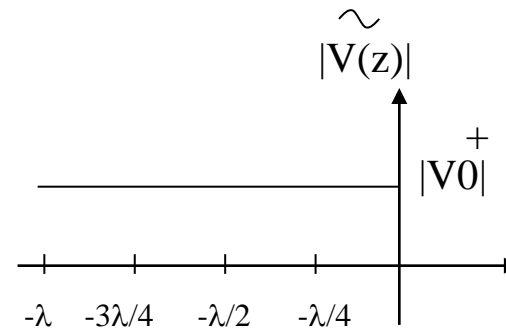
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$$|\tilde{V}(z)| = |V_0^+| [1 + |\Gamma|^2 + 2|\Gamma|\cos(2\beta z + \theta_r)]^{1/2}$$

Special cases

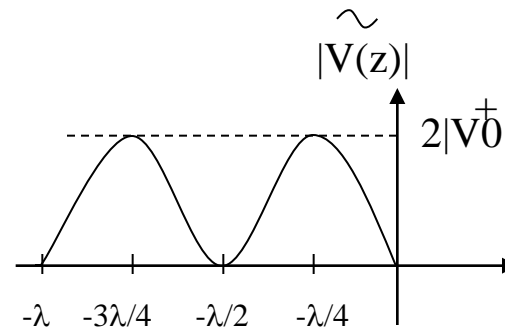
1. $Z_L = Z_0, \Gamma = 0$

$$|\tilde{V}(z)| = |V_0^+|$$



2. $Z_L = 0$, short circuit, $\Gamma = -1$

$$|\tilde{V}(z)| = |V_0^+| [2 + 2\cos(2\beta z + \pi)]^{1/2}$$



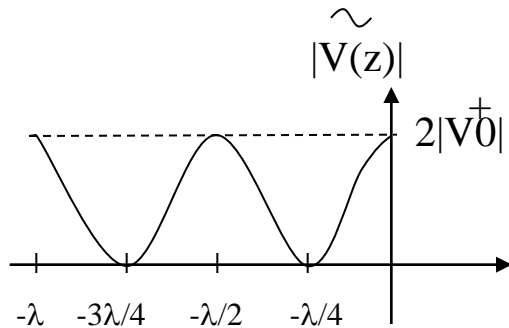
16.360 Lecture 6

$$|\tilde{V}(z)| = |V_0^+| [1 + |\Gamma|^2 + 2|\Gamma|\cos(2\beta z + \theta_r)]^{1/2}$$

Special cases

3. $Z_L = \infty$, open circuit, $\Gamma = 1$

$$|\tilde{V}(z)| = |V_0^+| [2 + 2\cos(2\beta z)]^{1/2}$$



16.360 Lecture 6

- Voltage maximum

$$|\tilde{V}(z)| = |V_0^+| [1 + |\Gamma|^2 + 2|\Gamma|\cos(2\beta z + \theta_r)]^{1/2}$$

$$|\tilde{V}(z)|_{\max} = |V_0^+| [1 + |\Gamma|], \quad \text{when } 2\beta z + \theta_r = 2n\pi.$$

$$-z = \lambda\theta_r/4\pi + n\lambda/2$$

$$n = 1, 2, 3, \dots, \text{ if } \theta_r < 0$$

$$n = 0, 1, 2, 3, \dots, \text{ if } \theta_r \geq 0$$

16.360 Lecture 6

- Voltage minimum

$$|\tilde{V}(z)| = |V_0^+| [1 + |\Gamma|^2 + 2|\Gamma|\cos(2\beta z + \theta_r)]^{1/2}$$

$$|\tilde{V}(z)|_{\min} = |V_0^+| [1 - |\Gamma|], \quad \text{when } 2\beta z + \theta_r = (2n+1)\pi.$$

$$-z = \lambda\theta_r/4\pi + n\lambda/2 + \lambda/4$$

Note:

voltage minimums occur $\lambda/4$ away from voltage maximum, because of the $2\beta z$, the special frequency doubled. _

16.360 Lecture 6

- Voltage standing-wave ratio VSWR or SWR

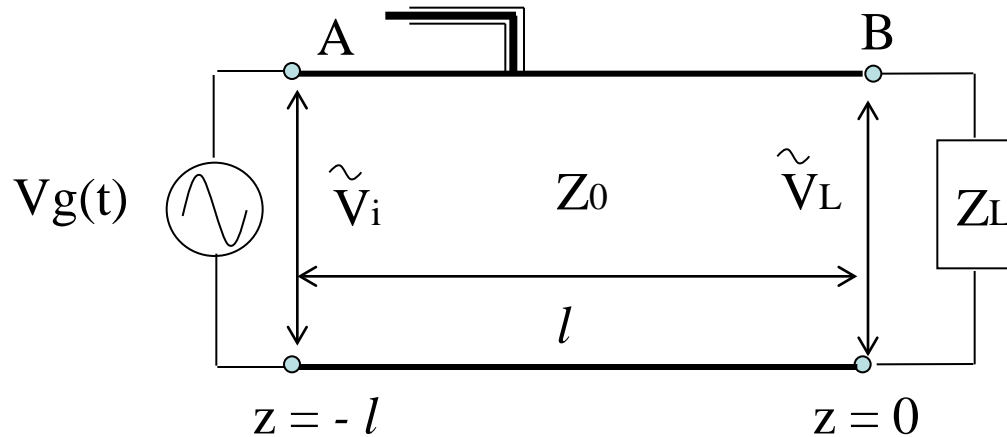
$$S \equiv \frac{|\tilde{V}(z)|_{\max}}{|\tilde{V}(z)|_{\min}} = \frac{1 + |\Gamma|}{1 - |\Gamma|}$$

$$S = 1, \text{ when } \Gamma = 0,$$

$$S = \infty, \text{ when } |\Gamma| = 1,$$

16.360 Lecture 6

- An example Voltage probe



$$S = 3, Z_0 = 50\Omega, \Delta l_{min} = 30cm, l_{min} = 12cm, Z_L = ?$$

Solution:

$$\Delta l_{min} = 30cm, \Rightarrow \lambda = 0.6m,$$

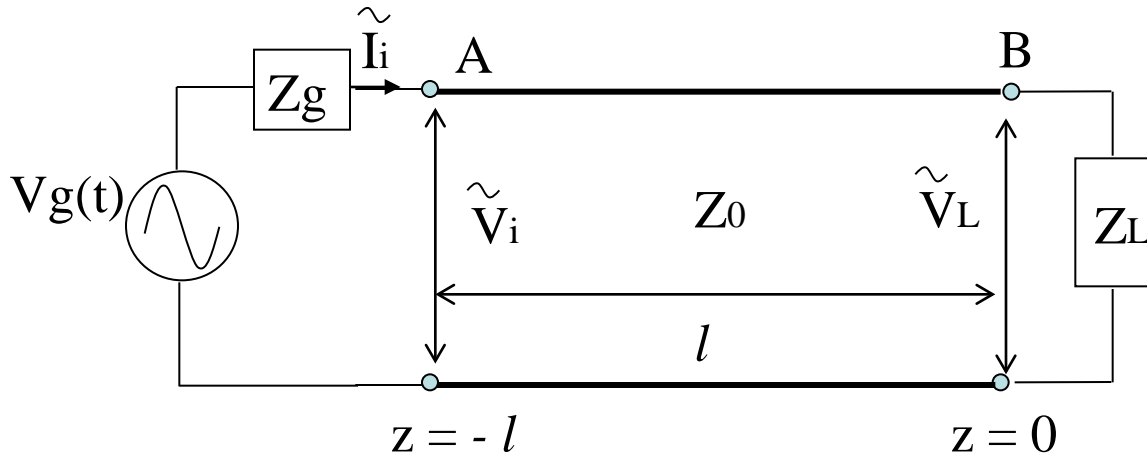
$$S = 3, \Rightarrow |\Gamma| = 0.5,$$

$$-2\beta l_{min} + \theta_r = -\pi, \Rightarrow \theta_r = -36^\circ,$$

$$\Rightarrow \Gamma, \text{ and } Z_L.$$

16.360 Lecture 6

- Input impedance



$$Z_{in}(z) = \frac{\tilde{V}(z)}{\tilde{I}(z)}$$

$$= \frac{v_0^+(e^{-j\beta z} + \Gamma e^{j\beta z})}{v_0^+(e^{-j\beta z} - \Gamma e^{j\beta z})} Z_0 = \frac{(1 + \Gamma e^{j2\beta z})}{(1 - \Gamma e^{j2\beta z})} Z_0$$

$$Z_{in}(-l) = \frac{(1 + \Gamma e^{-j2\beta l})}{(1 - \Gamma e^{-j2\beta l})} Z_0$$

16.360 Lecture 6

An example

A 1.05-GHz generator circuit with series impedance $Z_g = 10\text{-}\Omega$ and voltage source given by $V_g(t) = 10 \sin(\omega t + 30^\circ)$ is connected to a load $Z_L = 100 + j5\text{-}\Omega$ through a $50\text{-}\Omega$, 67-cm long lossless transmission line. The phase velocity is $0.7c$. Find $V(z,t)$ and $i(z,t)$ on the line.

Solution:

Since, $V_p = f\lambda$, $\lambda = V_p/f = 0.7c/1.05\text{GHz} = 0.2\text{m}$.

$$\beta = 2\pi/\lambda, \beta = 10 \pi.$$

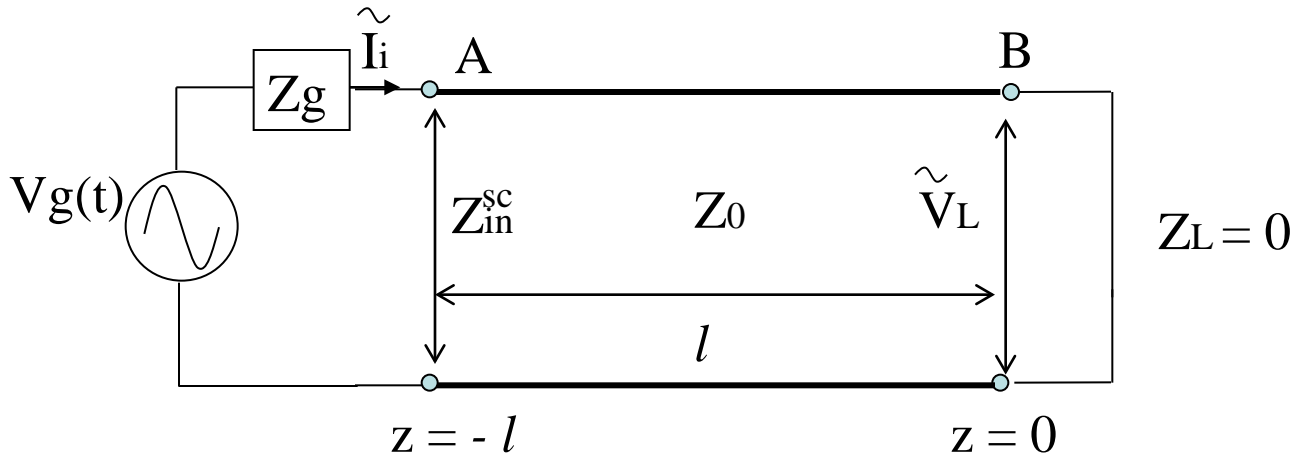
$$\Gamma = (Z_L - Z_0)/(Z_L + Z_0), \Gamma = 0.45 \exp(j26.6^\circ)$$

$$Z_{in}(-l) = \frac{(1 + \Gamma e^{-j2\beta l})}{(1 - \Gamma e^{-j2\beta l})} Z_0 = 21.9 + j17.4 \Omega$$

$$V_0^+ [\exp(-j\beta l) + \Gamma \exp(j\beta l)] = \frac{Z_{in}(-l)}{Z_{in}(-l) + Z_g} \tilde{V}_g$$

16.360 Lecture 7

short circuit line



$$Z_L = 0, \Gamma = -1, S = \infty$$

$$\begin{cases} \tilde{V}(z) = V_0(e^{-j\beta z} - e^{j\beta z}) = -2jV_0^+ \sin(\beta z) \\ \tilde{i}(z) = \frac{V_0^+}{Z_0}(e^{-j\beta z} + e^{j\beta z}) = 2V_0^+ \cos(\beta z)/Z_0 \end{cases}$$

$$Z_{in} = \frac{\tilde{V}(-l)}{\tilde{i}(-l)} = jZ_0 \tan(\beta l)$$

short circuit line

$$Z_{in} = \frac{\tilde{V}(-l)}{\tilde{i}(-l)} = jZ_0 \tan(\beta l)$$

- If $\tan(\beta l) \geq 0$, the line appears inductive, $j\omega L_{eq} = jZ_0 \tan(\beta l)$,
- If $\tan(\beta l) \leq 0$, the line appears capacitive, $1/j\omega C_{eq} = jZ_0 \tan(\beta l)$,
- The minimum length results in transmission line as a capacitor:

$$l = 1/\beta [\pi - \tan^{-1}(1/\omega C_{eq} Z_0)],$$

16.360 Lecture 7

An example:

Choose the length of a shorted $50\text{-}\Omega$ lossless line such that its input impedance at 2.25 GHz is equivalent to the reactance of a capacitor with capacitance $C_{eq} = 4\text{pF}$. The wave phase velocity on the line is $0.75c$.

Solution:

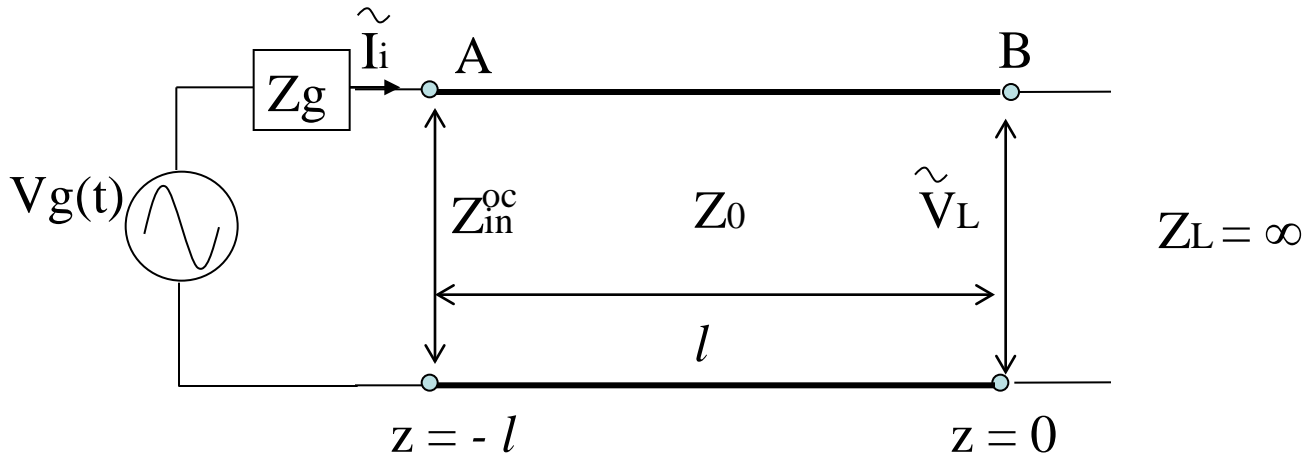
$$V_p = \lambda f, \Rightarrow \beta = 2\pi/\lambda = 2\pi f/V_p = 62.8 \text{ (rad/m)}$$

$$\tan(\beta l) = -1/\omega C_{eq} Z_0 = -0.354,$$

$$\begin{aligned} \beta l &= \tan^{-1}(-0.354) + n\pi, \\ &= -0.34 + n\pi, \end{aligned}$$

16.360 Lecture 7

open circuit line



$$Z_L = 0, \Gamma = 1, S = \infty$$

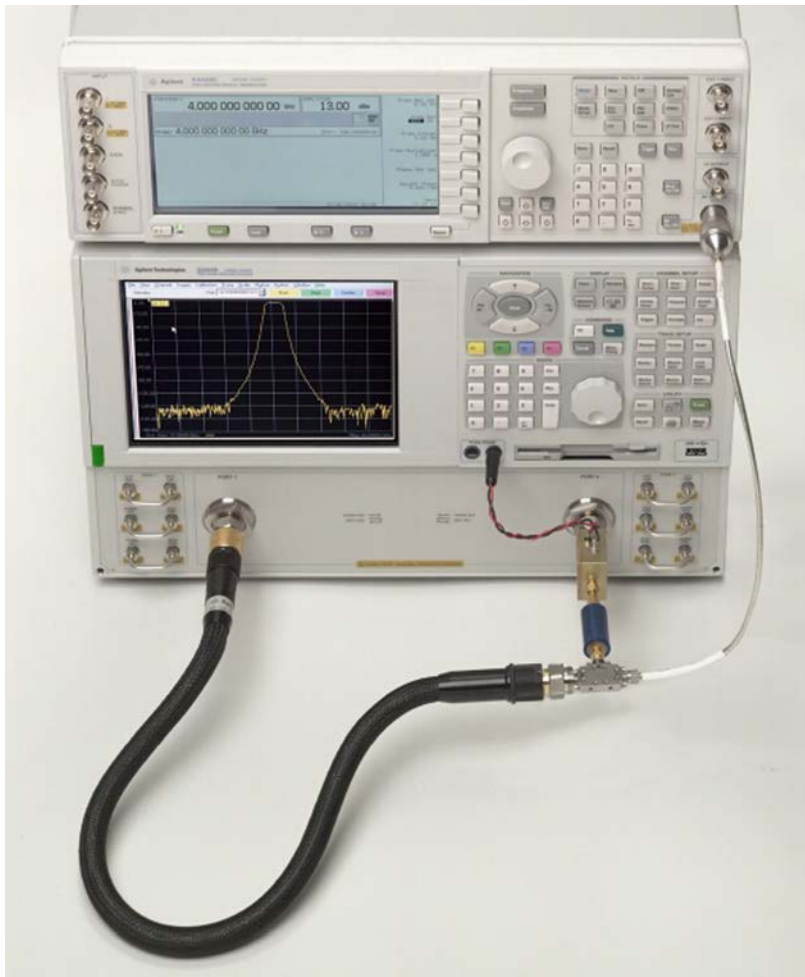
$$\begin{cases} \tilde{V}(z) = V_0(e^{-j\beta z} + e^{j\beta z}) = 2V_0^+ \cos(\beta z) \\ \tilde{i}(z) = \frac{V_0^+}{Z_0}(e^{-j\beta z} - e^{j\beta z}) = 2jV_0^+ \sin(\beta z)/Z_0 \end{cases}$$

$$Z_{in}^{oc} = \frac{\tilde{V}(-l)}{\tilde{i}(-l)} = -jZ_0 \cot(\beta l)$$

16.360 Lecture 7

Application for short-circuit and open-circuit

- Network analyzer



- Measure S parameters
- Measure Z_{in}^{sc} and Z_{in}^{oc}
- Calculate Z_0

$$Z_{in}^{sc} = jZ_0 \tan(\beta l)$$

$$Z_{in}^{oc} = -jZ_0 \cot(\beta l)$$

$$Z_0 = \sqrt{Z_{in}^{sc} Z_{in}^{oc}}$$

- Calculate βl

$$Z_0 = -j \sqrt{\frac{Z_{in}^{sc}}{Z_{in}^{oc}}}$$

16.360 Lecture 7

Line of length $l = n\lambda/2$

$$\tan(\beta l) = \tan((2\pi/\lambda)(n\lambda/2)) = 0,$$

$$Z_{in}(-l) = \frac{(1 + \Gamma e^{-j2\beta l})}{(1 - \Gamma e^{-j2\beta l})} Z_0 = Z_L$$

Any multiple of half-wavelength line doesn't modify the load impedance.

16.360 Lecture 7

Quarter-wave transformer $l = \lambda/4 + n\lambda/2$

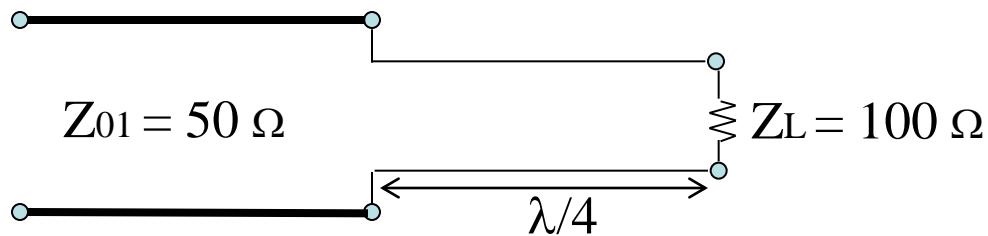
$$\beta l = (2\pi/\lambda)(\lambda/4 + n\lambda/2) = \pi/2 ,$$

$$\begin{aligned} Z_{\text{in}}(-l) &= \frac{(1 + \Gamma e^{-j2\beta l})}{(1 - \Gamma e^{-j2\beta l})} Z_0 = \frac{(1 + \Gamma e^{-j\pi})}{(1 - \Gamma e^{-j\pi})} Z_0 = \frac{(1 - \Gamma)}{(1 + \Gamma)} Z_0 \\ &= Z_0^2/Z_L \end{aligned}$$

16.360 Lecture 7

An example:

A $50\text{-}\Omega$ lossless transmission line is to be matched to a resistive load impedance with $Z_L = 100\ \Omega$ via a quarter-wave section, thereby eliminating reflections along the feed line. Find the characteristic impedance of the quarter-wave transformer.



$$Z_{in} = Z_0^2 / Z_L = 50\ \Omega$$

$$Z_0 = (Z_{in} Z_L)^{1/2} = (50 * 100)^{1/2}$$

16.360 Lecture 7

Matched transmission line:

1. $Z_L = Z_0$
2. $\Gamma = 0$
3. All incident power is delivered to the load.

16.360 Lecture 8

- Instantaneous power
- Time-average power

$$\begin{cases} \tilde{V}(z) = V_0^+(e^{-j\beta z} + \Gamma e^{j\beta z}) \\ \tilde{i}(z) = \frac{V_0^+}{Z_0} (e^{-j\beta z} - \Gamma e^{j\beta z}) \end{cases}$$

At load $z = 0$, the incident and reflected voltages and currents:

$$\tilde{V}^i = V_0^+ \quad \tilde{i}^i = \frac{V_0^+}{Z_0}$$

$$\tilde{V}^r = V_0^- \quad \tilde{i}^r = \frac{V_0^-}{Z_0}$$

16.360 Lecture 8

- Instantaneous power

$$\begin{aligned} P^i(t) &= v(t) i(t) = \operatorname{Re}[\tilde{V}^i \exp(j\omega t)] \operatorname{Re}[\tilde{i}^i \exp(j\omega t)] \\ &= \operatorname{Re}[|V_0^+| \exp(j\phi^+) \exp(j\omega t)] \operatorname{Re}[|V_0^+|/Z_0 \exp(j\phi^+) \exp(j\omega t)] \\ &= (|V_0^+|^2/Z_0) \cos^2(\omega t + \phi^+) \end{aligned}$$

$$\begin{aligned} P^r(t) &= v(t) i(t) = \operatorname{Re}[\tilde{V}^r \exp(j\omega t)] \operatorname{Re}[\tilde{i}^r \exp(j\omega t)] \\ &= \operatorname{Re}[|V_0^-| \exp(j\phi^+) \exp(j\omega t)] \operatorname{Re}[|V_0^-|/Z_0 \exp(j\phi^+) \exp(j\omega t)] \\ &= -|\Gamma|^2 (|V_0^+|^2/Z_0) \cos^2(\omega t + \phi^+ + \phi_r) \end{aligned}$$

16.360 Lecture 8

- Time-average

Time-domain approach:

$$\begin{aligned} P_{\text{av}}^{\text{i}} &= \frac{1}{T} \int_0^T P^{\text{i}}(t) dt = \frac{\omega}{2\pi} \int_0^T (|V_0^+|^2/Z_0) \cos^2(\omega t + \phi^+) dt \\ &= (|V_0^+|^2/2Z_0) \end{aligned}$$

$$P_{\text{av}}^{\text{r}} = -|\Gamma|^2 (|V_0^+|^2/2Z_0)$$

Net average power:

$$\begin{aligned} P_{\text{av}} &= P_{\text{av}}^{\text{i}} + P_{\text{av}}^{\text{r}} \\ &= (1-|\Gamma|^2) (|V_0^+|^2/2Z_0) \end{aligned}$$

16.360 Lecture 8

- Time-average

Phasor-domain approach

$$P_{\text{av}} = (1/2)\text{Re}[\tilde{V} \tilde{i}^*]$$

$$P_{\text{av}}^{\text{i}} = (1/2) \text{Re}[V_0^+ V_0^{+\ast} / Z_0] = (|V_0|^2 / 2Z_0)$$

$$P_{\text{av}}^{\text{r}} = -|\Gamma|^2 (|V_0^+|^2 / 2Z_0)$$

$$P_{\text{av}} = (1 - |\Gamma|^2) (|V_0|^2 / 2Z_0)$$

16.360 Lecture 9

Normalized admittance

$$y = \frac{Y}{Y_0} = \frac{G}{Y_0} + j \frac{B}{Y_0} = g + jb,$$

$$g = GZ_0, \quad b = BZ_0, \quad y = \frac{Y}{Y_0} = \frac{Z_0}{Z} = \frac{1}{z},$$

$$y_L = \frac{1}{Z_L} = \frac{1-\Gamma}{1+\Gamma}, \quad y_{in} = \frac{1}{Z_{in}} = \frac{1-\Gamma e^{-j\beta l}}{1+\Gamma e^{-j\beta l}},$$

$$z_{in} (l = \lambda / 4) = \frac{1-\Gamma}{1+\Gamma} = y_L,$$

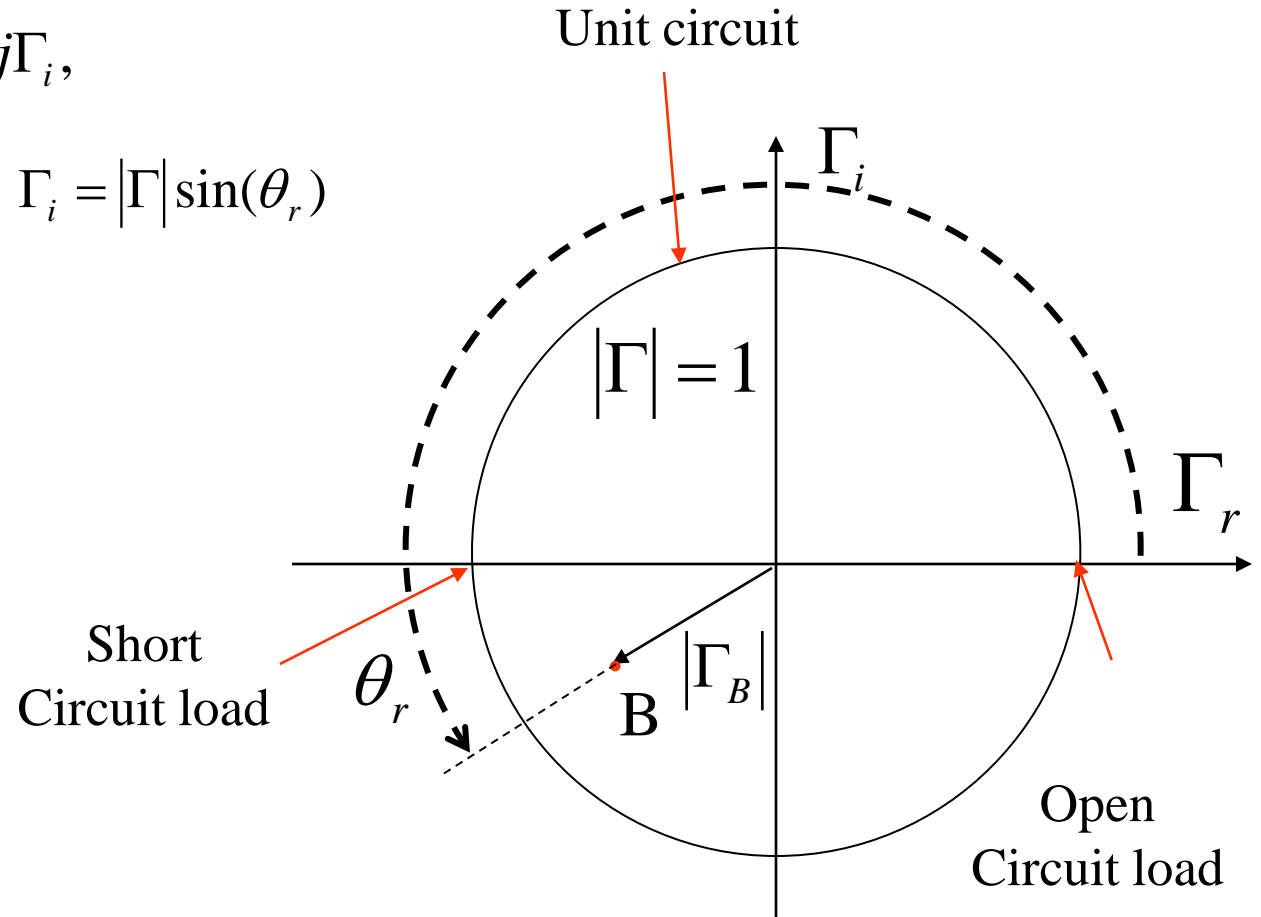
z and y are directly opposite each other on [Smith Chart](#)

16.360 Lecture 9

Parameter equations

$$\Gamma = |\Gamma|e^{j\theta_r} = \Gamma_r + j\Gamma_i,$$

$$\Gamma_r = |\Gamma| \cos(\theta_r) \quad \Gamma_i = |\Gamma| \sin(\theta_r)$$



16.360 Lecture 9

Normalized impedance

$$z_L = \frac{Z_L}{Z_0}, \quad \Gamma = \frac{Z_L - Z_0}{Z_L + Z_0} = \frac{Z_L/Z_0 - 1}{Z_L/Z_0 + 1} = \frac{z_L - 1}{z_L + 1},$$

$$z_L = \frac{1 + \Gamma}{1 - \Gamma}, \quad z_L = r_L + jx_L, \quad r_L + jx_L = \frac{(1 + \Gamma_r) + j\Gamma_i}{(1 - \Gamma_r) - j\Gamma_i}$$

Parameter equations

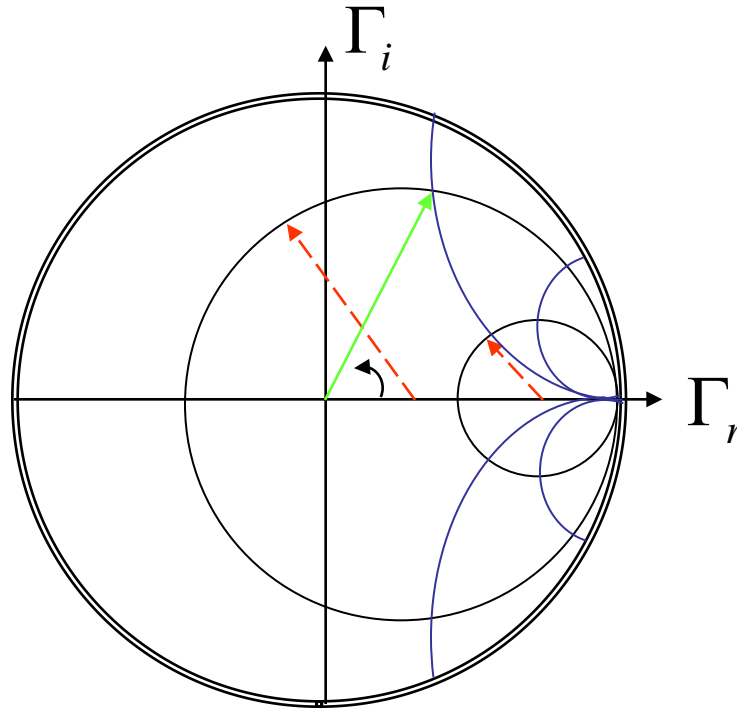
$$r_L = \frac{1 - \Gamma_r^2 - \Gamma_i^2}{(1 - \Gamma_r)^2 + \Gamma_i^2}, \quad x_L = \frac{2\Gamma_i}{(1 - \Gamma_r)^2 + \Gamma_i^2},$$

$$\left(\Gamma_r - \frac{r_L}{1 + r_L}\right)^2 + \Gamma_i^2 = 1, \quad (\Gamma_r - 1)^2 + \left(\Gamma_i - \frac{1}{x_L}\right)^2 = \left(\frac{1}{x_L}\right)^2,$$

16.360 Lecture 9

Parameter equations

$$\left(\Gamma_r - \frac{r_L}{1+r_L}\right)^2 + \Gamma_i^2 = 1, \quad (\Gamma_r - 1)^2 + \left(\Gamma_i - \frac{1}{x_L}\right)^2 = \left(\frac{1}{x_L}\right)^2,$$



16.360 Lecture 9

An example

$$z_L = 0.5 - j2,$$

Smith Chart

Input impedance

$$z_{in} = \frac{1 + \Gamma e^{-j\beta l}}{1 - \Gamma e^{-j\beta l}}, \quad \Gamma_l = \Gamma e^{-j\beta l} = |\Gamma| e^{j(\theta_r - \beta l)},$$

Wavelength toward generator (WTG)

Smith Chart

16.360 Lecture 9

An example

$$Z_L = 100 - j50, \quad \text{find } Z_{in}(-0.1\lambda)$$

$$z_L = 2 - j,$$

Smith Chart

$$S = \frac{1 + |\Gamma|}{1 - |\Gamma|}, \quad \text{Constant } |\Gamma| \text{ circle, SWR Circle}$$

16.360 Lecture 10

SWR, voltage maximum and minimum

$$z_L = r_L + jx_L, \quad \text{If } x_L = 0,$$

Smith Chart

$$S = \frac{1 + |\Gamma|}{1 - |\Gamma|}, \quad \Gamma_L = \frac{z_L - 1}{z_L + 1} = \frac{r_L - 1}{r_L + 1}, \quad r_L = \frac{1 + \Gamma_r}{1 - \Gamma_r}, \quad S = r_L,$$

Recall:

$$|\tilde{V}(z)|_{\max} = |V_0^+| [1 + |\Gamma|], \quad \text{when } 2\beta z + \theta_r = 2n\pi.$$

$$|\tilde{V}(z)|_{\min} = |V_0^+| [1 - |\Gamma|], \quad \text{when } 2\beta z + \theta_r = (2n+1)\pi.$$

16.360 Lecture 10

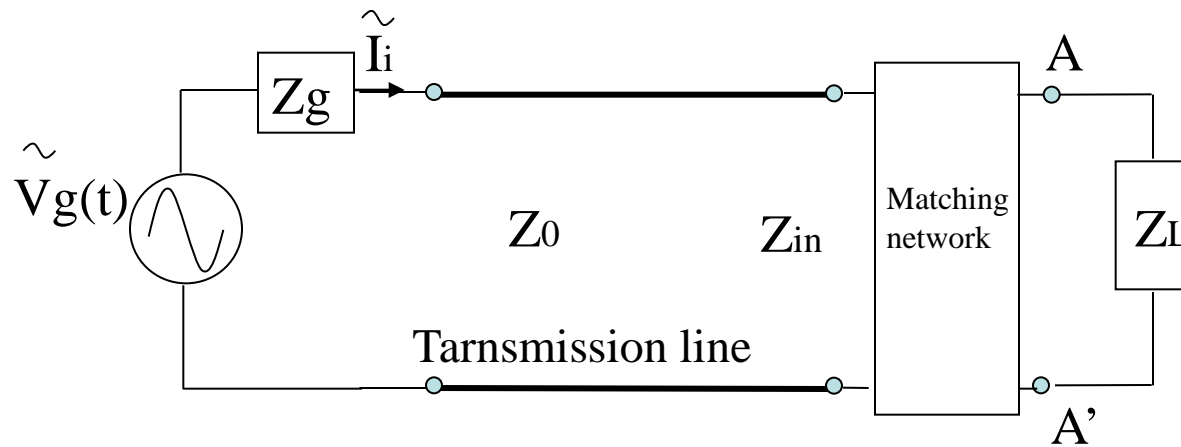
An example

A $50\text{-}\Omega$ lossless line is terminated in a load $Z_L = (25 + j50)\Omega$. Use the smith chart to find a) voltage reflection coefficient, b) the voltage standing-wave ratio, c) the distances of the first voltage maximum and first voltage minimum from the load, d) the input impedance of the line, given the line is 3.3λ , and e) the input admittance of the line.

Smith Chart

16.360 Lecture 11

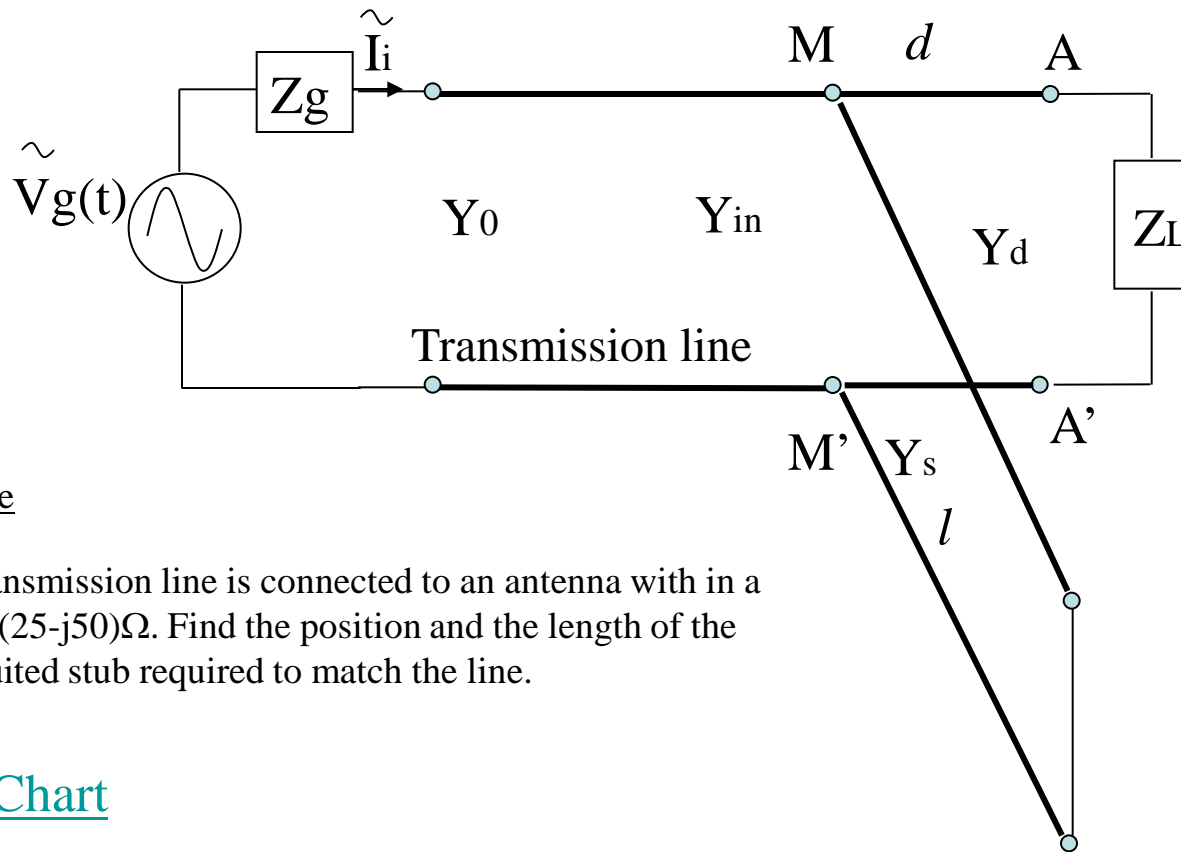
- impedance matching



$$Z_{in} = Z_0$$

16.360 Lecture 11

- single-stub impedance matching network



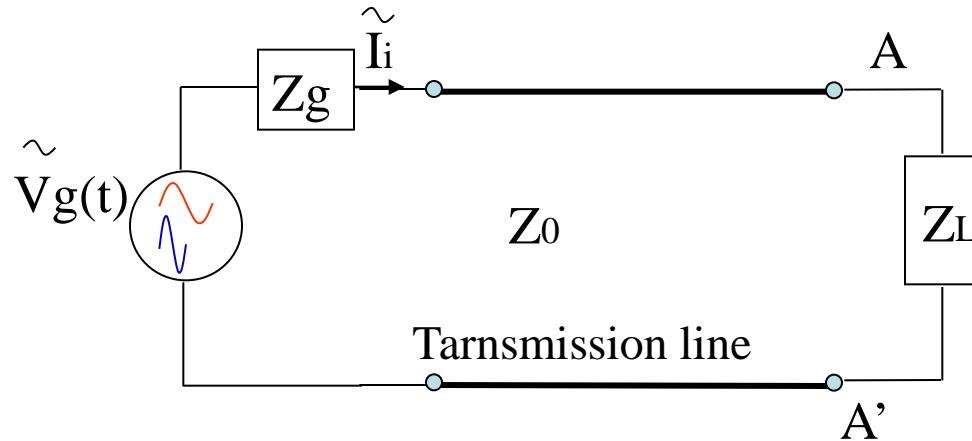
An example

A $50\text{-}\Omega$ transmission line is connected to an antenna with in a load $Z_L = (25 - j50)\Omega$. Find the position and the length of the short-circuited stub required to match the line.

Smith Chart

16.360 Lecture 12

- Transient on transmission line



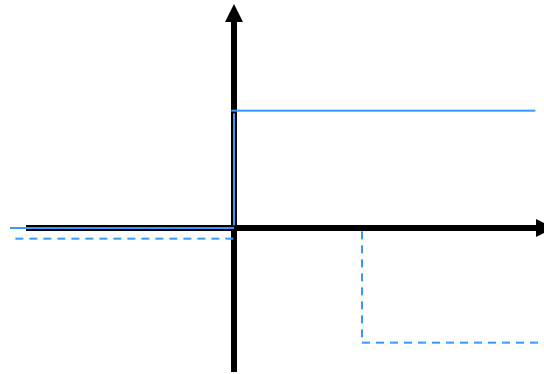
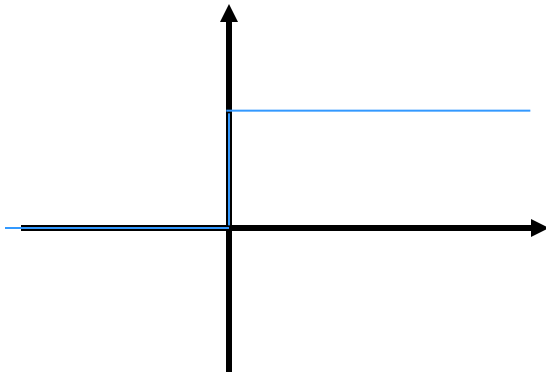
If $\omega_1, \omega_2, \dots, \omega_n$ are transmitted on the transmission line at the same time, each frequency has its own location of voltage distribution. The total voltage $V(z)$ is the sum of all these $V_{\omega_i}(z)$.

$$V(z) = \sum_{i=\omega_1}^{\omega_n} V_{\omega_i}(z)$$

Step function and pulse function

- step function $U(t)$

$$U(t) = 1, \text{ if } t \geq 0; U(t) = 0, \text{ if } t < 0$$

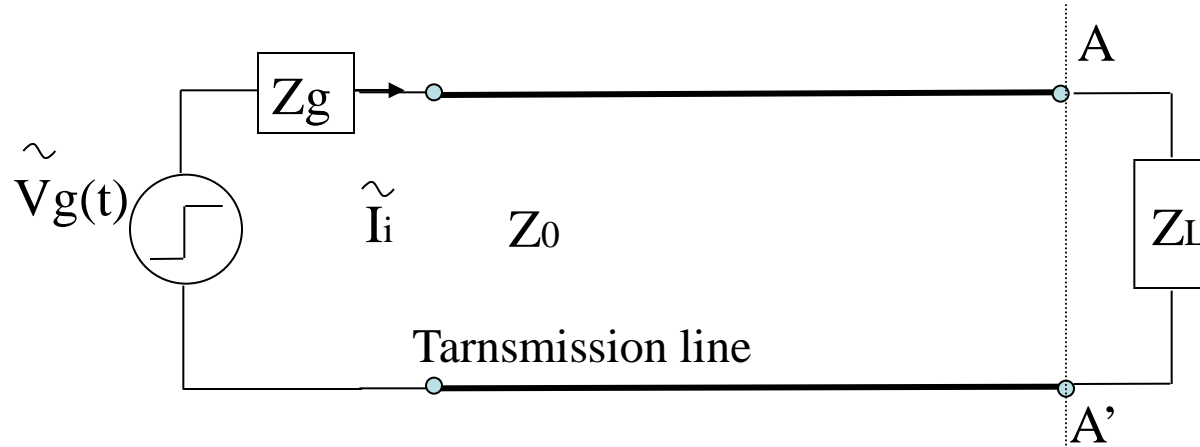


- single pulse function $V(t)$

$$V(t) = U(t) - U(t-t_0),$$

16.360 Lecture 12

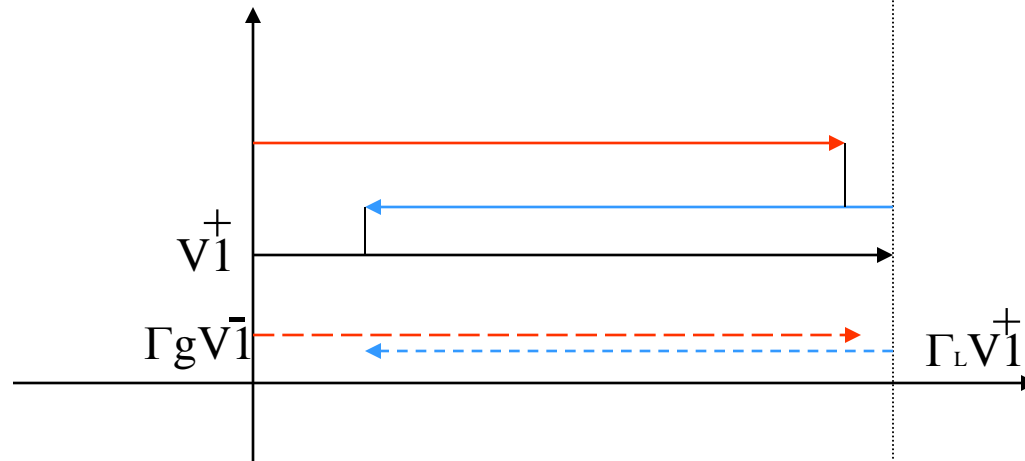
- transient of a step function



$$V\bar{1} = \Gamma_L V1^+$$

$$V2^+ = \Gamma_g V\bar{1}$$

$$V\bar{2} = \Gamma_L V2^+$$



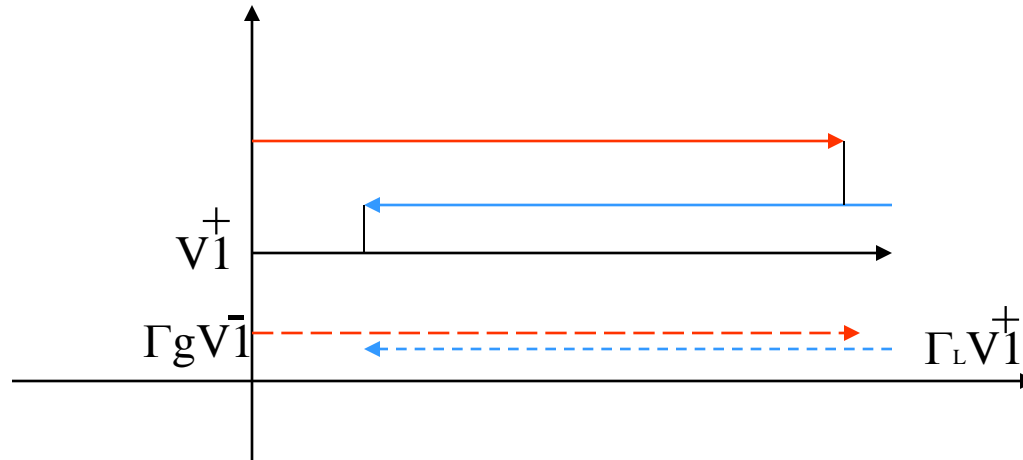
$$V = V1^+ + V\bar{1} + V2^+ + V\bar{2} + \dots$$

16.360 Lecture 12

$$V_1^- = \Gamma_L V_1^+$$

$$V_2^+ = \Gamma_g V_1^-$$

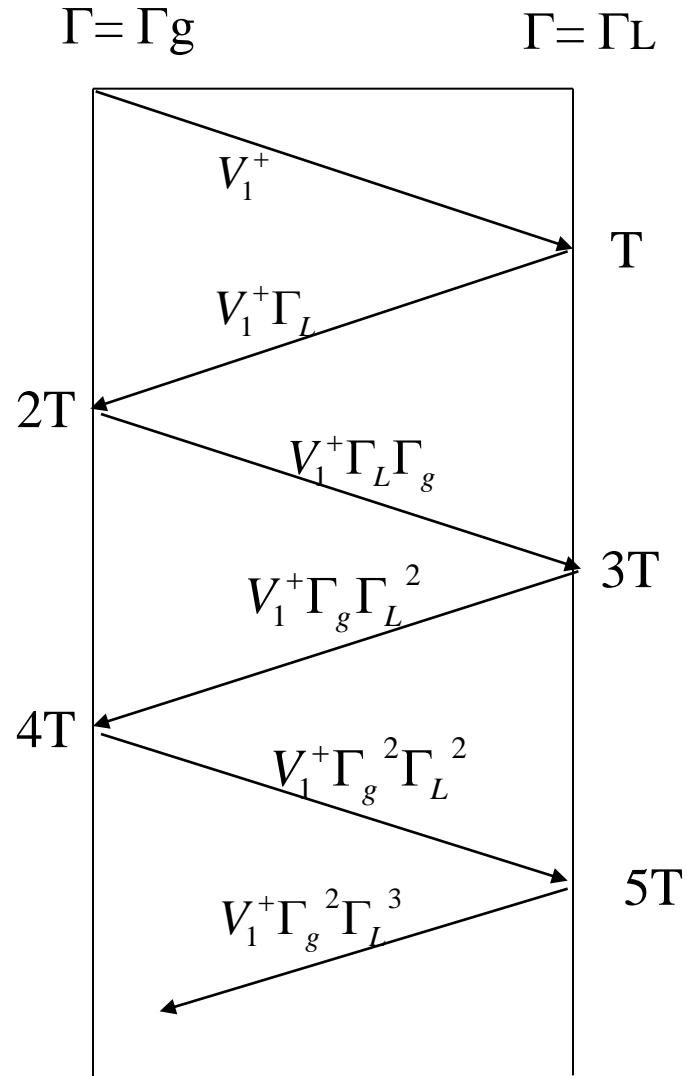
$$V_2^- = \Gamma_L V_2^+$$



$$\begin{aligned}
 V &= V_1^+ + V_1^- + V_2^+ + V_2^- + \dots \\
 &= V_1^+ + \Gamma_L V_1^+ + \Gamma_g (\Gamma_L V_1^+) + \Gamma_L (\Gamma_g (\Gamma_L V_1^+)) + \dots \\
 &= V_1^+ (1 + \Gamma_L) + \Gamma_g \Gamma_L (1 + \Gamma_L V_1^+) + \dots \\
 &= V_1^+ (1 + \Gamma_L) [1 + \Gamma_g \Gamma_L + (\Gamma_g \Gamma_L)^2 + \dots] \\
 &= V_1^+ (1 + \Gamma_L) \left[\frac{1}{1 - \Gamma_g \Gamma_L} \right]
 \end{aligned}$$

16.360 Lecture 12

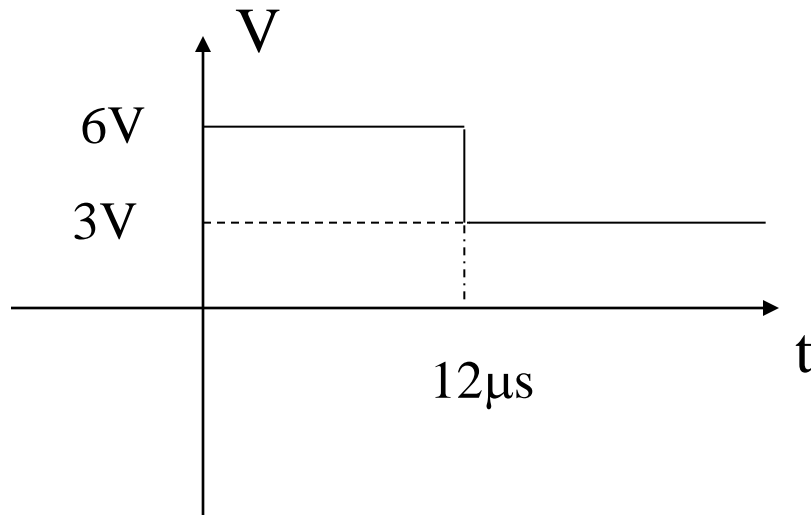
- Bounce Diagram



16.360 Lecture 12

- An example

$$Z_0 = 75\Omega, \epsilon_r = 2.1, V_g = ?, Z_{lf} = ?, L_f = ?$$



16.360 Lecture 13

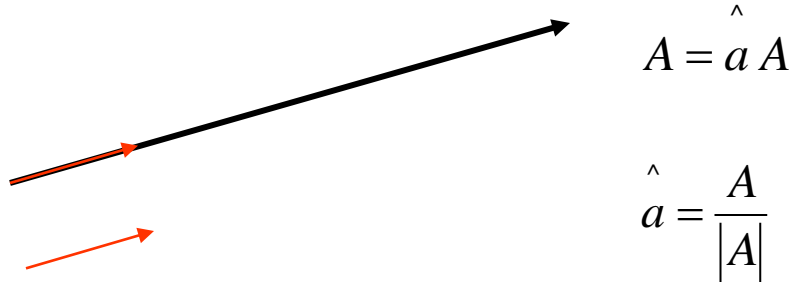
Basic Laws of Vector Algebra

Scalars:

e.g. 2 gallons, \$1,000, 35°C

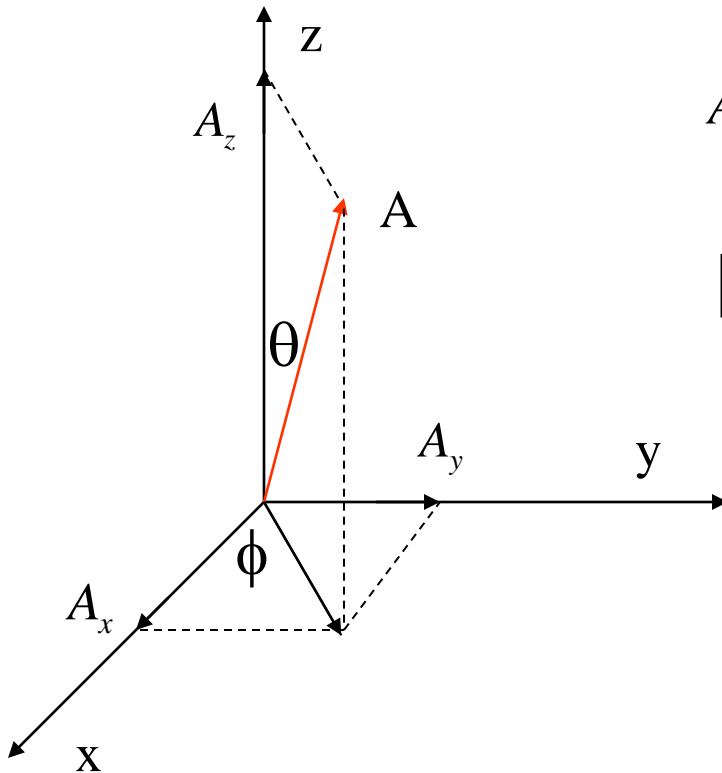
Vectors:

e.g. velocity: 35mph heading south
3N force toward center



16.360 Lecture 13

- Cartesian coordinate system



$$A = A_x \hat{x} + A_y \hat{y} + A_z \hat{z}$$

$$|A| = \sqrt{A_x^2 + A_y^2 + A_z^2}$$

$$A_z = |A| \cos(\theta)$$

$$A_x = |A| \sin(\theta) \cos(\phi)$$

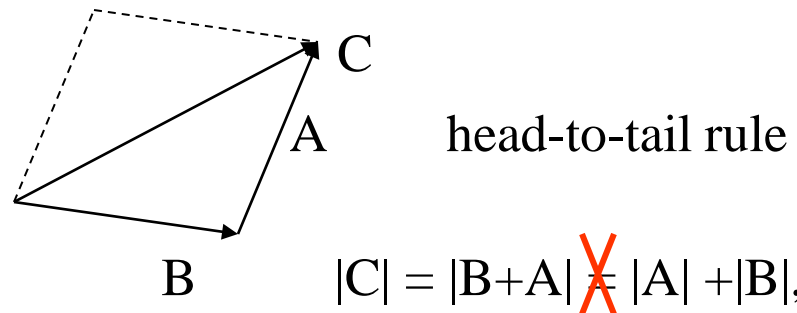
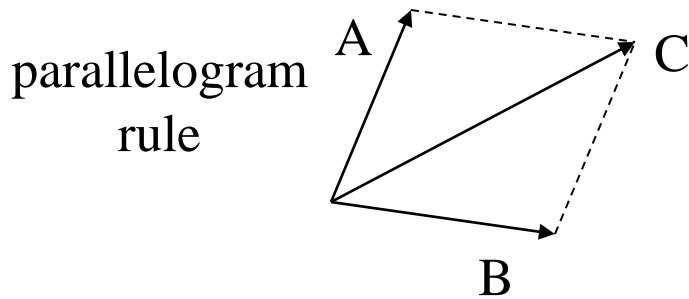
$$A_y = |A| \sin(\theta) \sin(\phi)$$

16.360 Lecture 13

- Vector addition and subtraction

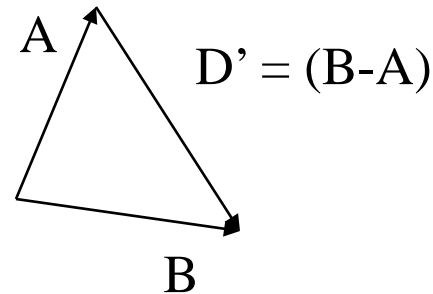
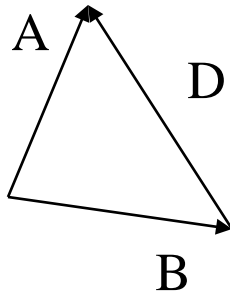
$$\mathbf{C} = \mathbf{B} + \mathbf{A} = \mathbf{A} + \mathbf{B},$$

$$= (\hat{x}A_x + \hat{y}A_y + \hat{z}A_z) + (\hat{x}B_x + \hat{y}B_y + \hat{z}B_z) = \hat{x}(A_x + B_x) + \hat{y}(A_y + B_y) + \hat{z}(A_z + B_z)$$



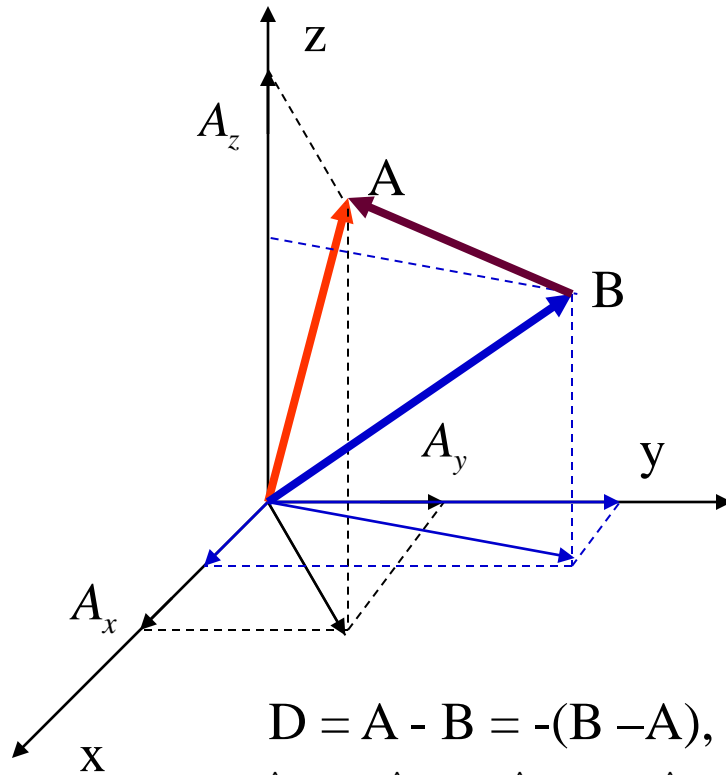
$$\mathbf{D} = \mathbf{A} - \mathbf{B} = -(\mathbf{B} - \mathbf{A}),$$

$$= (\hat{x}A_x + \hat{y}A_y + \hat{z}A_z) - (\hat{x}B_x + \hat{y}B_y + \hat{z}B_z) = \hat{x}(A_x - B_x) + \hat{y}(A_y - B_y) + \hat{z}(A_z - B_z)$$



16.360 Lecture 13

- position and distance



$$\begin{aligned} \mathbf{D} &= \mathbf{A} - \mathbf{B} = -(\mathbf{B} - \mathbf{A}), \\ &= (\hat{x} A_x + \hat{y} A_y + \hat{z} A_z) - (\hat{x} B_x + \hat{y} B_y + \hat{z} B_z) = \hat{x}(A_x - B_x) + \hat{y}(A_y - B_y) + \hat{z}(A_z - B_z) \end{aligned}$$

$$|\mathbf{D}| = \sqrt{(A_x - B_x)^2 + (A_y - B_y)^2 + (A_z - B_z)^2}$$

16.360 Lecture 13

- Vector multiplication

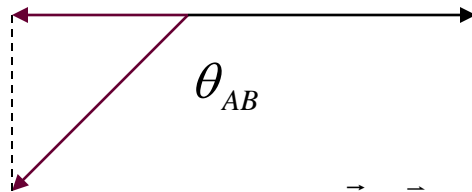
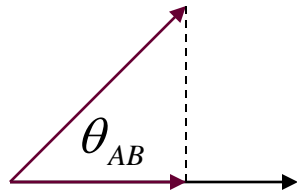
1. simple product

$$\vec{B} = k\vec{A} = \hat{x}(kA_x) + \hat{y}(kA_y) + \hat{z}(kA_z)$$

$$\vec{A} = \hat{x}(A_x) + \hat{y}(A_y) + \hat{z}(A_z)$$

2. scalar product (dot product)

$$\vec{B} \bullet \vec{A} = |\vec{A}||\vec{B}| \cos \theta_{AB}$$



$$\vec{B} = \hat{x}(B_x) + \hat{y}(B_y) + \hat{z}(B_z)$$
$$\vec{A} = \hat{x}(A_x) + \hat{y}(A_y) + \hat{z}(A_z)$$

$$\vec{B} \bullet \vec{A} = [\hat{x}(B_x) + \hat{y}(B_y) + \hat{z}(B_z)] \bullet [\hat{x}(A_x) + \hat{y}(A_y) + \hat{z}(A_z)]$$
$$= [\hat{x}(B_x A_x) + \hat{y}(B_y A_y) + \hat{z}(B_z A_z)]$$

Properties of scalar product (dot product)

a) commutative property

$$\vec{B} \bullet \vec{A} = \vec{A} \bullet \vec{B} \quad \vec{A} \bullet \vec{A} = |\vec{A}|^2$$

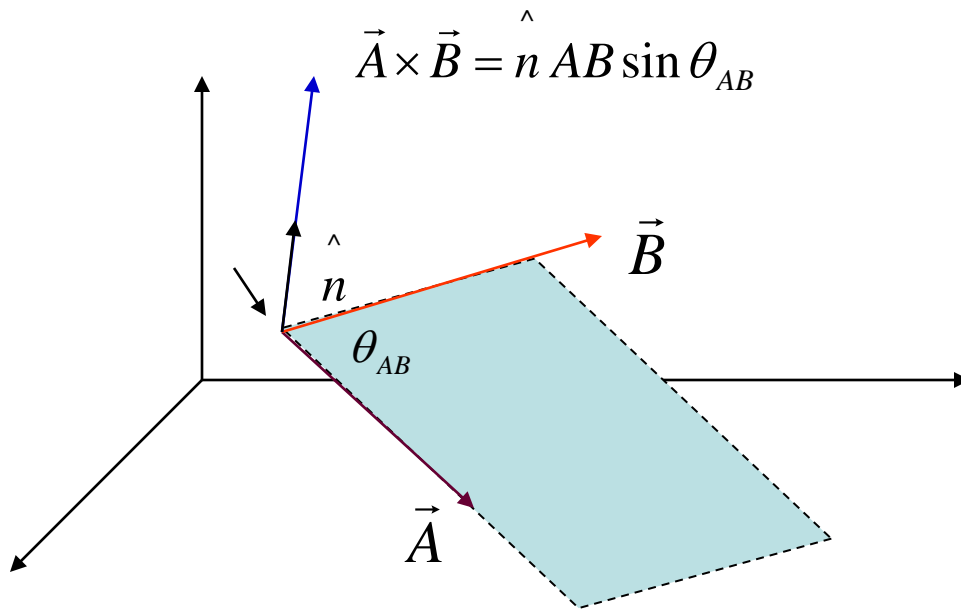
b) Distributive property

$$\vec{B} \bullet (\vec{A} + \vec{C}) = \vec{B} \bullet \vec{A} + \vec{B} \bullet \vec{C}$$

16.360 Lecture 13

3. vector product (cross product)

$$\vec{A} \times \vec{B} = \hat{n} AB \sin \theta_{AB}$$



a) anticommutative property

$$\vec{A} \times \vec{B} = -\vec{B} \times \vec{A}$$

b) Distributive property

$$\vec{A} \times (\vec{B} + \vec{C}) = \vec{A} \times \vec{B} + \vec{A} \times \vec{C}$$

c) $\vec{A} \times \vec{A} = 0$

$\hat{x} \times \hat{y} = \hat{z},$	$\hat{y} \times \hat{z} = \hat{x},$
$\hat{z} \times \hat{x} = \hat{y},$	

16.360 Lecture 13

3. vector product (cross product)

$$\begin{aligned}\vec{A} \times \vec{B} &= (\hat{x} A_x + \hat{y} A_y + \hat{z} A_z) \times (\hat{x} B_x + \hat{y} B_y + \hat{z} B_z) \\ &= \hat{x}(A_y B_z - A_z B_y) + \hat{y}(A_z B_x - A_x B_z) + \hat{z}(A_x B_y - A_y B_x)\end{aligned}$$

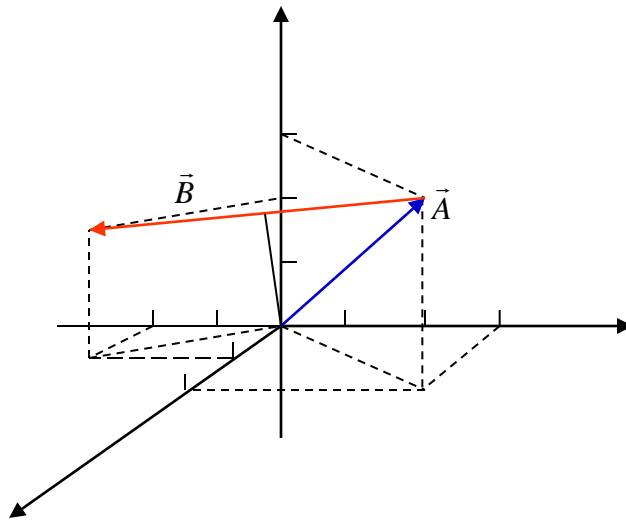
$$\vec{A} \times \vec{B} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ A_x & A_y & A_z \\ B_x & B_y & B_z \end{vmatrix}$$

16.360 Lecture 13

Example vectors and angles

In Cartesian coordinate, vector A is directed from origin to point P1(2,3,3), and vector B is directed from P1 to point P2(1,-2,2). Find:

- Vector A, its magnitude $|A|$, and unit vector \hat{a}
- the angle that A makes with the y-axis
- Vector B
- the angle between A and B
- perpendicular distance from origin to vector B



$$\vec{A} = \hat{x}2 + \hat{y}3 + \hat{z}3 \quad |A| = \sqrt{2^2 + 3^2 + 3^2}$$
$$\hat{a} = \frac{\vec{A}}{|A|} = (\hat{x}2 + \hat{y}3 + \hat{z}3) / \sqrt{22}$$

$$\vec{A} \cdot \hat{y} = |A| |\hat{y}| \cos(\theta_{Ay})$$

$$\vec{B} = \vec{P}_2 - \vec{P}_1 = \hat{x}(1-2) + \hat{y}(-2-3) + \hat{z}(2-3)$$

$$\vec{A} \cdot \vec{B} = |A| |B| \cos(\theta_{AB})$$

$$d = |A| \sin(180^\circ - \theta_{AB})$$

16.360 Lecture 13

4. Scalar and vector triple product

a) scalar triple product

$$\vec{A} \bullet (\vec{B} \times \vec{C}) = \vec{B} \bullet (\vec{C} \times \vec{A}) = \vec{C} \bullet (\vec{A} \times \vec{B})$$

$$\vec{A} \bullet (\vec{B} \times \vec{C}) = \begin{vmatrix} A_x & A_y & A_z \\ B_x & B_y & B_z \\ C_x & C_y & C_z \end{vmatrix}$$

$$\vec{A} \times (\vec{B} \times \vec{C}) = \vec{B}(\vec{A} \bullet \vec{C}) - \vec{C}(\vec{A} \bullet \vec{B}) \times \begin{vmatrix} A_x & A_y & A_z \\ B_x & B_y & B_z \\ C_x & C_y & C_z \end{vmatrix} \vec{B} \bullet (\vec{C} \times \vec{A}) = \vec{C} \bullet (\vec{A} \times \vec{B})$$

b) vector triple product

$$\vec{A} \times (\vec{B} \times \vec{C}) \neq (\vec{A} \times \vec{B}) \times \vec{C}$$

$$\vec{A} \times (\vec{B} \times \vec{C}) = \vec{B}(\vec{A} \bullet \vec{C}) - \vec{C}(\vec{A} \bullet \vec{B})$$

16.360 Lecture 13

Example vector triple product

$$\vec{A} = \hat{x} - \hat{y} + \hat{z} \quad \vec{B} = \hat{y} + \hat{z} \quad \vec{C} = -2\hat{x} + \hat{z}$$

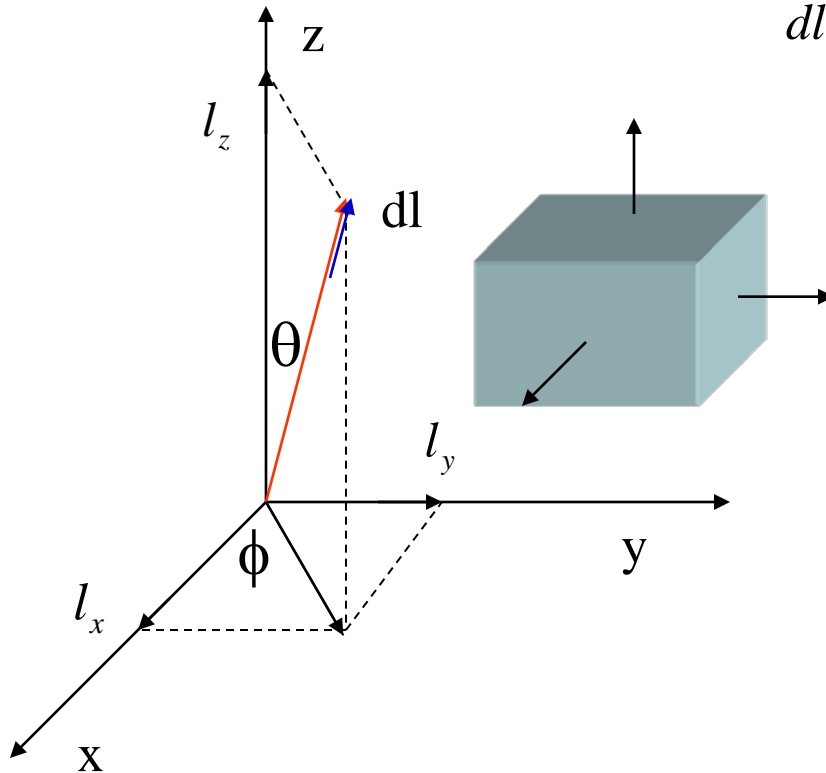
$$\vec{A} \times (\vec{B} \times \vec{C}) = ? \quad (\vec{A} \times \vec{B}) \times \vec{C} = ?$$

$$\vec{B} \times \vec{C} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ B_x & B_y & B_z \\ C_x & C_y & C_z \end{vmatrix} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 0 & 1 & 1 \\ -2 & 0 & 3 \end{vmatrix} = \hat{x}3 - 2\hat{y} + 2\hat{z}$$

$$\vec{A} \times (\vec{B} \times \vec{C}) = 3 \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 1 & -1 & 2 \\ -2 & -2 & 2 \end{vmatrix} = \hat{x}2 - 6\hat{y} + \hat{z}0$$

16.360 Lecture 14

- Cartesian coordinate system



$$d\vec{l} = dl_x \hat{x} + dl_y \hat{y} + dl_z \hat{z} = dx \hat{x} + dy \hat{y} + dz \hat{z}$$

$$d\vec{s}_x = \hat{x} dl_y dl_z = x dy dz$$

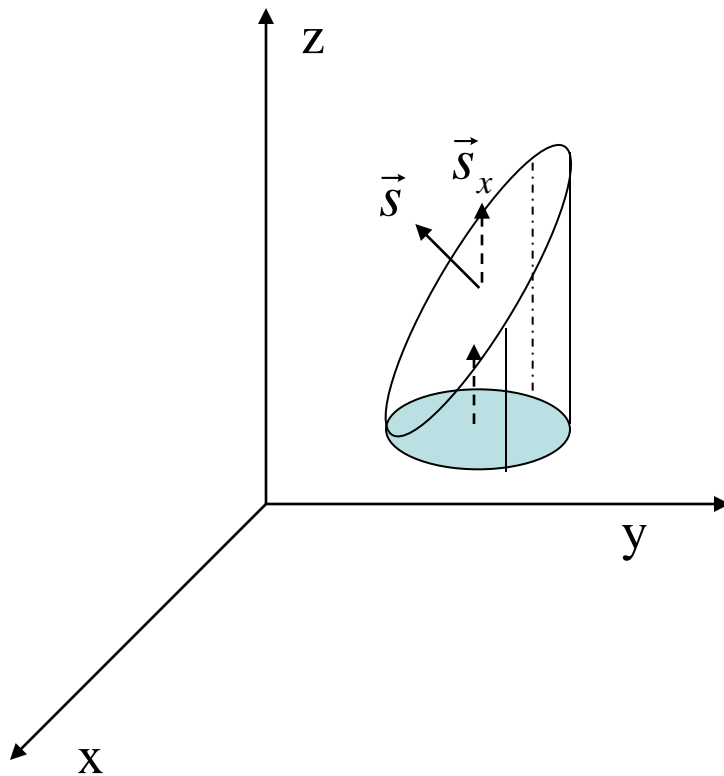
$$d\vec{s}_y = \hat{y} dl_x dl_z = y dx dz$$

$$d\vec{s}_z = \hat{z} dl_x dl_y = z dx dy$$

$$dv = dx dy dz$$

16.360 Lecture 14

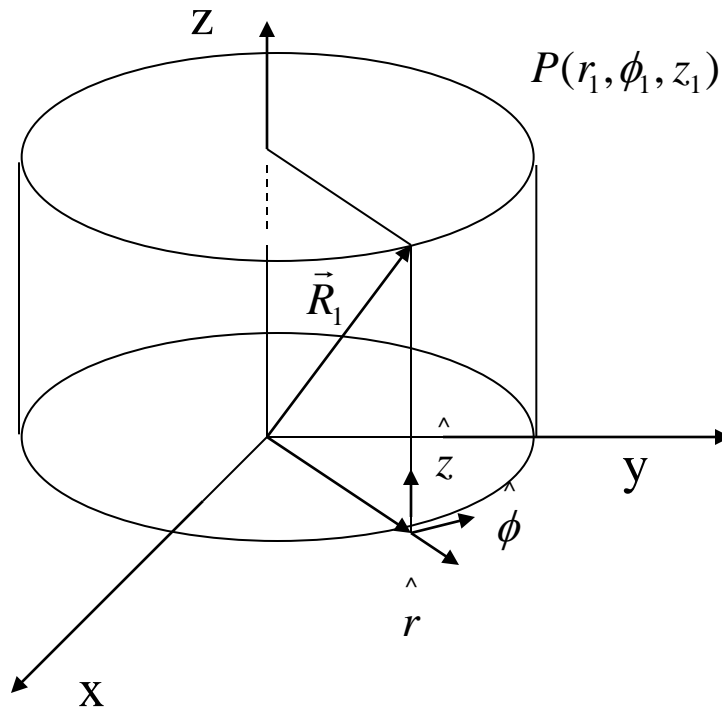
- Cartesian coordinate system



directions of area

16.360 Lecture 14

- Cylindrical coordinate system



$\hat{r} \times \hat{\phi} = \hat{z}$	$\hat{\phi} \times \hat{z} = \hat{r}$	$\hat{z} \times \hat{r} = \hat{\phi}$
---------------------------------------	---------------------------------------	---------------------------------------

$$\hat{r} \times \hat{r} = \mathbf{0} \quad \hat{r} \cdot \hat{r} = 1$$

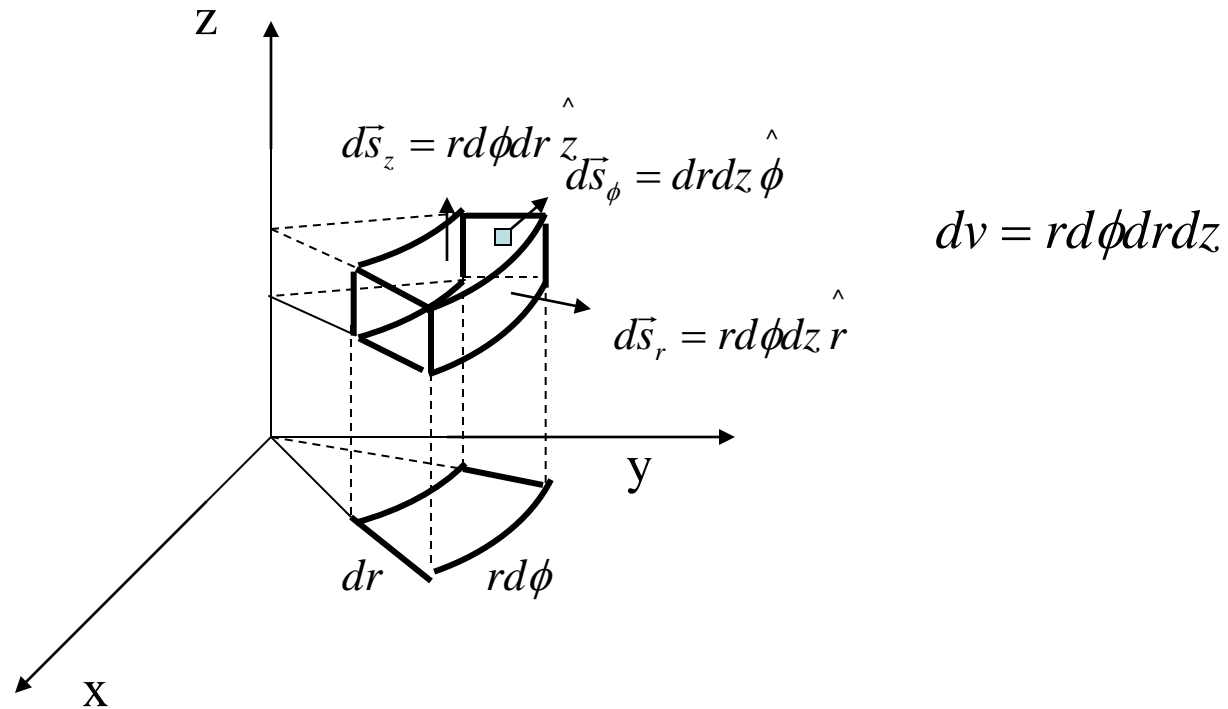
$$\vec{A} = \hat{r} A_r + \hat{\phi} A_\phi + \hat{z} A_z$$

$$|\vec{A}| = \sqrt{A_r^2 + A_\phi^2 + A_z^2}$$

$$\vec{R}_1 = \hat{r} r_1 + \hat{z} z_1$$

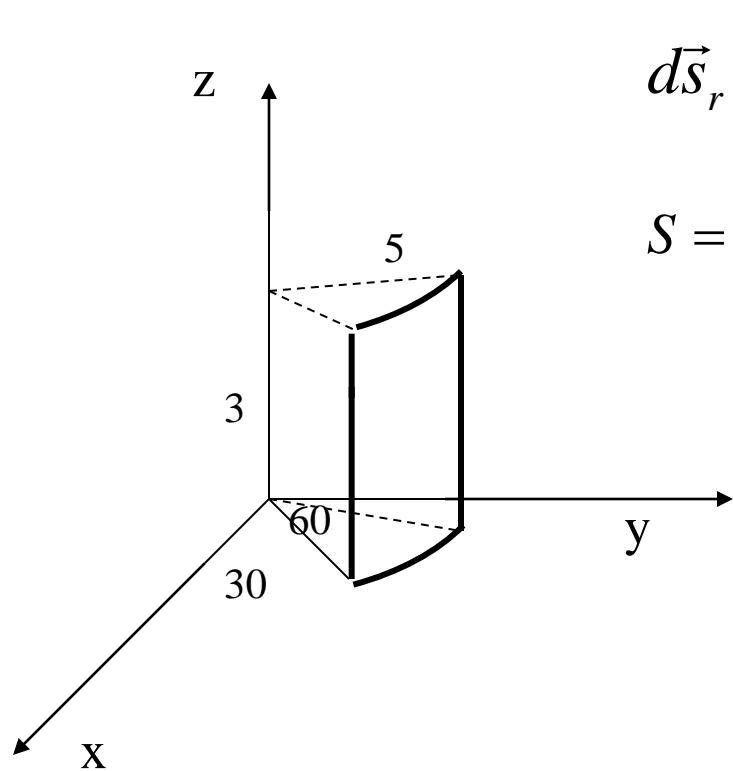
16.360 Lecture 14

- the differential areas and volume



16.360 Lecture 14

Example: cylindrical area

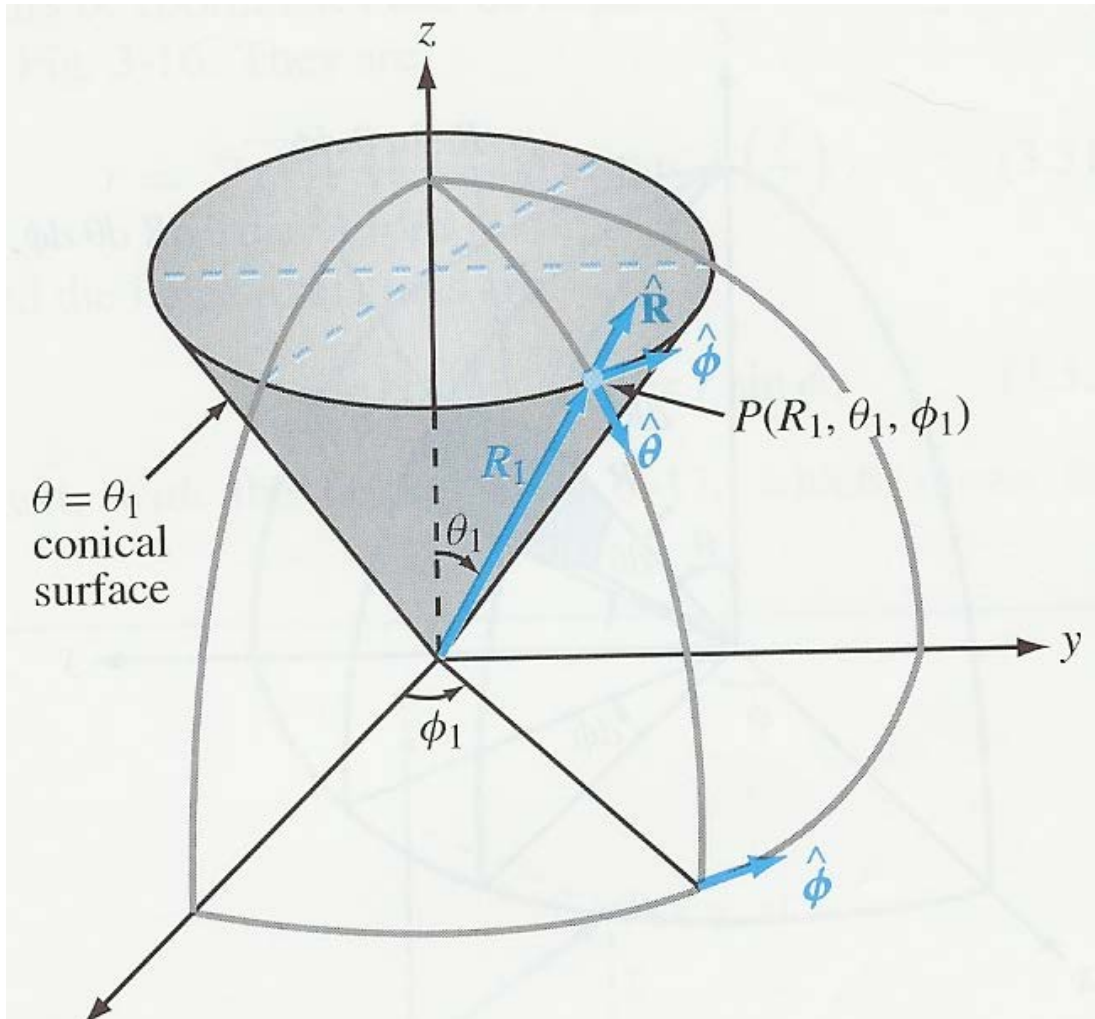


$$d\vec{s}_r = rd\phi dz \hat{r}$$

$$S = \int_{30}^{60} r d\phi \int_0^3 dz = 5 \cdot 3 \int_{\pi/6}^{\pi/3} d\phi$$

16.360 Lecture 14

• Spherical coordinate system



$$\hat{r} \times \hat{\theta} = \hat{\phi} \quad \hat{\theta} \times \hat{\phi} = \hat{R} \quad \hat{\phi} \times \hat{R} = \hat{\theta}$$

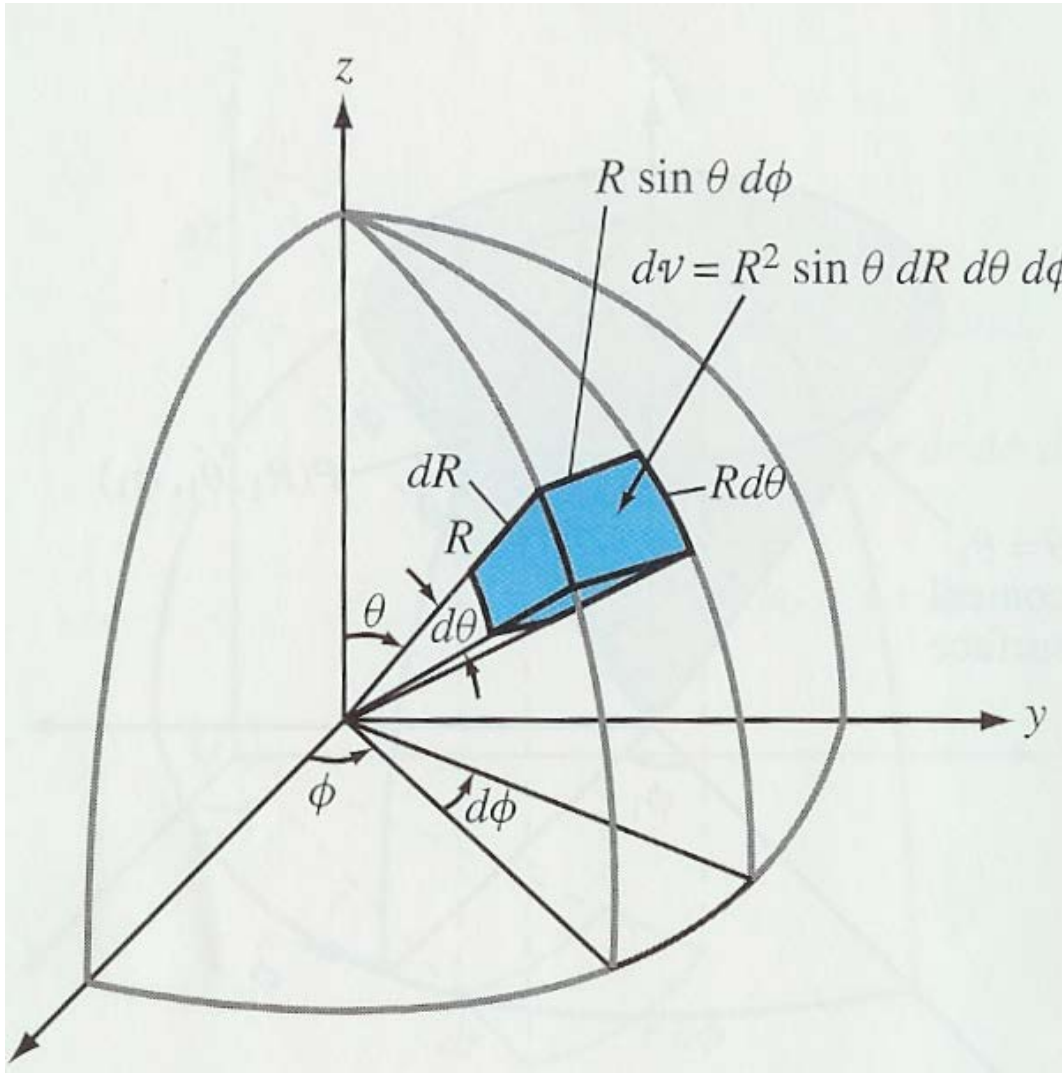
$$\hat{r} \times \hat{r} = 0 \quad \hat{r} \cdot \hat{r} = 1$$

$$\vec{A} = \hat{R} A_R + \hat{\theta} A_\theta + \hat{\phi} A_\phi$$

$$|\vec{A}| = \sqrt{A_R^2 + A_\theta^2 + A_\phi^2}$$

16.360 Lecture 14

- differential volume in Spherical coordinate system



$$\begin{aligned} d\vec{l} &= \hat{R} dl_R + \hat{\theta} dl_\theta + \hat{\phi} dl_\phi \\ &= \hat{R} dR + \hat{\theta} R d\theta + \hat{\phi} R \sin \theta d\phi \end{aligned}$$

$$d\vec{s}_R = \hat{R} dl_\theta dl_\phi = \hat{R} R^2 \sin \theta d\theta d\phi$$

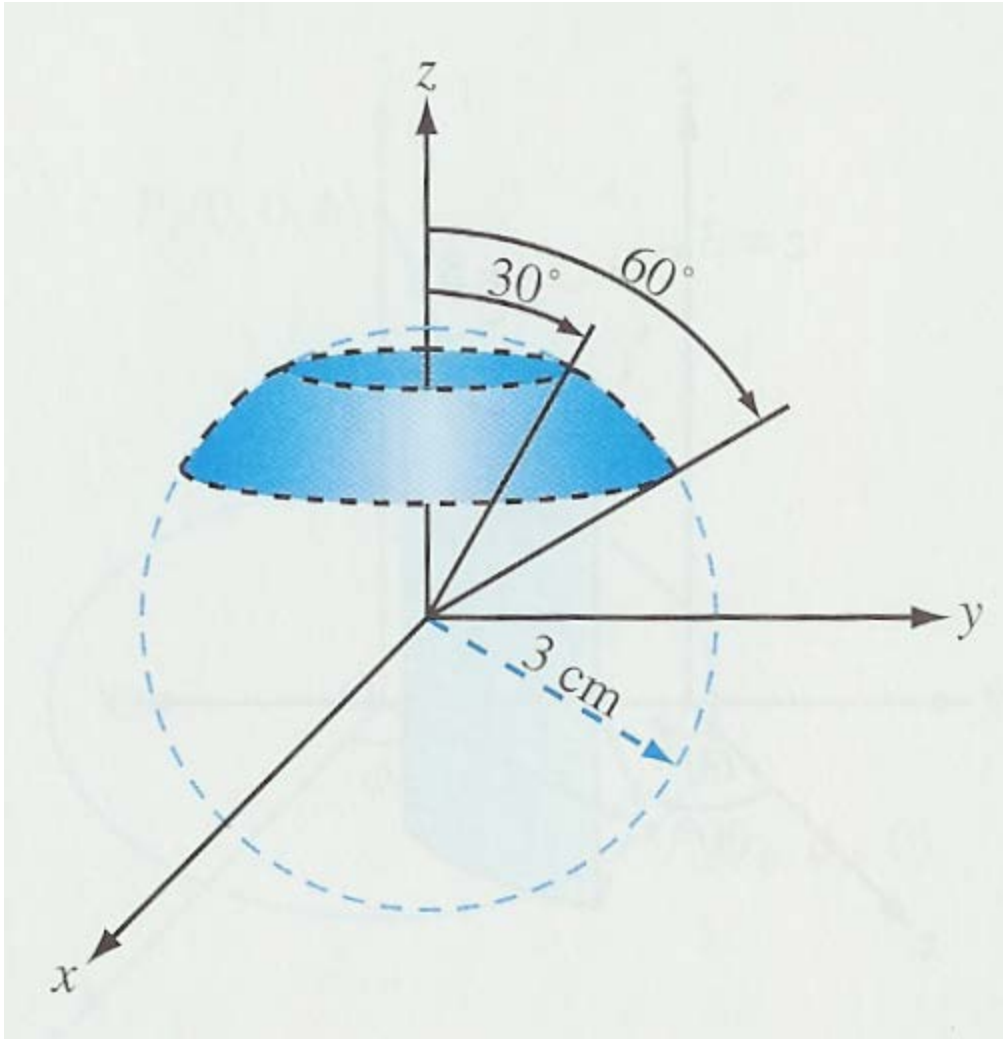
$$d\vec{s}_\theta = \hat{\theta} dl_R dl_\phi = \hat{\theta} R \sin \theta dR d\phi$$

$$d\vec{s}_\phi = \hat{\phi} dl_\theta dl_R = \hat{\phi} R d\theta dR$$

$$dv = dR dl_\theta dl_\phi = R^2 \sin \theta dR d\theta d\phi$$

16.360 Lecture 14

• Examples



(1) Find the area of the strip

$$S = \int d\vec{S}_R = \int_{30}^{60} \int_0^{2\pi} R^2 \sin \theta d\theta d\phi$$

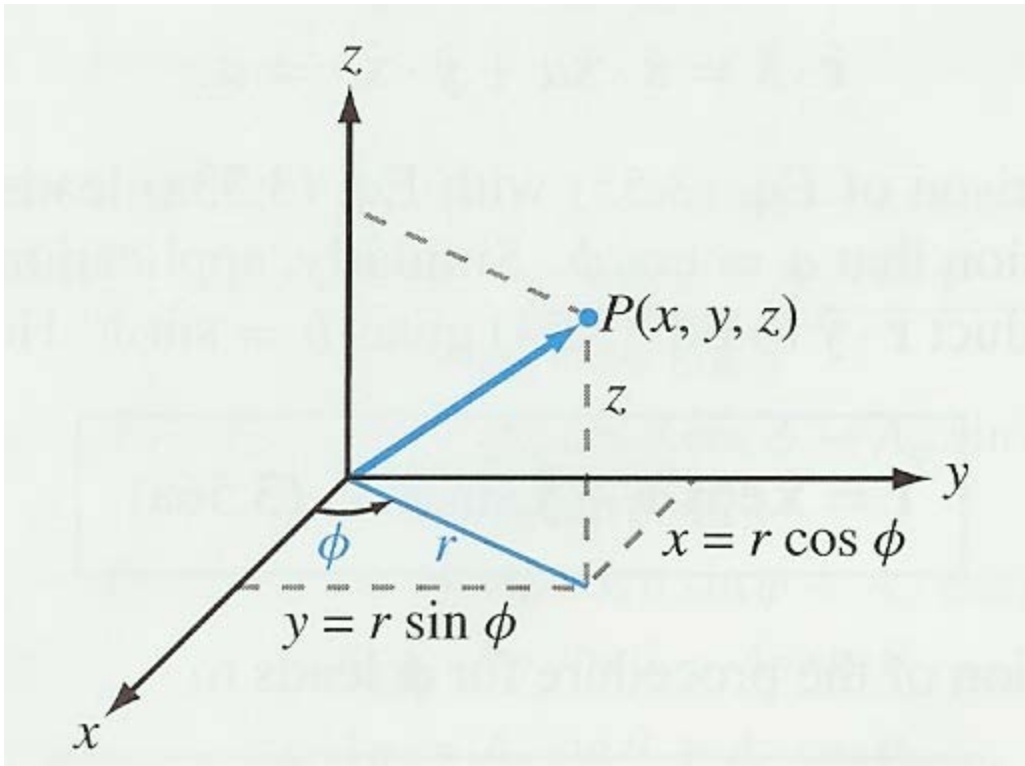
(2) A sphere of radius 2cm contains a volume charge density $\rho_v = 4 \cos^2 \theta$

Find the total charge contained in the sphere

$$\begin{aligned} Q &= \iiint dv \rho_v = \int_0^{2/100} dR \int_0^{180} d\theta \int_0^{260} d\phi \rho_v (R^2 \sin \theta) \\ &= \int_0^{2/100} dR R^2 \int_0^{180} d\theta \int_0^{260} d\phi (4 \cos^2 \theta) (\sin \theta) \\ &= \int_0^{2/100} dR R^2 \int_0^{180} d\theta (4 \cos^2 \theta) (\sin \theta) \int_0^{260} d\phi \end{aligned}$$

16.360 Lecture 15

- Cartesian to cylindrical transformation



$$r = \sqrt{x^2 + y^2},$$

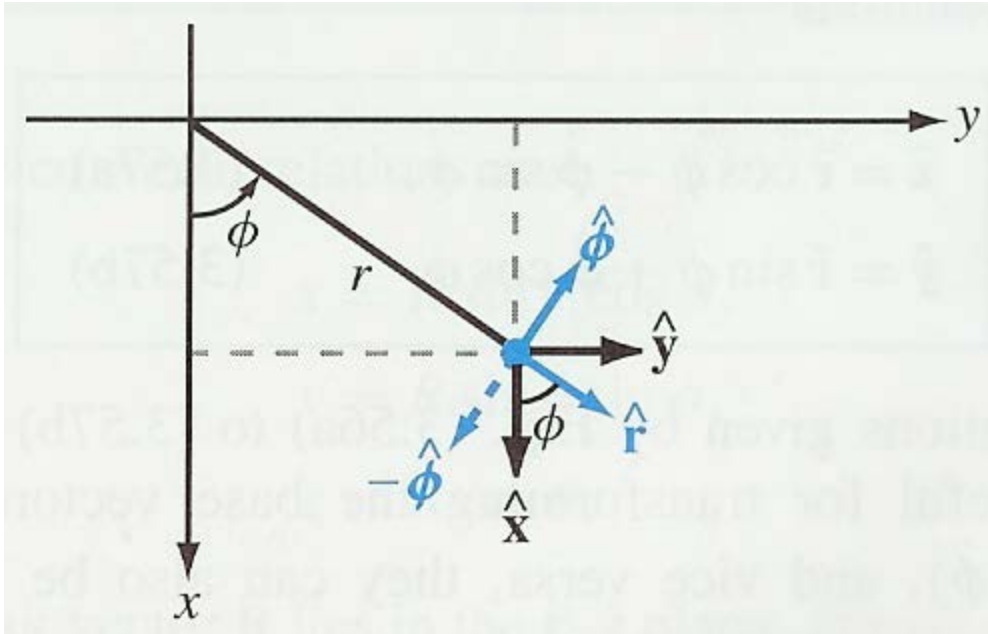
$$\phi = \tan^{-1}\left(\frac{y}{x}\right),$$

$$x = r \cos(\phi),$$

$$y = r \sin(\phi),$$

16.360 Lecture 15

• Cartesian to cylindrical transformation



$$\hat{r} \cdot \hat{x} = \cos(\phi), \quad \hat{r} \cdot \hat{y} = \sin(\phi),$$

$$\hat{\phi} \cdot \hat{x} = -\sin(\phi), \quad \hat{\phi} \cdot \hat{y} = \cos(\phi),$$

$$\hat{r} \cdot \hat{x} = (a\hat{x} + b\hat{y}) \cdot \hat{x} = a = \cos(\phi),$$

$$\hat{r} \cdot \hat{y} = (a\hat{x} + b\hat{y}) \cdot \hat{y} = b = \sin(\phi),$$

$$\hat{r} = \hat{x} \cos \phi + \hat{y} \sin \phi,$$

$$\hat{\phi} = -\hat{x} \sin \phi + \hat{y} \cos \phi,$$

$$\hat{x} = \hat{r} \cos \phi - \hat{\phi} \sin \phi,$$

$$\hat{y} = \hat{r} \sin \phi + \hat{\phi} \cos \phi,$$

16.360 Lecture 15

- Cartesian to cylindrical transformation

$$\vec{A} = A_x \hat{x} + A_y \hat{y} + A_z \hat{z},$$

$$\hat{r} \bullet \hat{x} = \cos(\phi), \quad \hat{r} \bullet \hat{y} = \sin(\phi),$$

$$\longrightarrow \vec{A} = A_r \hat{r} + A_\phi \hat{\phi} + A_z \hat{z},$$

$$\hat{\phi} \bullet \hat{x} = -\sin(\phi), \quad \hat{\phi} \bullet \hat{y} = \cos(\phi),$$

$$A_r = A_x \cos \phi + A_y \sin \phi,$$

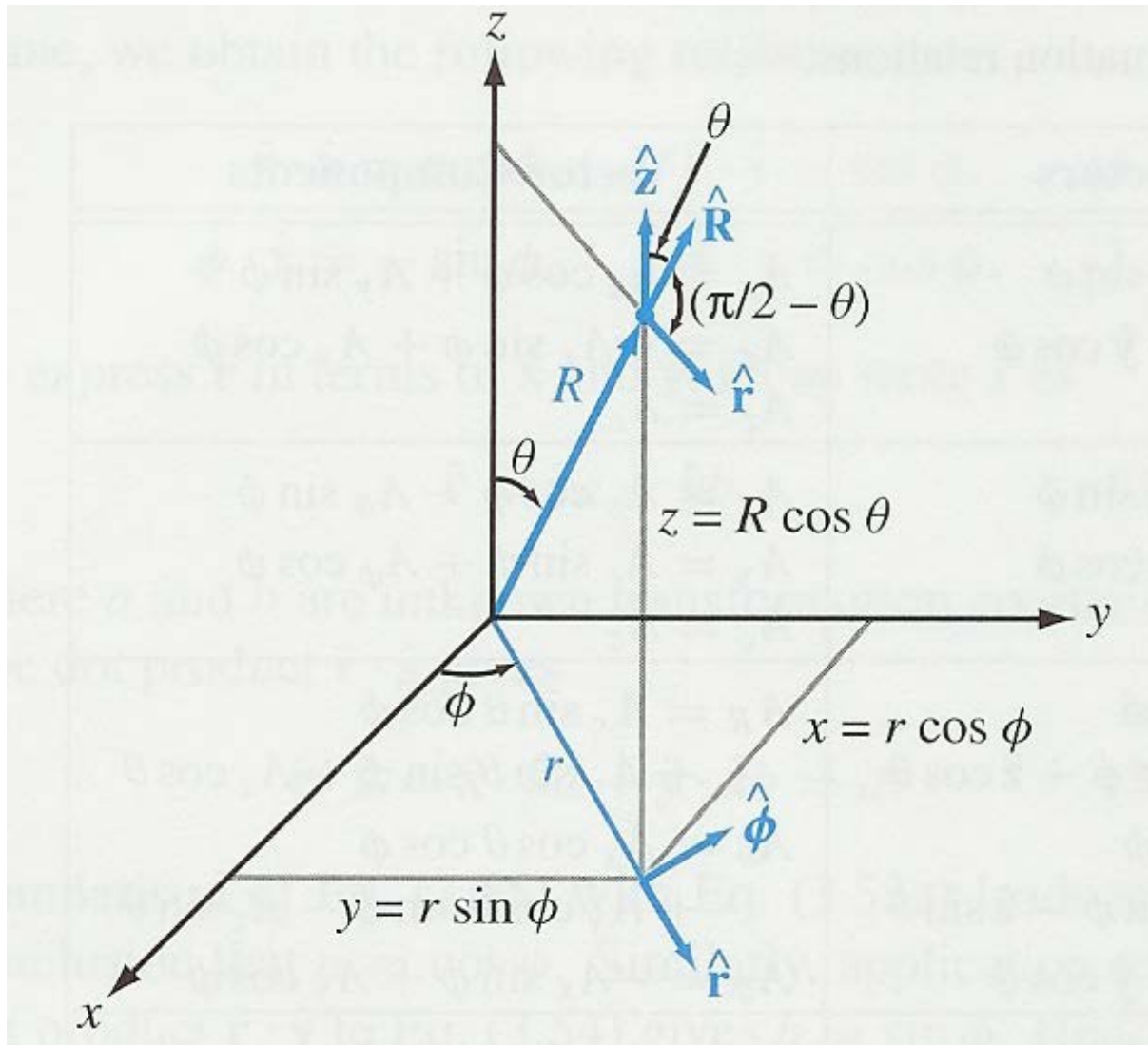
$$A_\phi = -A_x \sin \phi + A_y \cos \phi,$$

$$A_x = A_r \cos \phi - A_\phi \sin \phi,$$

$$A_y = A_r \sin \phi + A_\phi \cos \phi,$$

16.360 Lecture 15

• Cartesian to Spherical transformation



$$r = \sqrt{x^2 + y^2 + z^2},$$

$$\theta = \tan^{-1}\left(\frac{\sqrt{x^2 + y^2}}{z}\right),$$

$$\phi = \tan^{-1}\left(\frac{y}{x}\right),$$

$$x = R \sin \theta \cos \phi,$$

$$y = R \sin \theta \sin \phi,$$

$$z = R \cos \theta,$$

16.360 Lecture 15

- Cartesian to Spherical transformation

$$\hat{R} = \hat{r}a + \hat{z}b,$$

$$\hat{R} \bullet \hat{r} = a = \cos \theta, \quad \hat{r} = \hat{x} \cos \phi + \hat{y} \sin \phi,$$

$$\hat{R} \bullet \hat{z} = b = \sin \theta,$$

$$\hat{R} = \hat{x} \sin \theta \cos \phi + \hat{y} \sin \theta \sin \phi + \hat{z} \cos \theta,$$

$$\hat{\theta} = \hat{r}c + \hat{z}d, \quad \hat{\theta} \bullet \hat{r} = c = \cos \theta, \quad \hat{\theta} \bullet \hat{z} = d = -\sin \theta,$$

$$\hat{\theta} = \hat{x} \cos \theta \cos \phi + \hat{y} \cos \theta \sin \phi - \hat{z} \sin \theta,$$

$$\hat{\phi} = -\hat{x} \sin \phi + \hat{y} \cos \phi,$$

16.360 Lecture 15

- Cartesian to Spherical transformation

$$\hat{x} = \hat{R}c + \hat{\theta}d + \hat{\phi}e,$$

$$\hat{R} \bullet \hat{x} = \sin \theta \cos \phi, \quad \hat{\theta} \bullet \hat{x} = \cos \theta \cos \phi, \quad \hat{\phi} \bullet \hat{x} = e = -\sin \phi,$$

$$\hat{R} \bullet \hat{z} = b = \sin \theta, \quad \hat{r} = \hat{x} \cos \phi + \hat{y} \sin \phi,$$

$$\hat{x} = \hat{R} \sin \theta \cos \phi + \hat{\theta} \cos \theta \cos \phi - \hat{\phi} \sin \phi,$$

$$\hat{\theta} = \hat{r}c + \hat{z}d, \quad \hat{\theta} \bullet \hat{r} = c = \cos \theta, \quad \hat{\theta} \bullet \hat{r} = c = -\sin \theta,$$

$$\hat{y} = \hat{R} \sin \theta \sin \phi + \hat{\theta} \cos \theta \sin \phi + \hat{\phi} \cos \phi,$$

$$\hat{z} = \hat{R} \cos \theta - \hat{\theta} \sin \theta,$$

16.360 Lecture 15

- Distance between two points:

$$d_{12} = |\mathbf{R}_{12}| = \left[(x_2 - x_1)^2 + (y_2 - y_1)^2 + (z_2 - z_1)^2 \right]^{1/2},$$

$$d_{12} = |\mathbf{R}_{12}| = \left[(r_2 \cos \phi_2 - r_1 \cos \phi_1)^2 + (r_2 \sin \phi_2 - r_1 \sin \phi_1)^2 + (z_2 - z_1)^2 \right]^{1/2},$$

$$\begin{aligned} d_{12} &= |\mathbf{R}_{12}| = \left[(R_2 \sin \theta_2 \cos \phi_2 - R_1 \sin \theta_1 \cos \phi_1)^2 + (R_2 \sin \theta_2 \sin \phi_2 - R_1 \sin \theta_1 \sin \phi_1)^2 + (R_2 \cos \theta_2 - R_1 \cos \theta_1)^2 \right]^{1/2} \\ &= \left\{ R_2^2 + R_1^2 - 2R_1R_2[\cos \theta_2 \cos \theta_1 + \sin \theta_2 \sin \theta_1 \cos(\phi_2 - \phi_1)] \right\}^{1/2}, \end{aligned}$$

Gradient in Cartesian Coordinates

Gradient: differential change of a scalar

$$\begin{aligned}dT(x, y, z) &= \frac{\partial T}{\partial x} dx + \frac{\partial T}{\partial y} dy + \frac{\partial T}{\partial z} dz, \\&= \frac{\partial T}{\partial x} \hat{x} \bullet d\vec{l} + \frac{\partial T}{\partial y} \hat{y} \bullet d\vec{l} + \frac{\partial T}{\partial z} \hat{z} \bullet d\vec{l} \\&= \left(\frac{\partial T}{\partial x} \hat{x} + \frac{\partial T}{\partial y} \hat{y} + \frac{\partial T}{\partial z} \hat{z} \right) \bullet d\vec{l}, \\&= \nabla T \bullet d\vec{l},\end{aligned}$$

$$\nabla T \equiv \left(\frac{\partial T}{\partial x} \hat{x} + \frac{\partial T}{\partial y} \hat{y} + \frac{\partial T}{\partial z} \hat{z} \right),$$

$$\Delta T = \int dT(x, y, z) = \int \nabla T \bullet d\vec{l}$$

The direction of ∇T is along the maximum increase of T.

16.360 Lecture 16

Example of Gradient in Cartesian Coordinates

Find the directional derivative of $T = x^2 + y^2z$, along the direction $\hat{x}2 + \hat{y}3 - \hat{z}2$ and evaluate it at $(1, -1, 2)$.

$$\frac{dT}{dl} = \nabla T \cdot \hat{a}_l,$$

$$\nabla T = \left(\hat{x} \frac{\partial}{\partial x} + \hat{y} \frac{\partial}{\partial y} + \hat{z} \frac{\partial}{\partial z} \right) (x^2 + y^2z),$$

16.360 Lecture 16

Gradient operator in cylindrical Coordinates

$$\nabla T = \hat{x} \frac{\partial T}{\partial x} + \hat{y} \frac{\partial T}{\partial y} + \hat{z} \frac{\partial T}{\partial z},$$

$$\begin{aligned} \hat{x} &= \hat{r} \cos \phi - \hat{\phi} \sin \phi, \\ \hat{y} &= \hat{r} \sin \phi + \hat{\phi} \cos \phi, \end{aligned}$$

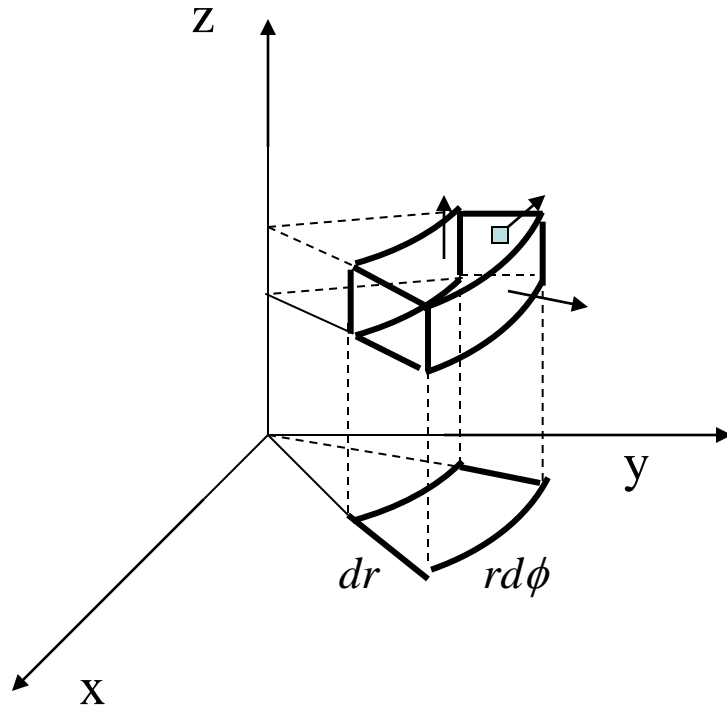
$$\begin{aligned} \frac{\partial T}{\partial x} &= \frac{\partial T}{\partial r} \frac{\partial r}{\partial x} + \frac{\partial T}{\partial \phi} \frac{\partial \phi}{\partial x} + \frac{\partial T}{\partial z} \frac{\partial z}{\partial x}, \\ &= \frac{\partial T}{\partial r} \cos \phi + \frac{\partial T}{\partial \phi} \left(-\frac{1}{r} \sin \phi\right), \end{aligned}$$

$$\begin{aligned} \frac{\partial T}{\partial y} &= \frac{\partial T}{\partial r} \frac{\partial r}{\partial y} + \frac{\partial T}{\partial \phi} \frac{\partial \phi}{\partial y} + \frac{\partial T}{\partial z} \frac{\partial z}{\partial y}, \\ &= \frac{\partial T}{\partial r} \sin \phi + \frac{\partial T}{\partial \phi} \left(\frac{1}{r} \cos \phi\right), \end{aligned}$$

$$\begin{aligned} \nabla T &= \hat{x} \frac{\partial T}{\partial x} + \hat{y} \frac{\partial T}{\partial y} + \hat{z} \frac{\partial T}{\partial z} \\ &= \hat{r} \left(\frac{\partial T}{\partial r} \cos^2 \phi - \frac{\partial T}{\partial \phi} \frac{1}{r} \cos \phi \sin \phi + \frac{\partial T}{\partial r} \sin^2 \phi + \frac{\partial T}{\partial \phi} \frac{1}{r} \cos \phi \sin \phi \right) \\ &\quad + \hat{\phi} \left(-\frac{\partial T}{\partial r} \sin \phi \cos \phi + \frac{\partial T}{\partial \phi} \frac{1}{r} \sin^2 \phi + \frac{\partial T}{\partial \phi} \frac{1}{r} \cos^2 \phi + \frac{\partial T}{\partial r} \sin \phi \cos \phi \right) \\ &= \hat{r} \frac{\partial T}{\partial r} + \hat{\phi} \frac{1}{r} \frac{\partial T}{\partial \phi} + \hat{z} \frac{\partial T}{\partial z} \end{aligned}$$

16.360 Lecture 16

Gradient operator in cylindrical Coordinates



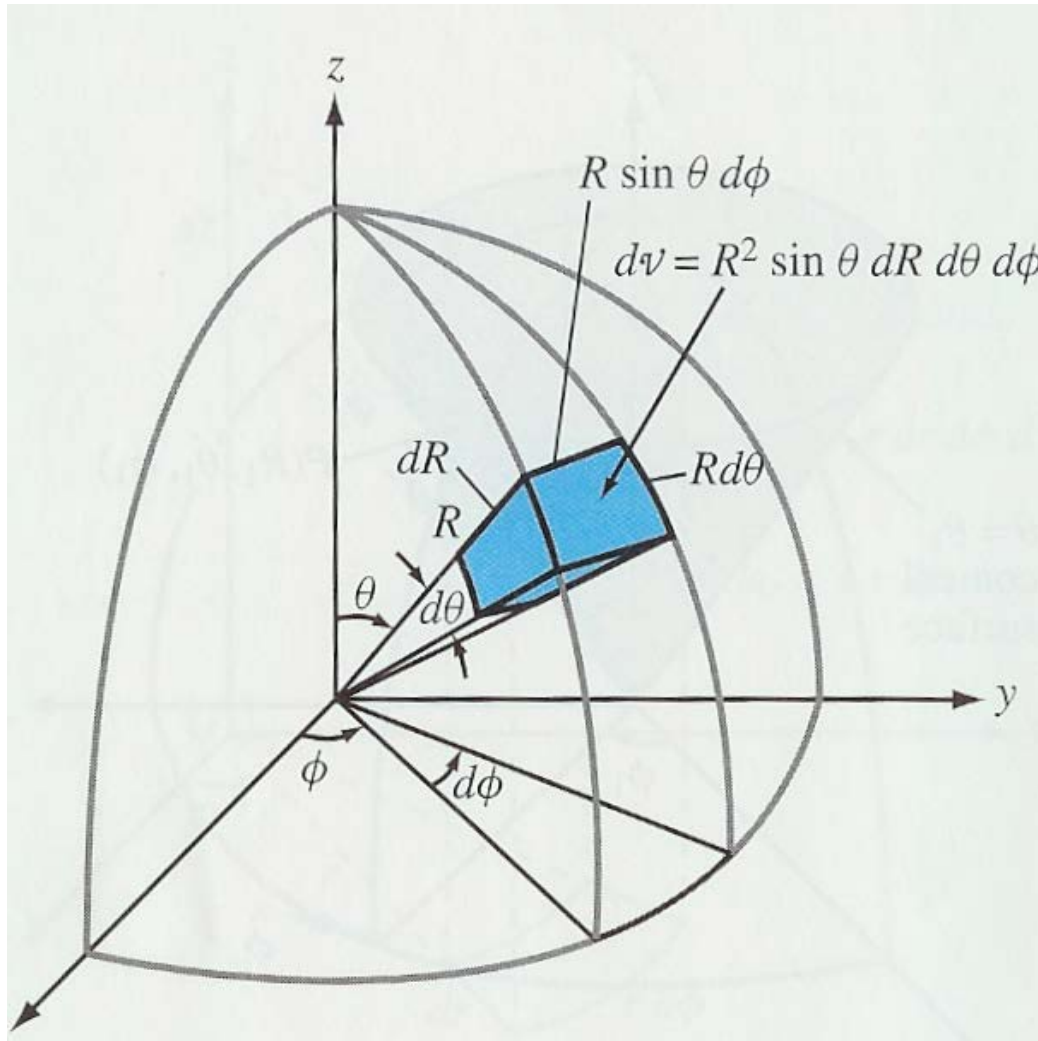
$$\nabla T = \hat{r} \frac{\partial T}{\partial l_r} + \hat{\phi} \frac{\partial T}{\partial l_\phi} + \hat{z} \frac{\partial T}{\partial l_z}$$

$$l_r = dr, \quad l_\phi = rd\phi, \quad l_z = dz,$$

$$\begin{aligned} \nabla T &= \hat{r} \frac{\partial T}{\partial l_r} + \hat{\phi} \frac{\partial T}{\partial l_\phi} + \hat{z} \frac{\partial T}{\partial l_z} \\ &= \hat{r} \frac{\partial T}{\partial r} + \hat{\phi} \frac{1}{r} \frac{\partial T}{\partial \phi} + \hat{z} \frac{\partial T}{\partial z} \end{aligned}$$

16.360 Lecture 16

Gradient operator in Spherical Coordinates



$$\nabla T = \hat{R} \frac{\partial T}{\partial l_R} + \hat{\theta} \frac{\partial T}{\partial l_\theta} + \hat{\phi} \frac{\partial T}{\partial l_\phi},$$

$$l_R = dR, \quad l_\theta = R d\theta, \quad l_\phi = R \sin \theta d\phi,$$

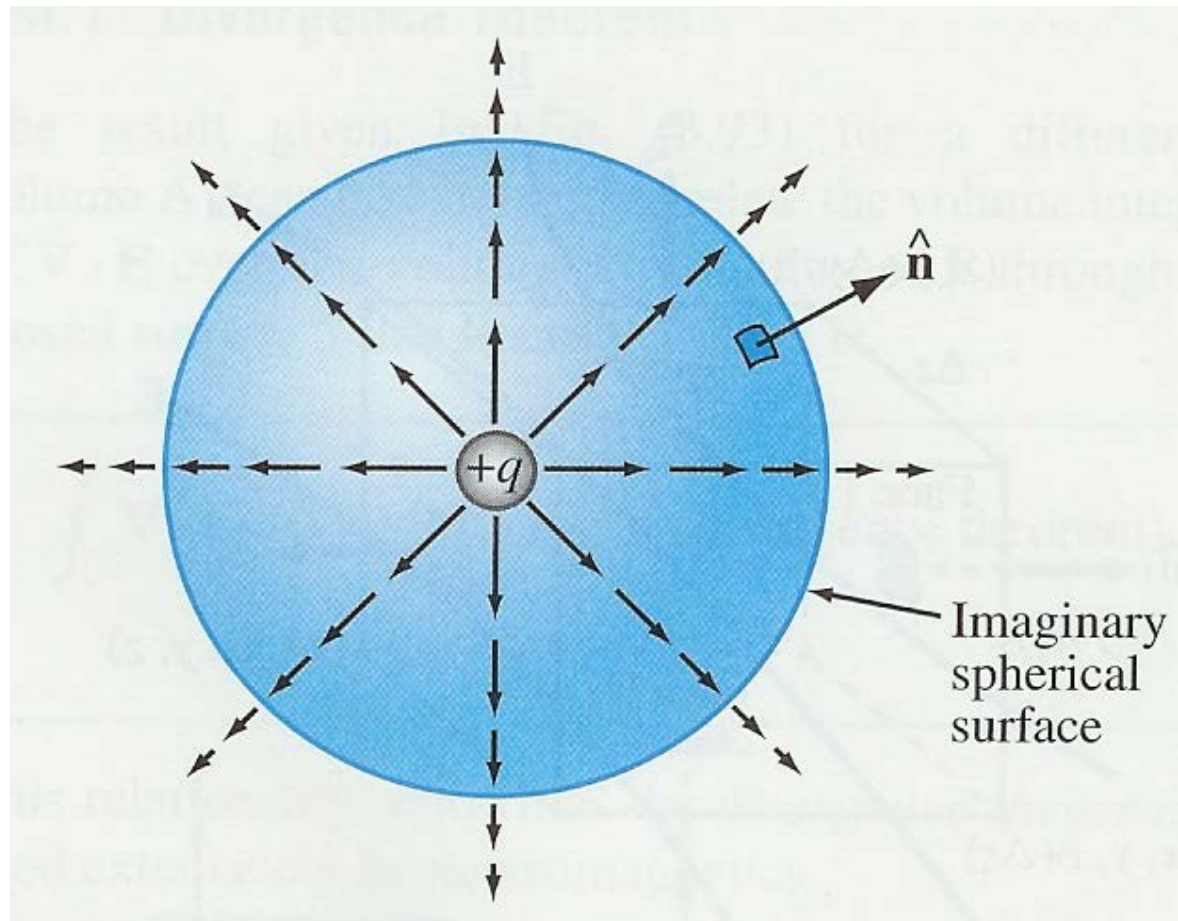
$$\begin{aligned} \nabla T &= \hat{R} \frac{\partial T}{\partial l_R} + \hat{\theta} \frac{\partial T}{\partial l_\theta} + \hat{\phi} \frac{\partial T}{\partial l_\phi} \\ &= \hat{R} \frac{\partial T}{\partial R} + \hat{\theta} \frac{1}{R} \frac{\partial T}{\partial \theta} + \hat{\phi} \frac{1}{R \sin \theta} \frac{\partial T}{\partial \phi} \end{aligned}$$

Properties of the Gradient operator

$$\nabla(U + V) = \nabla U + \nabla V,$$

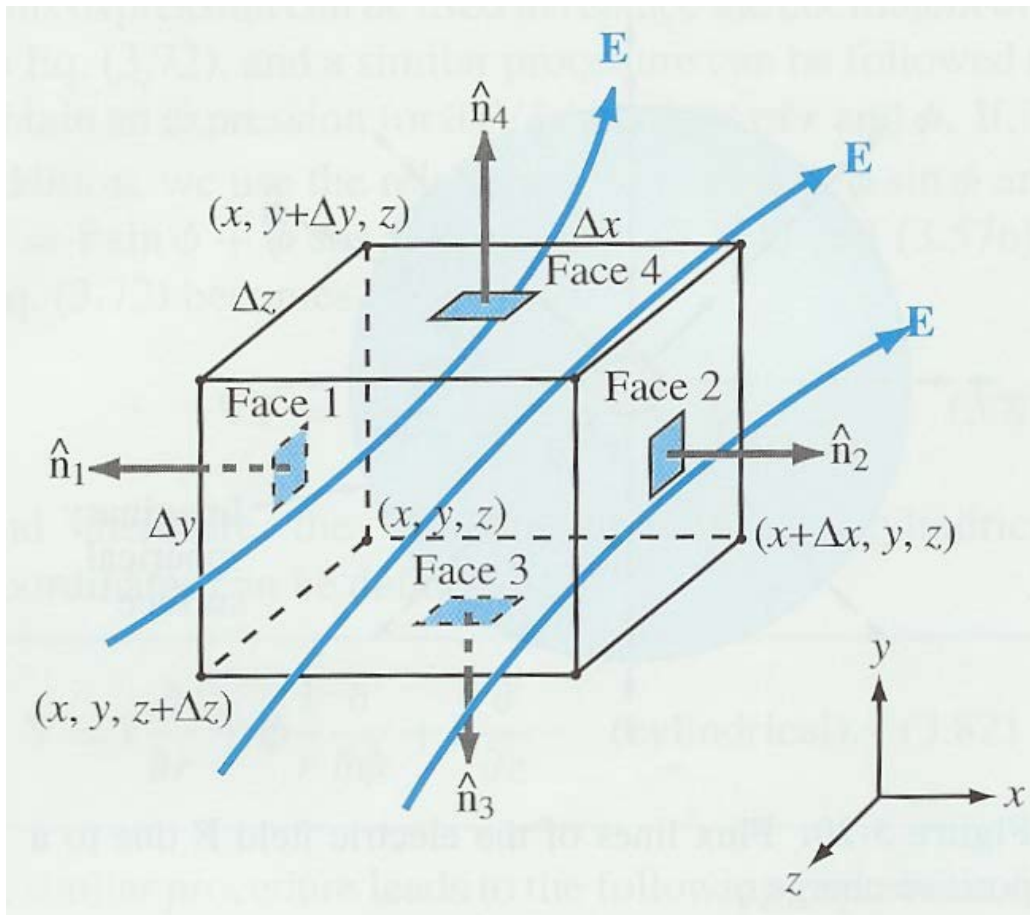
$$\nabla(UV) = U\nabla V + V\nabla U,$$

$$\nabla(V^n) = nV^{n-1}\nabla V,$$

Flux in Cartesian Coordinates

$$\text{Total flux} = \oint \vec{E} \cdot d\vec{s},$$

Flux in Cartesian Coordinates



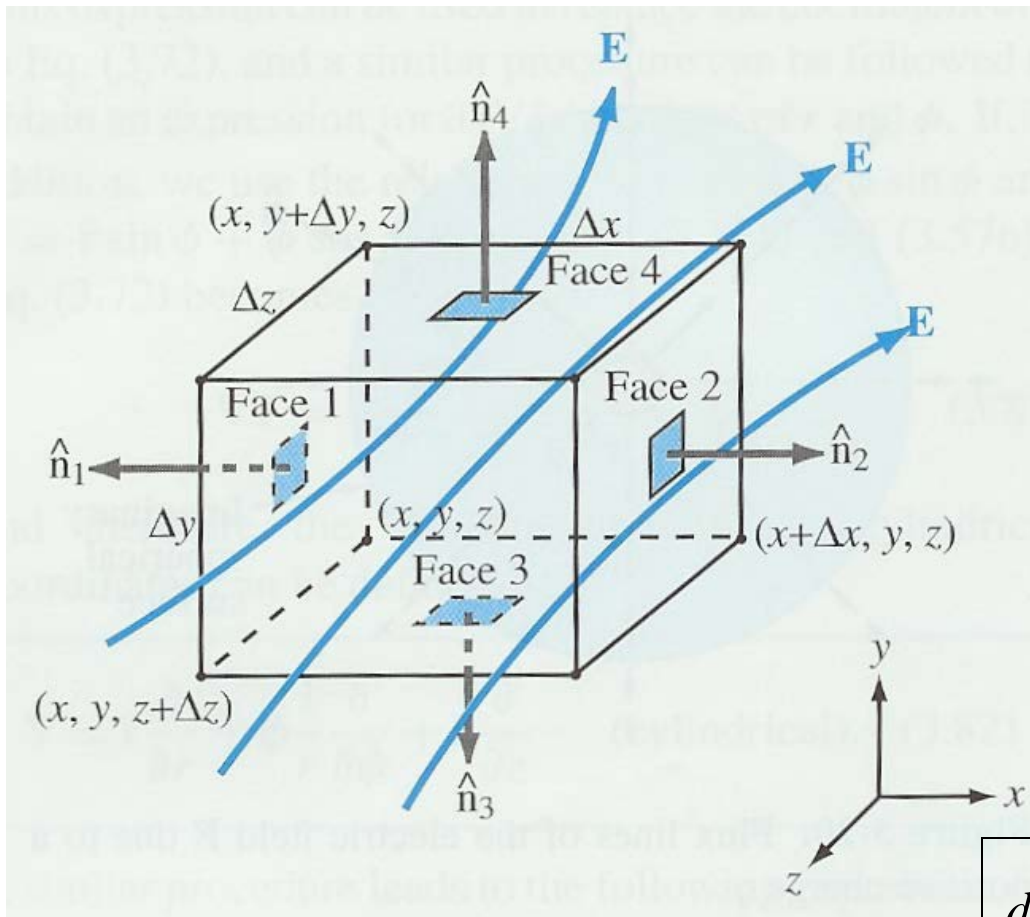
$$\begin{aligned}
 F_1 &= \int_{\text{face1}} \vec{E} \cdot \vec{n}_1 ds, \\
 &= \int_{\text{face1}} (\hat{x}E_x + \hat{y}E_y + \hat{z}E_z) \cdot (-\hat{x}) dydz \\
 &= -E_x(1)\Delta y\Delta z,
 \end{aligned}$$

$$\begin{aligned}
 F_2 &= \int_{\text{face2}} \vec{E} \cdot \vec{n}_2 ds, \\
 &= \int_{\text{face2}} (\hat{x}E_x + \hat{y}E_y + \hat{z}E_z) \cdot (\hat{x}) dydz \\
 &= E_x(2)\Delta y\Delta z,
 \end{aligned}$$

$$\begin{aligned}
 F_1 + F_2 &= [E_x(2) - E_x(1)]\Delta y\Delta z \\
 &= \frac{[E_x(2) - E_x(1)]}{\Delta x} \Delta x\Delta y\Delta z \\
 &= \frac{\partial E_x}{\partial x} \Delta x\Delta y\Delta z,
 \end{aligned}$$

16.360 Lecture 17

Definition of divergence in Cartesian Coordinates



$$F_1 + F_2 = [E_x(2) - E_x(1)]\Delta y\Delta z$$

$$= \frac{[E_x(2) - E_x(1)]}{\Delta x} \Delta x\Delta y\Delta z$$

$$= \frac{\partial E_x}{\partial x} \Delta x\Delta y\Delta z,$$

$$F_3 + F_4 = \frac{\partial E_y}{\partial y} \Delta x\Delta y\Delta z,$$

$$F_5 + F_6 = \frac{\partial E_z}{\partial z} \Delta x\Delta y\Delta z,$$

$$\oint_s \vec{E} \cdot d\vec{s} = \left(\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} \right) \Delta x\Delta y\Delta z,$$

$$\text{div} \vec{E} \equiv \nabla \cdot \vec{E} = \left(\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} \right),$$

Properties of divergence

$$\nabla \cdot (\vec{E}_1 + \vec{E}_2) = \nabla \cdot \vec{E}_1 + \nabla \cdot \vec{E}_2,$$

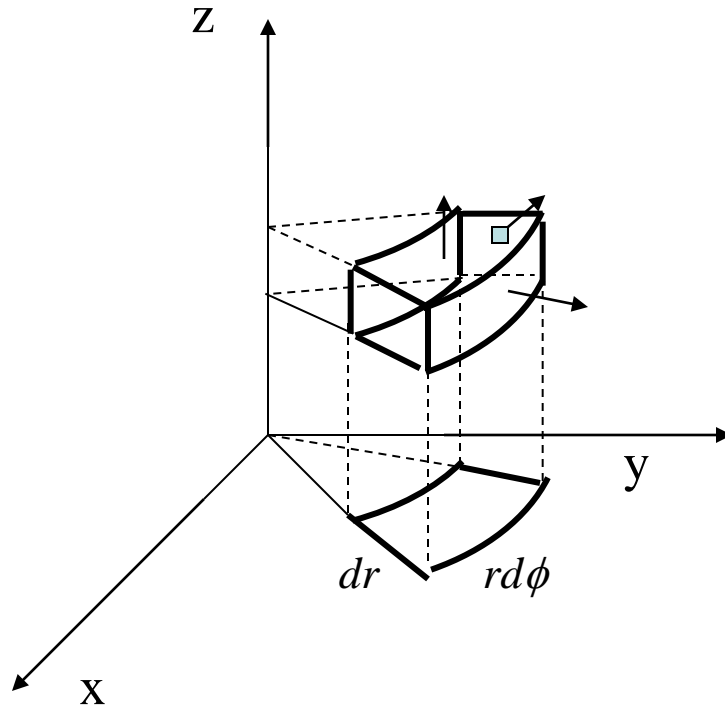
If $\nabla \cdot \vec{E}_1 = 0$, No net flux on any closed surface.

Divergence theorem

$$\begin{aligned} \oint_S \vec{E} \cdot d\vec{s} &= \left(\frac{\partial E_x}{\partial z} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} \right) \Delta x \Delta y \Delta z \\ &= \int_v \nabla \cdot \vec{E} dx dy dz, \end{aligned}$$

16.360 Lecture 17

Divergence in Cylindrical Coordinates



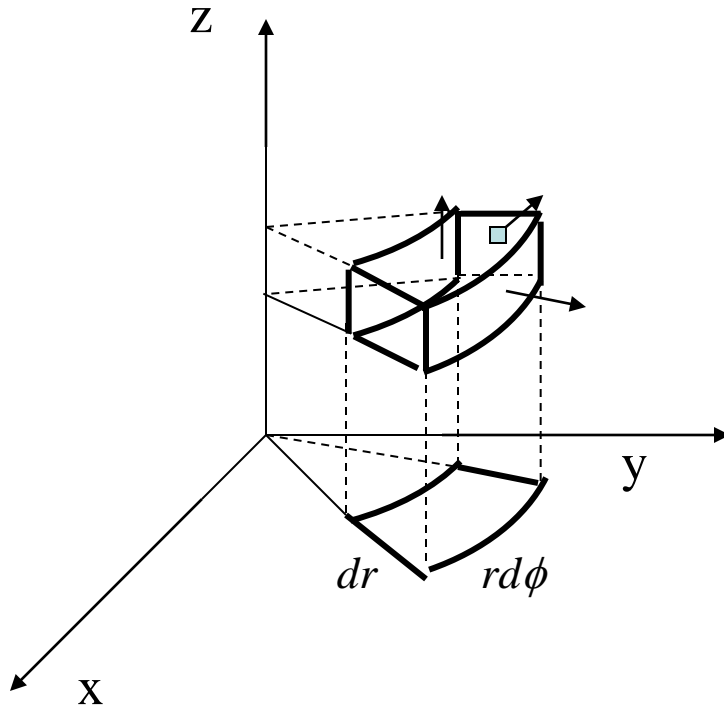
$$\begin{aligned}
 F_1 &= \int_{\text{face1}} \vec{E} \cdot \vec{n}_r ds_r, \\
 &= \int_{\text{face1}} (\hat{r}E_r + \hat{\phi}E_\phi + \hat{z}E_z) \cdot (-\hat{r})r\Delta\phi\Delta z \\
 &= -E_r(1)r\Delta\phi\Delta z,
 \end{aligned}$$

$$\begin{aligned}
 F_2 &= \int_{\text{face2}} \vec{E} \cdot \vec{n}_r ds_r, \\
 &= \int_{\text{face1}} (\hat{r}E_r + \hat{\phi}E_\phi + \hat{z}E_z) \cdot (\hat{r})r\Delta\phi\Delta z \\
 &= E_r(2)r\Delta\phi\Delta z,
 \end{aligned}$$

$$\begin{aligned}
 F_1 + F_2 &= [E_r(2) - E_r(1)]r\Delta\phi\Delta z \\
 &= \frac{[rE_r(2) - rE_r(1)]}{\Delta r} \Delta r\Delta\phi\Delta z \\
 &= \frac{\partial}{\partial r}(rE_r)\Delta r\Delta\phi\Delta z,
 \end{aligned}$$

16.360 Lecture 17

Divergence in Cylindrical Coordinates



$$\begin{aligned}
 F_5 + F_6 &= [E_z(2) - E_z(1)]r\Delta\phi\Delta r \\
 &= \frac{[E_z(2) - E_z(1)]}{\Delta z} r\Delta r\Delta\phi\Delta z \\
 &= \frac{\partial}{\partial z}(E_z)r\Delta r\Delta\phi\Delta z,
 \end{aligned}$$

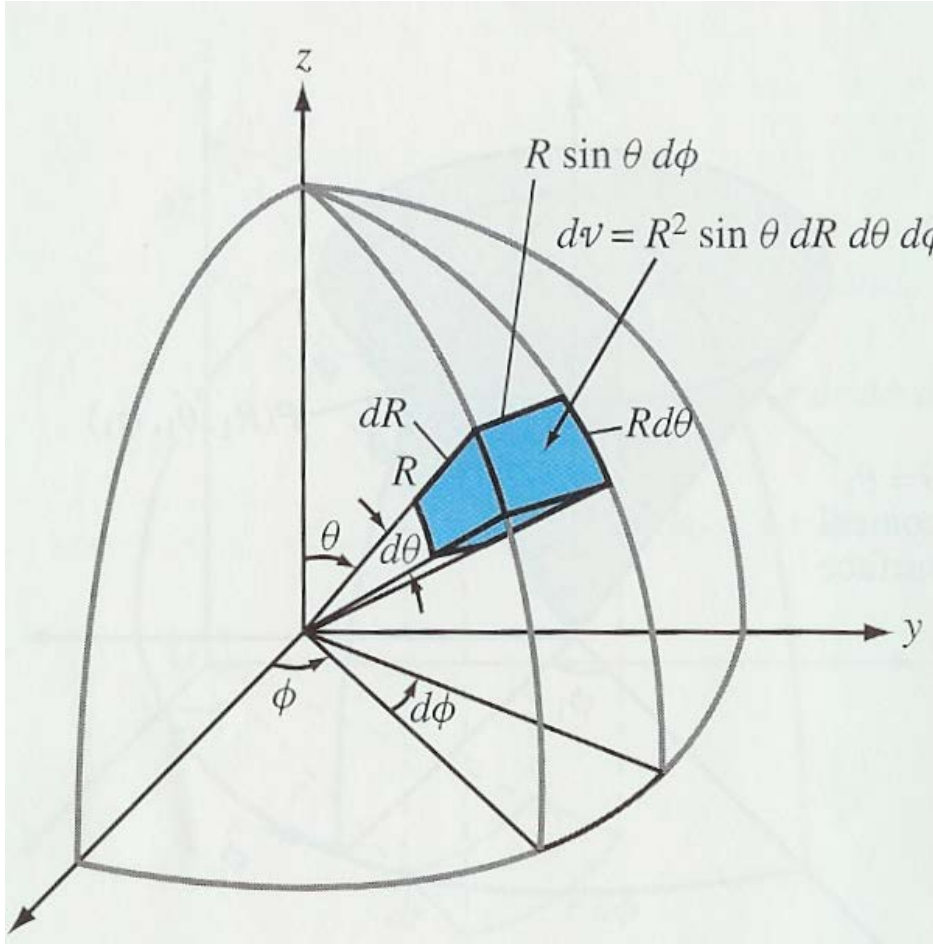
$$\begin{aligned}
 F_1 + F_2 &= [E_r(2) - E_r(1)]r\Delta\phi\Delta z \\
 &= \frac{[rE_r(2) - rE_r(1)]}{\Delta r} \Delta r\Delta\phi\Delta z \\
 &= \frac{\partial}{\partial r}(rE_r)\Delta r\Delta\phi\Delta z,
 \end{aligned}$$

$$\begin{aligned}
 F_3 + F_4 &= [E_\phi(2) - E_\phi(1)]\Delta r\Delta z \\
 &= \frac{[E_\phi(2) - E_\phi(1)]}{\Delta\phi} \Delta r\Delta\phi\Delta z \\
 &= \frac{\partial}{\partial\phi}(E_\phi)\Delta r\Delta\phi\Delta z,
 \end{aligned}$$

$$\begin{aligned}
 \oint_s \vec{E} \cdot d\vec{s} &= \left(\frac{\partial(rE_r)}{r\partial r} + \frac{\partial E_\phi}{r\partial\phi} + \frac{\partial E_z}{\partial z} \right) r\Delta r\Delta\phi\Delta z \\
 &= \nabla \cdot \vec{E} \Delta v,
 \end{aligned}$$

16.360 Lecture 17

Divergence in Spherical Coordinates



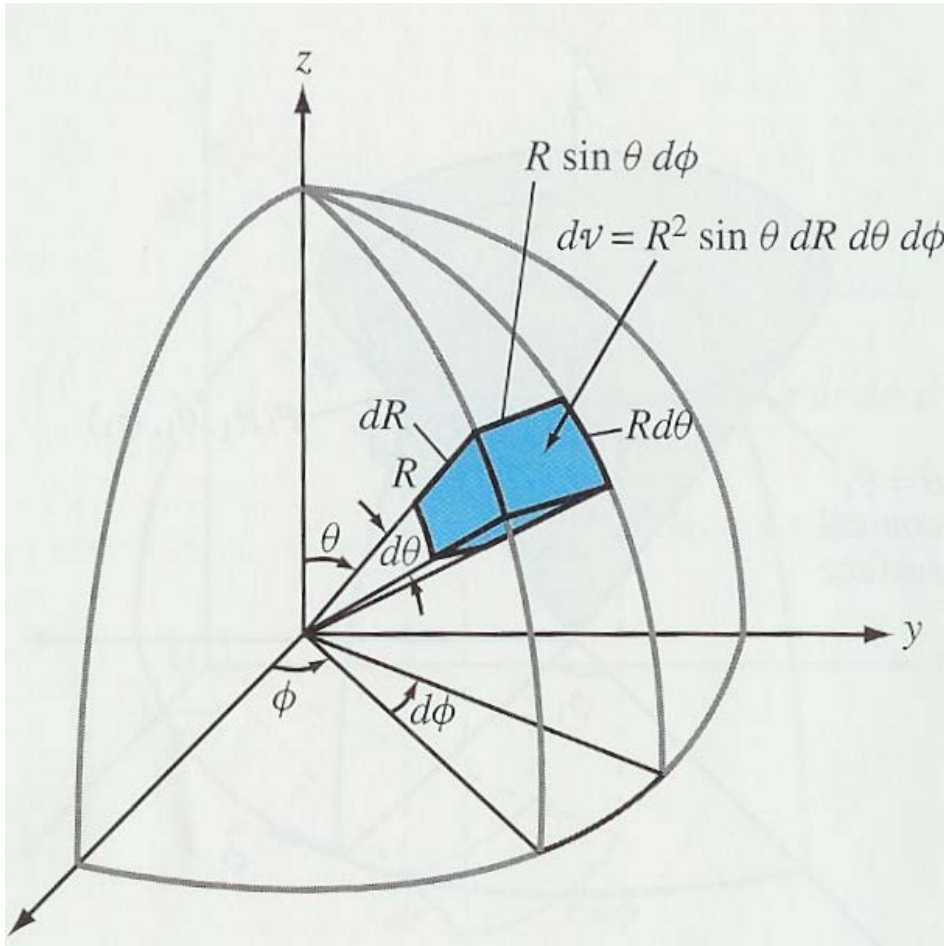
$$\begin{aligned}
 F_1 &= \int_{\text{face1}} \vec{E} \cdot \vec{n}_R ds_R, \\
 &= \int_{\text{face1}} (\hat{R}E_R + \hat{\theta}E_\theta + \hat{\phi}E_\phi) \cdot (-\hat{R})R \sin \theta \Delta \phi R \Delta \theta \\
 &= -E_R(1)R^2 \sin \theta \Delta \phi \Delta \theta,
 \end{aligned}$$

$$\begin{aligned}
 F_2 &= \int_{\text{face2}} \vec{E} \cdot \vec{n}_R ds_R, \\
 &= \int_{\text{face2}} (\hat{R}E_R + \hat{\theta}E_\theta + \hat{\phi}E_\phi) \cdot (\hat{R})R \sin \theta \Delta \phi R \Delta \theta \\
 &= E_R(2)R^2 \sin \theta \Delta \phi \Delta \theta,
 \end{aligned}$$

$$\begin{aligned}
 F_1 + F_2 &= [E_R(2) - E_R(1)]R^2 \sin \theta \Delta \theta \Delta \phi \\
 &= \frac{[R^2 E_R(2) - R^2 E_R(1)]}{\Delta R} \sin \theta \Delta R \Delta \phi \Delta \theta \\
 &= \frac{\partial}{\partial R} (R^2 E_R) \sin \theta \Delta R \Delta \phi \Delta \theta,
 \end{aligned}$$

16.360 Lecture 17

Divergence in Spherical Coordinates



$$\begin{aligned}
 F_1 + F_2 &= [E_R(2) - E_R(1)]R^2 \sin \theta \Delta \theta \Delta \phi \\
 &= \frac{[R^2 E_R(2) - R^2 E_R(1)]}{\Delta R} \sin \theta \Delta R \Delta \phi \Delta \theta \\
 &= \frac{\partial}{\partial R} (R^2 E_R) \sin \theta \Delta R \Delta \phi \Delta \theta,
 \end{aligned}$$

$$\begin{aligned}
 F_3 + F_4 &= [E_\theta(2) - E_\theta(1)]R \sin \theta \Delta R \Delta \phi \\
 &= \frac{[\sin \theta E_\theta(2) - \sin \theta E_\theta(1)]}{\Delta \theta} R \Delta R \Delta \phi \Delta \theta \\
 &= \frac{\partial}{\partial \theta} (\sin \theta E_R) R \Delta R \Delta \phi \Delta \theta,
 \end{aligned}$$

$$\begin{aligned}
 F_5 + F_6 &= [E_\phi(2) - E_\phi(1)]R \Delta \theta \Delta R \\
 &= \frac{[E_\phi(2) - E_\phi(1)]}{\Delta \phi} R \Delta R \Delta \phi \Delta \theta \\
 &= \frac{\partial}{\partial \phi} (E_\phi) \Delta R \Delta \phi \Delta \theta,
 \end{aligned}$$

16.360 Lecture 17

Divergence in Spherical Coordinates

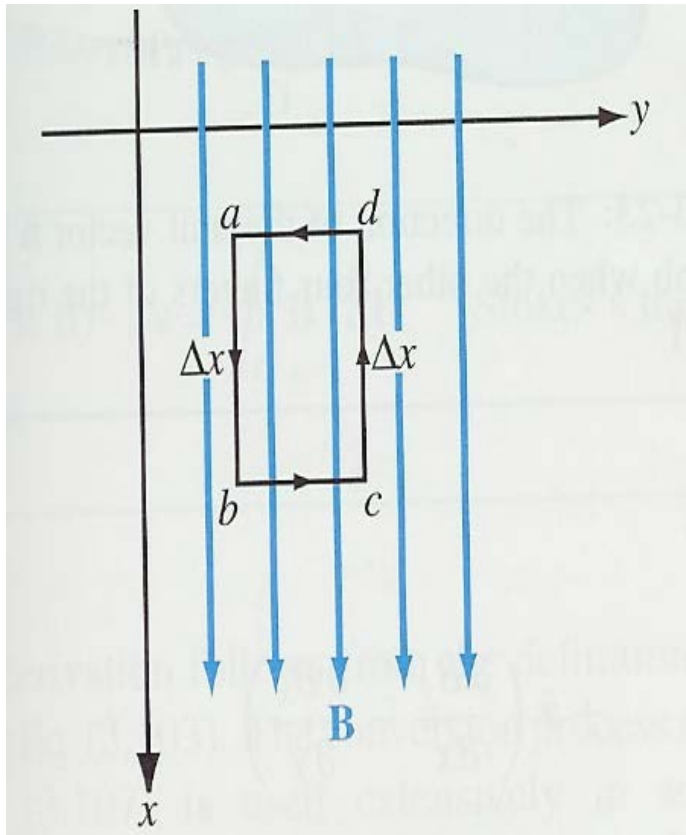
$$F_1 + F_2 = \frac{\partial}{\partial R} (R^2 E_R) \sin \theta \Delta R \Delta \phi \Delta \theta, \quad F_3 + F_4 = \frac{\partial}{\partial \theta} (\sin \theta E_R) R \Delta R \Delta \phi \Delta \theta,$$

$$F_5 + F_6 = \frac{\partial}{\partial \phi} (E_\phi) R \Delta R \Delta \phi \Delta \theta,$$

$$\begin{aligned} \oint_S \vec{E} \cdot d\vec{s} &= \left[\frac{\partial}{\partial R} (R^2 E_R) + R \frac{\partial}{\partial \theta} (\sin \theta E_\theta) + R \frac{\partial}{\partial \phi} (E_\phi) \right] \Delta R \Delta \phi \Delta \theta, \\ &= \left[\frac{1}{R^2} \frac{\partial}{\partial R} (R^2 E_R) + \frac{1}{R \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta E_\theta) + \frac{1}{R \sin \theta} \frac{\partial}{\partial \phi} (E_\phi) \right] R^2 \sin \theta \Delta R \Delta \phi \Delta \theta \\ &= \nabla \cdot \vec{E} \Delta v \end{aligned}$$

16.360 Lecture 18

Circulation of a Vector

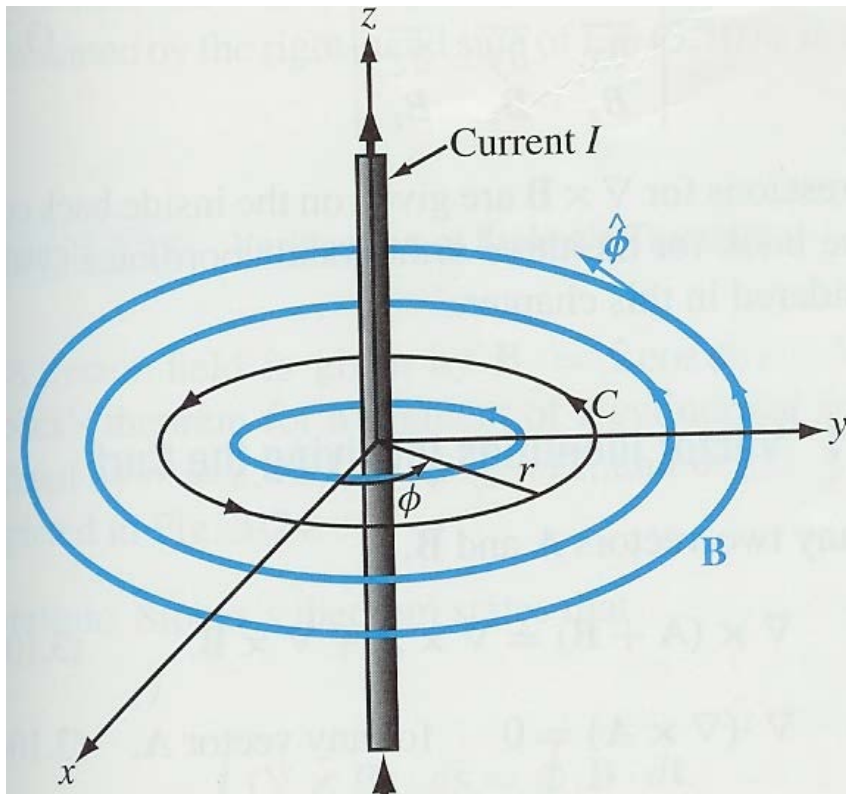


$$\text{Circulation} = \oint \vec{B} \cdot d\vec{l},$$

$$\text{Circulation} = \int_a^b \hat{x}B_0 \cdot \hat{x}dx + \int_b^c \hat{x}B_0 \cdot \hat{y}dy$$

$$+ \int_c^d \hat{x}B_0 \cdot \hat{x}dx + \int_d^a \hat{x}B_0 \cdot \hat{y}dy$$

$$= 0,$$

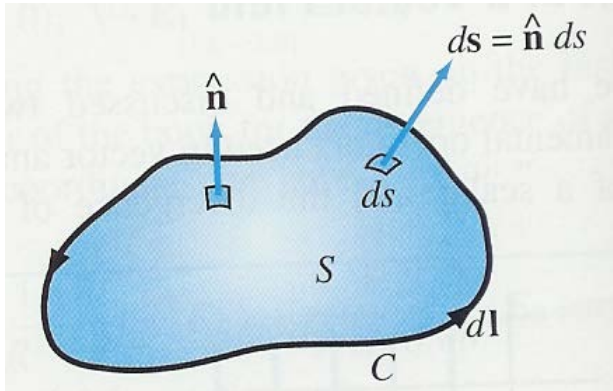
Circulation of a Vector

$$\vec{B} = \hat{\phi} \frac{\mu_0 I}{2\pi r},$$

$$\begin{aligned} \text{Circulation} &= \oint_C \vec{B} \cdot d\vec{l} = \int_0^{2\pi} \hat{\phi} \frac{\mu_0 I}{2\pi r} \cdot \hat{\phi} d\phi, \\ &= \mu_0 I, \end{aligned}$$

16.360 Lecture 18

Curl in Cartesian Coordinates



$$\int_{\text{face2}} \vec{n}_2 \vec{B} \cdot d\vec{l}_{y,z},$$

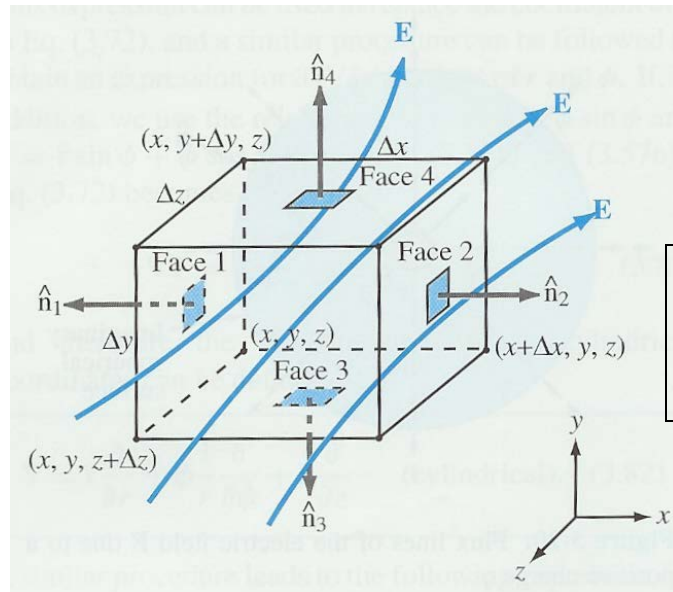
$$= \int_{\text{face2}} \hat{x} (B_{y+\Delta y} \Delta z - B_y \Delta z - B_{z+\Delta z} \Delta y + B_z \Delta y)$$

$$= \int \hat{x} \left(\frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z} \right) \Delta y \Delta z,$$

$$\oint_C \vec{n} \vec{B} \cdot d\vec{l} = \int \hat{x} \left(\frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z} \right) \Delta y \Delta z + \Delta x \Delta z + \int \hat{z} \left(\frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x} \right) \Delta y \Delta x$$

$$= \int \left[\hat{x} \left(\frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z} \right) + \hat{y} \left(\frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x} \right) + \hat{z} \left(\frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x} \right) \right] \cdot d\vec{s}$$

$$\int \hat{n} (\nabla \times \vec{B}) \Delta s$$



$$\nabla \times \vec{B} \equiv \hat{x} \left(\frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z} \right) + \hat{y} \left(\frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x} \right) + \hat{z} \left(\frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x} \right)$$

Vector identities involving the curl

$$\nabla \times (\vec{A} + \vec{B}) = \nabla \times \vec{A} + \nabla \times \vec{B}$$

$$\nabla \cdot (\nabla \times \vec{A}) = 0$$

$$\nabla \cdot (\nabla V) = 0$$

Stokes's theorem

$$\oint_c \vec{n} \vec{B} \cdot d\vec{l} = \int (\nabla \times B) \cdot d\vec{s}$$

16.360 Lecture 18

Curls in Rectangular, Cylindrical and Spherical Coordinates

$$\nabla \times \vec{B} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ B_x & B_y & B_z \end{vmatrix}$$

$$\nabla \times \vec{B} = \frac{1}{r} \begin{vmatrix} \hat{r} & r\hat{\phi} & \hat{z} \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \phi} & \frac{\partial}{\partial z} \\ A_r & rA_\phi & A_z \end{vmatrix}$$

$$\nabla \times \vec{B} = \frac{1}{R^2 \sin \theta} \begin{vmatrix} \hat{R} & R\hat{\theta} & R \sin \theta \hat{\phi} \\ \frac{\partial}{\partial R} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ A_R & RA_\theta & R \sin \theta A_\phi \end{vmatrix}$$

Laplacian Operator of a scalar

$$\nabla V = \hat{x} \frac{\partial V}{\partial x} + \hat{y} \frac{\partial V}{\partial y} + \hat{z} \frac{\partial V}{\partial z} = \hat{x} A_x + \hat{y} A_y + \hat{z} A_z$$

$$\nabla \bullet (\nabla V) = \frac{\partial}{\partial x} A_x + \frac{\partial}{\partial y} A_y + \frac{\partial}{\partial z} A_z$$

$$= \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}$$

$$\nabla^2 V \equiv \nabla \bullet (\nabla V) = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}$$

Laplacian Operator of a vector

$$\vec{E} = \hat{x}E_x + \hat{y}E_y + \hat{z}E_z$$

$$\nabla^2 \vec{E} = \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \vec{E} = \hat{x} \frac{\partial^2 E_x}{\partial x^2} + \hat{y} \frac{\partial^2 E_y}{\partial y^2} + \hat{z} \frac{\partial^2 E_z}{\partial z^2}$$

$$\nabla^2 \vec{E} = \nabla(\nabla \cdot \vec{E}) - \nabla \times (\nabla \times \vec{E})$$

16.360 Lecture 19

Maxwell equations

$$\begin{aligned}\nabla \cdot \vec{D} &= \rho_v, & \nabla \cdot \vec{B} &= 0, \\ \nabla \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t}, & \nabla \times \vec{H} &= \vec{J} + \frac{\partial \vec{D}}{\partial t},\end{aligned}$$

E: electric field intensity

D: electric flux intensity

H: magnetic field intensity

B: magnetic flux intensity

$\vec{D} = \epsilon \vec{E}$, $\vec{B} = \mu \vec{H}$, ϵ : electrical permittivity; μ : magnetic permeability

ρ_v : electric charge density per unit volume; \vec{J} : current density per unit area.

Electrostatics

$$\begin{aligned}\nabla \cdot \vec{D} &= \rho_v, \\ \nabla \times \vec{E} &= 0,\end{aligned}$$

Magnetostatics

$$\begin{aligned}\nabla \cdot \vec{B} &= 0, \\ \nabla \times \vec{H} &= \vec{J},\end{aligned}$$

16.360 Lecture 19

Electrostatics

$$\nabla \cdot \vec{D} = \rho_v, \quad \nabla \times \vec{E} = 0,$$

Volume charge density

$$\rho_v = \lim_{\Delta v \rightarrow 0} \frac{\Delta q}{\Delta v} = \frac{dq}{dv}, \quad Q = \int_v \rho_v dv,$$

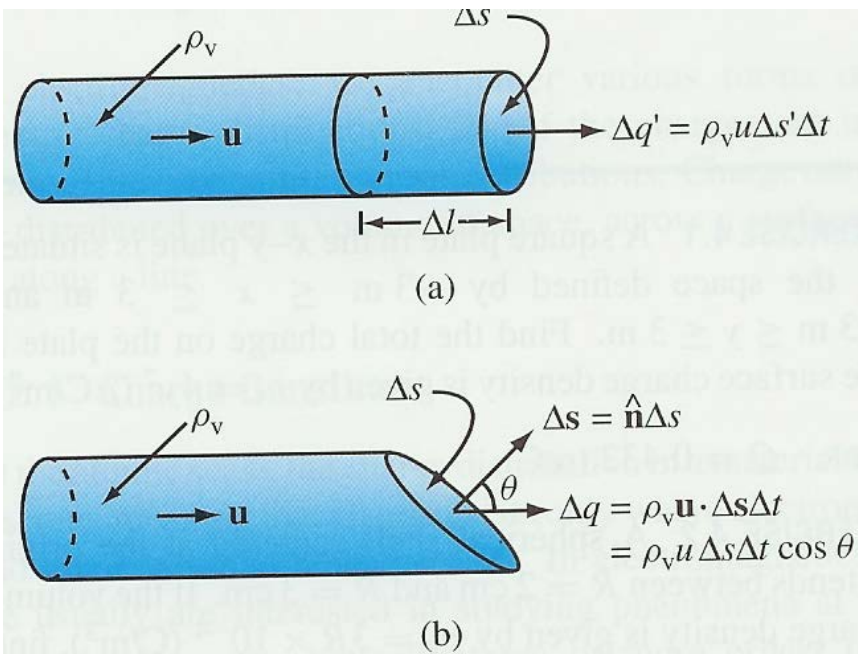
Surface charge density

$$\rho_s = \lim_{\Delta s \rightarrow 0} \frac{\Delta q}{\Delta s} = \frac{dq}{ds},$$

Line charge density

$$\rho_l = \lim_{\Delta l \rightarrow 0} \frac{\Delta q}{\Delta l} = \frac{dq}{dl},$$

Current density \mathbf{J}



$$\Delta q = \rho_v \vec{u} \cdot \Delta \vec{s} \Delta t,$$

$$\Delta I = \frac{\Delta q}{\Delta t} = \rho_v \vec{u} \cdot \Delta \vec{s},$$

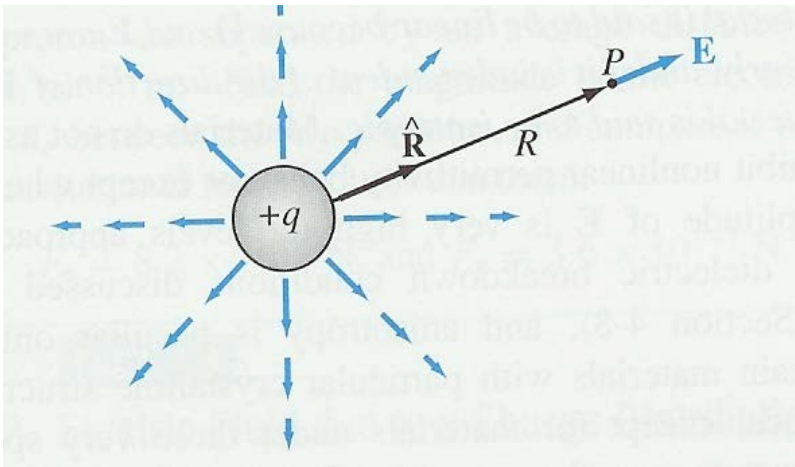
$$\vec{J} = \rho_v \vec{u},$$

$$I = \int_s \vec{J} \cdot d\vec{s},$$

Figure 4-2: Charges with velocity \mathbf{u} moving through a cross section $\Delta s'$ in (a) and Δs in (b).

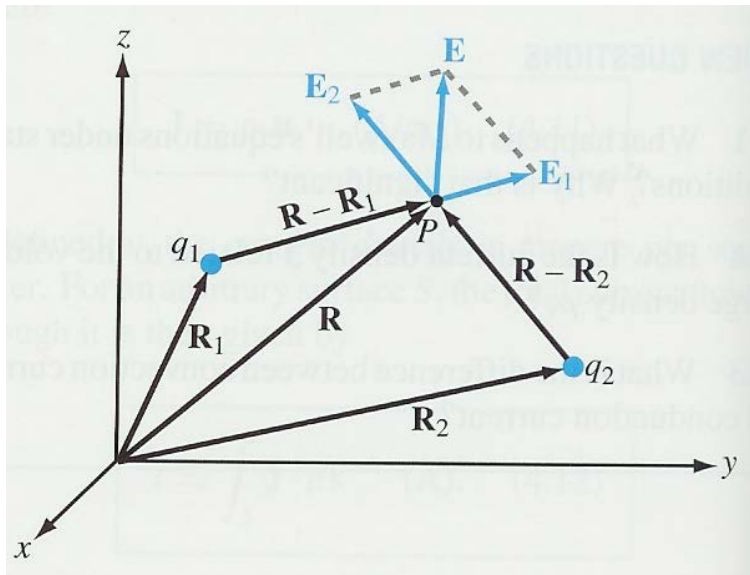
16.360 Lecture 19

Coulomb's law



$$\vec{E} = \vec{R} \frac{q}{4\pi\epsilon R^2},$$

$$\vec{F} = q' \vec{E},$$



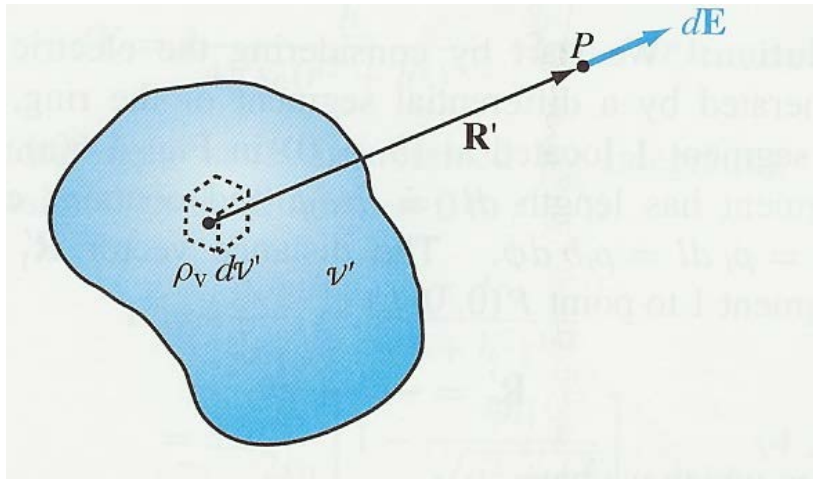
$$\vec{E}_1 = \frac{q_1(\vec{R} - \vec{R}_1)}{4\pi\epsilon |\vec{R} - \vec{R}_1|^3}, \quad \vec{E}_2 = \frac{q_2(\vec{R} - \vec{R}_2)}{4\pi\epsilon |\vec{R} - \vec{R}_2|^3},$$

$$\vec{E} = \vec{E}_1 + \vec{E}_2 = \frac{1}{4\pi\epsilon} \left[\frac{q_1(\vec{R} - \vec{R}_1)}{|\vec{R} - \vec{R}_1|^3} + \frac{q_2(\vec{R} - \vec{R}_2)}{|\vec{R} - \vec{R}_2|^3} \right],$$

$$\vec{E} = \frac{1}{4\pi\epsilon} \sum_{i=1}^N \frac{q_i(\vec{R} - \vec{R}_i)}{|\vec{R} - \vec{R}_i|^3},$$

16.360 Lecture 19

Electric field due to a charge distribution



$$d\vec{E} = \hat{R}' \frac{dq}{4\pi\epsilon R'^2},$$

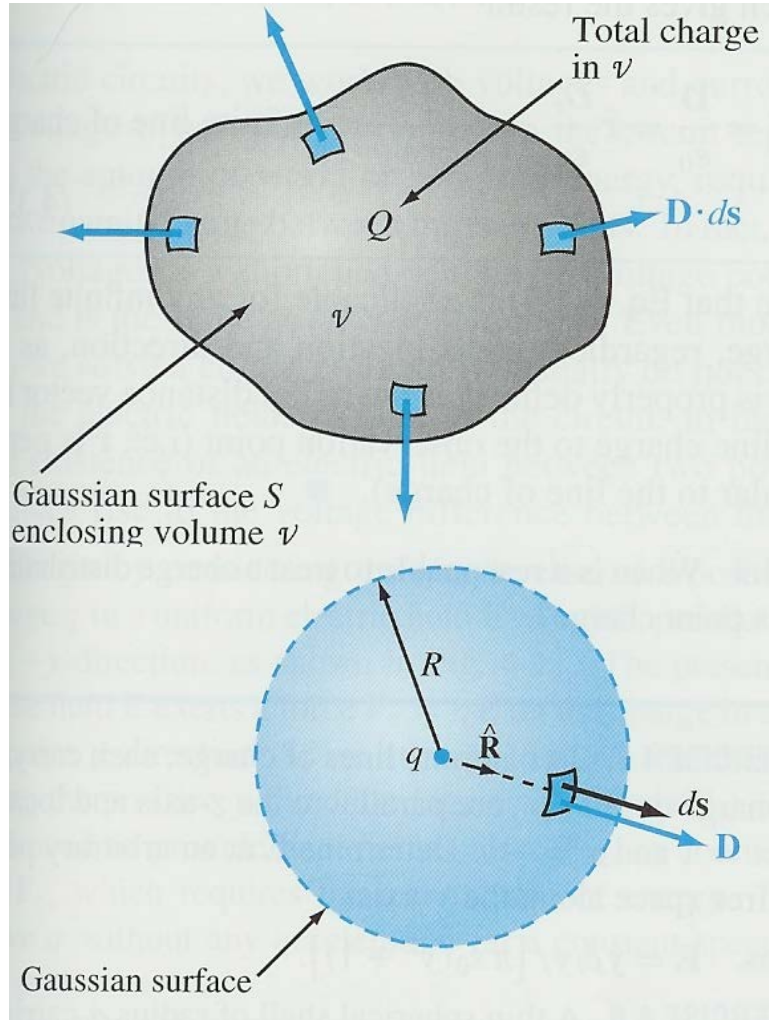
$$\vec{E} = \int d\vec{E} = \frac{1}{4\pi\epsilon} \int_{V'} \hat{R}' \frac{\rho_v dv'}{R'^2},$$

$$\vec{E} = \int d\vec{E} = \frac{1}{4\pi\epsilon} \int_s \hat{R}' \frac{\rho_s ds'}{R'^2},$$

$$\vec{E} = \int d\vec{E} = \frac{1}{4\pi\epsilon} \int_l \hat{R}' \frac{\rho_l dl'}{R'^2},$$

16.360 Lecture 19

Gauss's law



$$\nabla \cdot \vec{D} = \rho_v, \quad \int_v \nabla \cdot \vec{D} dv = \int_v \rho_v dv = Q,$$

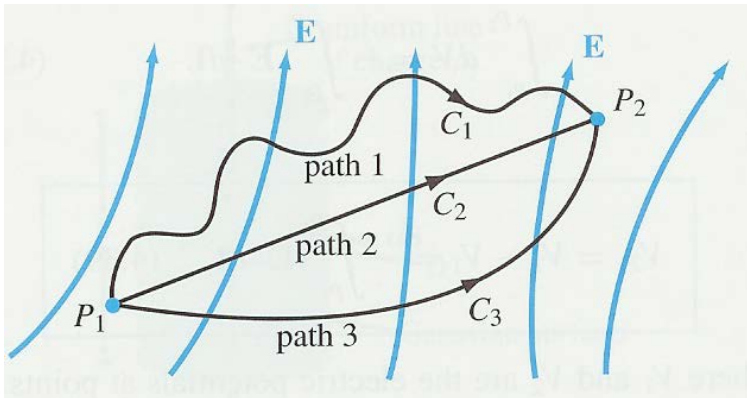
$$\int_v \nabla \cdot \vec{D} dv = \oint_s \vec{D} \cdot d\vec{s},$$

$$\boxed{\oint_s \vec{D} \cdot d\vec{s} = Q,}$$

Gauss's law

16.360 Lecture 19

Electrical scalar potential



$$dW = -\vec{F}_e \cdot d\vec{l} = -q\vec{E} \cdot d\vec{l},$$

$$dW = qdV,$$

$$dV = -\vec{E} \cdot d\vec{l},$$

$$V_{21} = V_2 - V_1 = \int_{P_1}^{P_2} dV = \int_{P_1}^{P_2} -\vec{E} \cdot d\vec{l},$$

$$0 = \oint_C dV = -\oint_C \vec{E} \cdot d\vec{l} = 0,$$

$$\oint_s (\nabla \times \vec{E}) \cdot d\vec{s} = \oint_C \vec{E} \cdot d\vec{l} = 0,$$

16.360 Lecture 19

Electrical potential due to point charge

$$\vec{E} = \vec{R} \frac{q}{4\pi\epsilon R^2}, \quad V = \int_{\infty}^R -\vec{E} \cdot d\vec{l}$$

$$V = \int_{\infty}^R -\vec{E} \cdot d\vec{l} = -\int_{\infty}^R \hat{R} \frac{q}{4\pi\epsilon R^2} \cdot \hat{R} dR = \frac{q}{4\pi\epsilon R},$$

$$V(\vec{R}) = \frac{q}{4\pi\epsilon |\vec{R} - \vec{R}_1|},$$

Electrical potential due to continuous distributions

$$V(\vec{R}) = \frac{1}{4\pi\epsilon} \int_{v'} \frac{\rho_{v'}}{|\vec{R} - \vec{R}'|} dv', \quad V(\vec{R}) = \frac{1}{4\pi\epsilon} \int_{s'} \frac{\rho_{s'}}{|\vec{R} - \vec{R}'|} ds',$$

$$V(\vec{R}) = \frac{1}{4\pi\epsilon} \int_{l'} \frac{\rho_{l'}}{|\vec{R} - \vec{R}'|} dl',$$

16.360 Lecture 19

Electric field as a function of Electrical potential

$$dV = -\vec{E} \cdot d\vec{l}, \quad dV = \nabla V \cdot d\vec{l},$$

$$\vec{E} = -\nabla V,$$

Poisson's equation

$$\nabla \cdot \vec{D} = \rho_v, \quad \nabla \cdot \vec{E} = \frac{\rho_v}{\epsilon},$$

$$\nabla \cdot \nabla V = -\frac{\rho_v}{\epsilon}, \quad \boxed{\nabla^2 V = -\frac{\rho_v}{\epsilon}}, \quad \text{Poisson's equation}$$

$$\boxed{\nabla^2 V = 0}, \quad \text{Laplace's equation}$$

16.360 Lecture 19

Electrical properties of material

- conductor
- dielectric
- semiconductor

Material	Conductivity, σ (S/m)
<i>Conductors</i>	
Silver	6.2×10^7
Copper	5.8×10^7
Gold	4.1×10^7
Aluminum	3.5×10^7
Iron	10^7
Mercury	10^6
Carbon	3×10^4
<i>Semiconductors</i>	
Pure germanium	2.2
Pure silicon	4.4×10^{-4}
<i>Insulators</i>	
Glass	10^{-12}
Paraffin	10^{-15}
Mica	10^{-15}
Fused quartz	10^{-17}

Conductors

Electron drift velocity $u_e = -\mu_e \vec{E}$

Hole drift velocity $u_h = \mu_h \vec{E}$

Conducting current

$$\vec{J} = \vec{J}_e + \vec{J}_h = \rho_{ve} u_e + \rho_{vh} u_h = (-\rho_{ve} \mu_e + \rho_{vh} \mu_h) \vec{E},$$

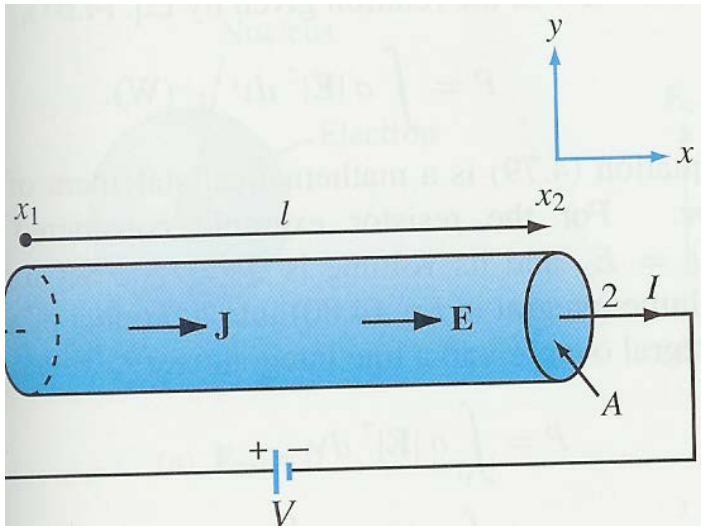
$$\sigma = -\rho_{ve} \mu_e + \rho_{vh} \mu_h,$$

Point form of Ohm's law

$$\boxed{\vec{J} = \sigma \vec{E},}$$

16.360 Lecture 20

Resistance



$$V = V_1 - V_2 = -\int_{x_2}^{x_1} \vec{E} \cdot d\vec{l} = E_x l,$$

$$I = \int_A \vec{J} \cdot d\vec{s} = \int_A \sigma \vec{E} \cdot d\vec{s} = \sigma E_x A,$$

$$R = \frac{V}{I} = \frac{l}{\sigma A},$$

General form

$$R = \frac{V}{I} = \frac{-\int_{x_2}^{x_1} \vec{E} \cdot d\vec{l}}{\int_A \vec{J} \cdot d\vec{s}} = \frac{-\int_{x_2}^{x_1} \vec{E} \cdot d\vec{l}}{\int_A \sigma \vec{E} \cdot d\vec{s}},$$

Joule's law

$$\Delta W = F_e \Delta l_e + F_h \Delta l_h,$$

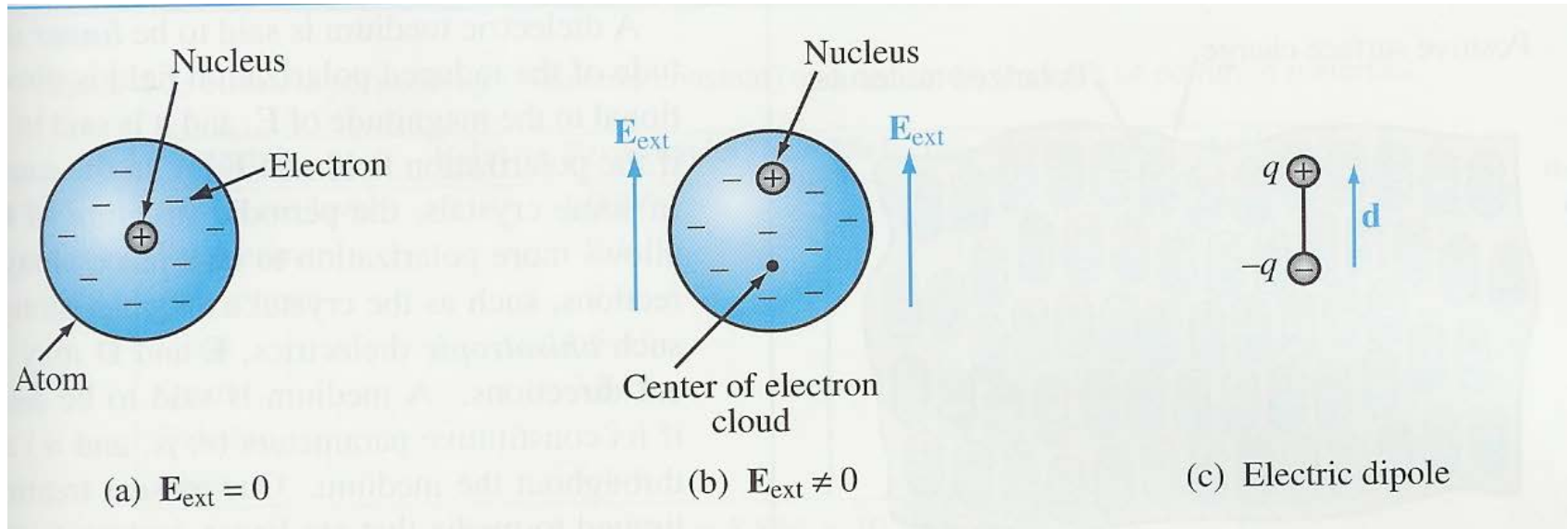
$$\begin{aligned}\Delta P &= \frac{\Delta W}{\Delta t} = F_e \frac{\Delta l_e}{\Delta t} + F_h \frac{\Delta l_h}{\Delta t} \\ &= F_e u_e + F_h u_h = (\rho_{ve} u_e \vec{E} + \rho_{vh} u_h \vec{E}) \Delta v \\ &= \vec{J} \cdot \vec{E} \Delta v,\end{aligned}$$

General form

$$P = \int_v \vec{J} \cdot \vec{E} dv,$$

16.360 Lecture 20

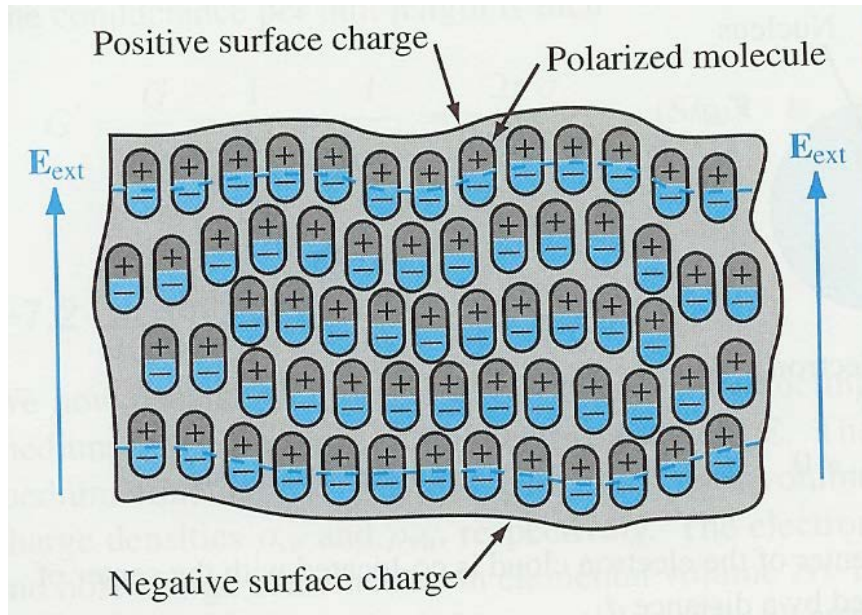
Dielectrics



Electrical field induced polarization

16.360 Lecture 20

Dielectrics



$$\vec{D} = \epsilon_0 \vec{E} + \vec{P},$$

\vec{P} : electric polarization field

For homogeneous material:

$$\vec{P} = \epsilon_0 \chi_e \vec{E},$$

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P} = \epsilon_0 \vec{E} + \epsilon_0 \chi_e \vec{E} = \epsilon \vec{E},$$

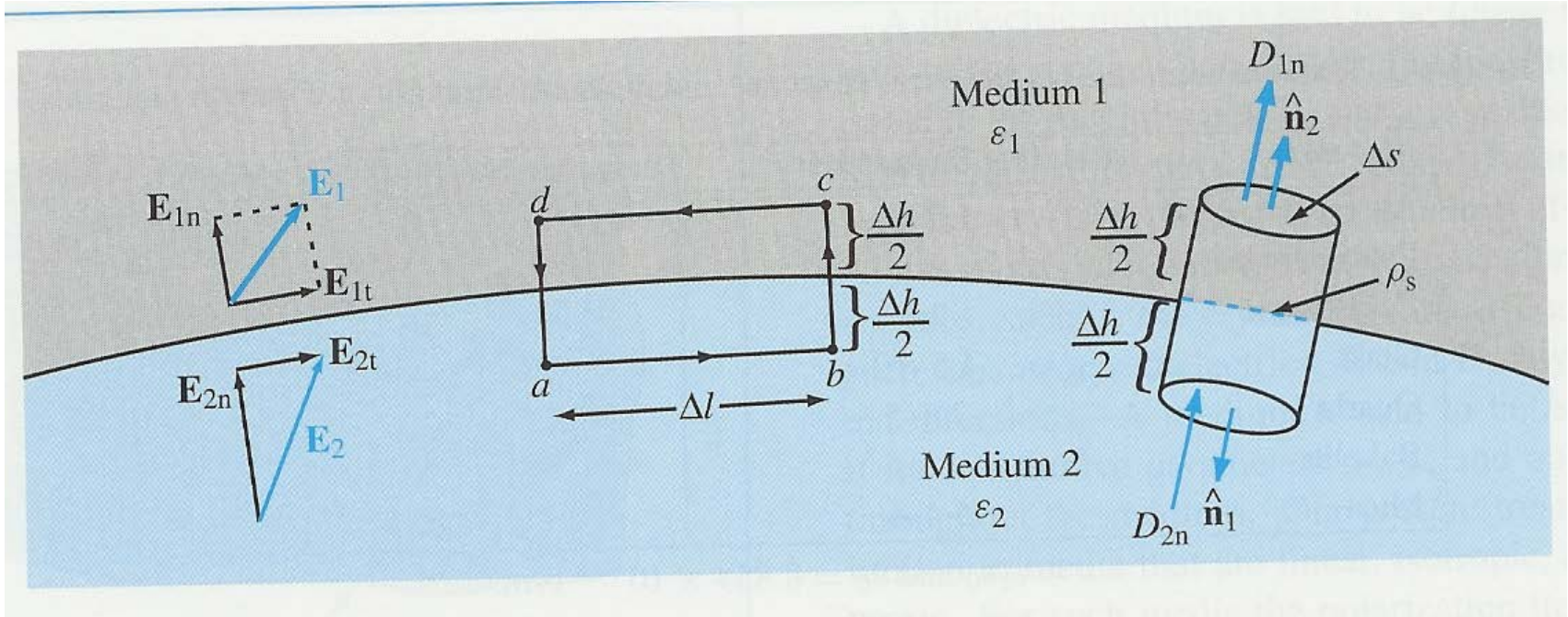
$\epsilon = \epsilon_0 (1 + \chi_e)$, Electric susceptibility

Relative permittivity: $\epsilon_r = \frac{\epsilon}{\epsilon_0} = (1 + \chi_e),$

Dielectric breakdown

16.360 Lecture 20

Electric boundary condition



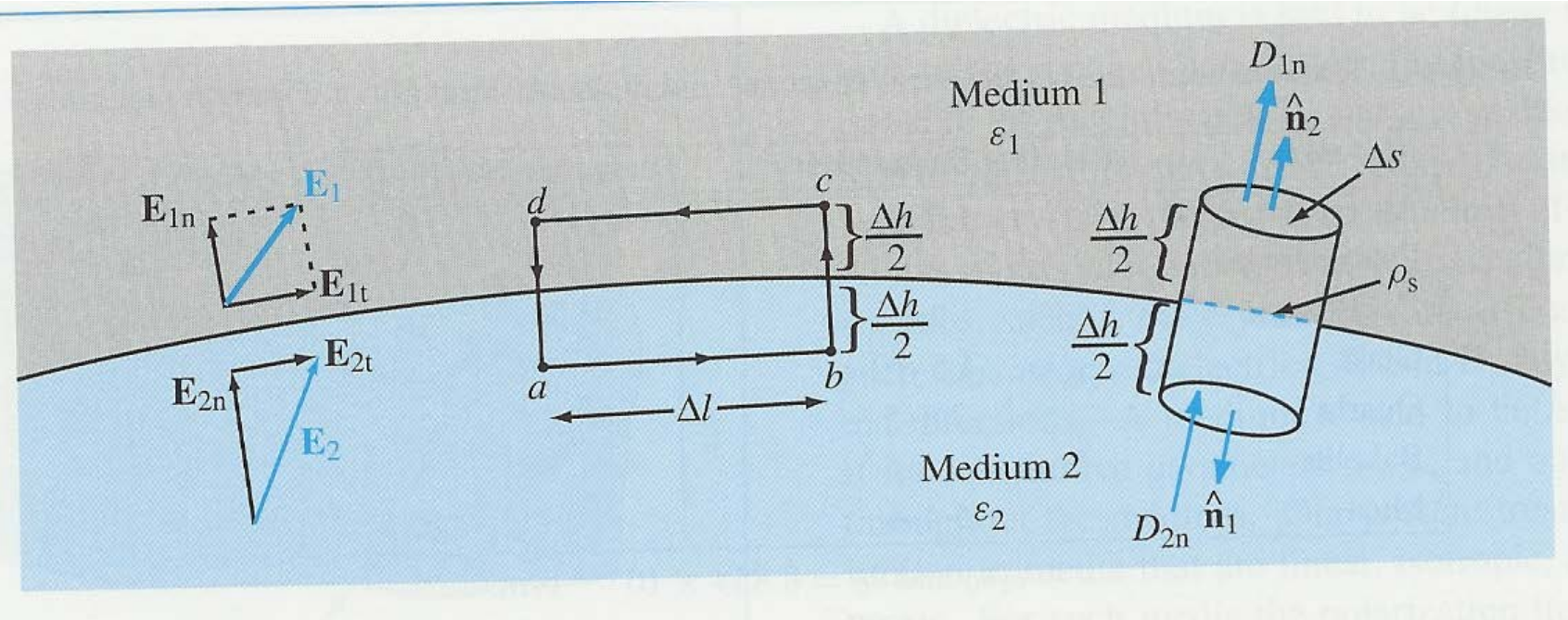
$\oint_C \vec{E} \cdot d\vec{l} = \lim_{\Delta h \rightarrow 0} \left[\int_a^b \vec{E}_2 \cdot d\vec{l} + \int_c^d \vec{E}_1 \cdot d\vec{l} \right] = 0;$ the tangential component is continuous across the boundary of two media.

$$\vec{E}_1 = \vec{E}_{1t} + \vec{E}_{1n}, \quad \vec{E}_{1t} \Delta l - \vec{E}_{2t} \Delta l = 0,$$

$$\vec{E}_2 = \vec{E}_{2t} + \vec{E}_{2n}, \quad \vec{E}_{1t} = \vec{E}_{2t},$$

16.360 Lecture 20

Electric boundary condition



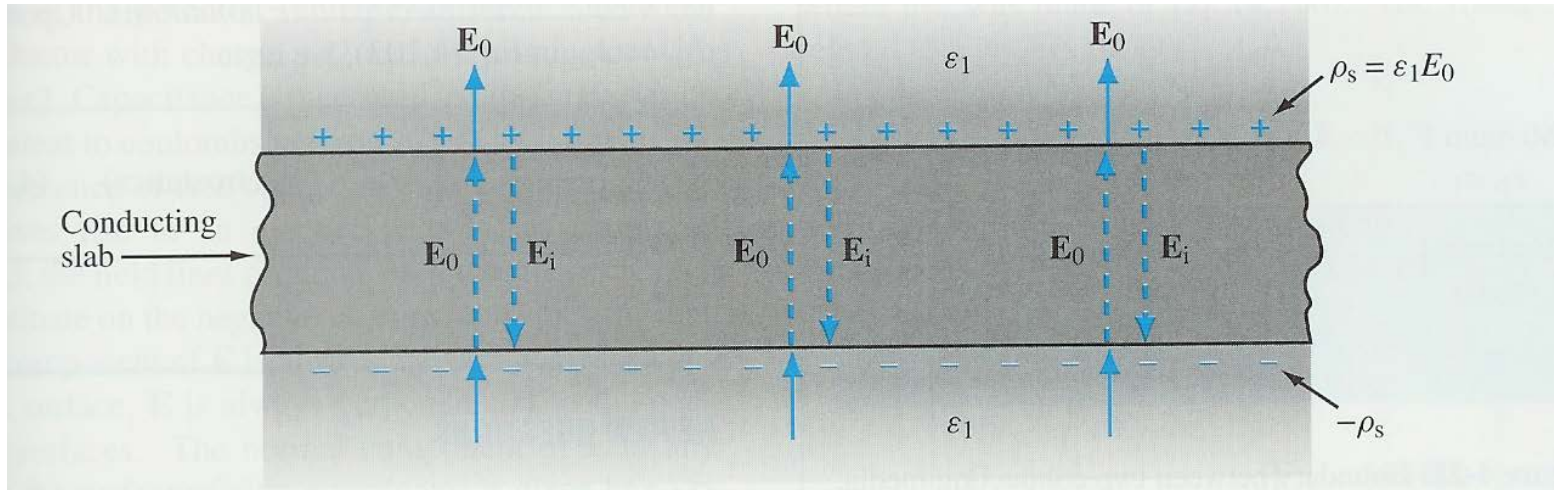
$\oint_C \vec{D} \cdot d\vec{s} = \lim_{\Delta h \rightarrow 0} \left[\int_{top} \vec{D} \cdot d\vec{s} + \int_{bottom} \vec{D} \cdot d\vec{s} \right] = \rho_s \Delta s$; the normal component of D changes, the amount of change is equal to the surface Charge density.

$$\vec{D}_{1n} \Delta s - \vec{D}_{2n} \Delta s = \rho_s \Delta s,$$

$$\vec{D}_{1n} - \vec{D}_{2n} = \rho_s,$$

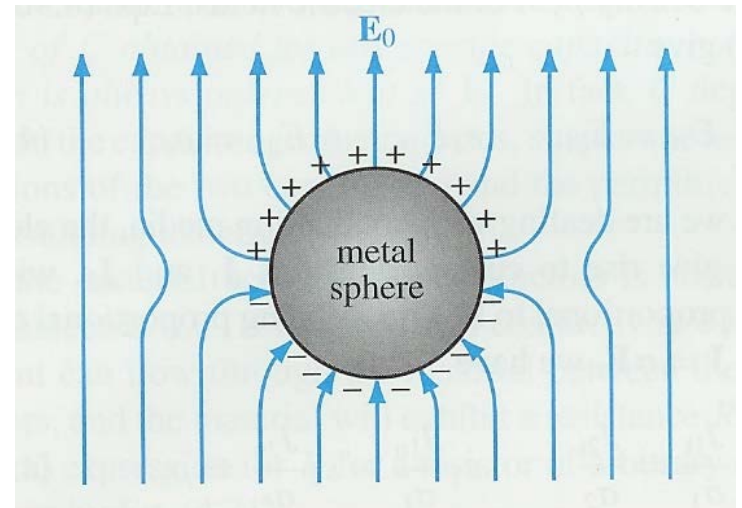
16.360 Lecture 20

Dielectric-Conductor boundary

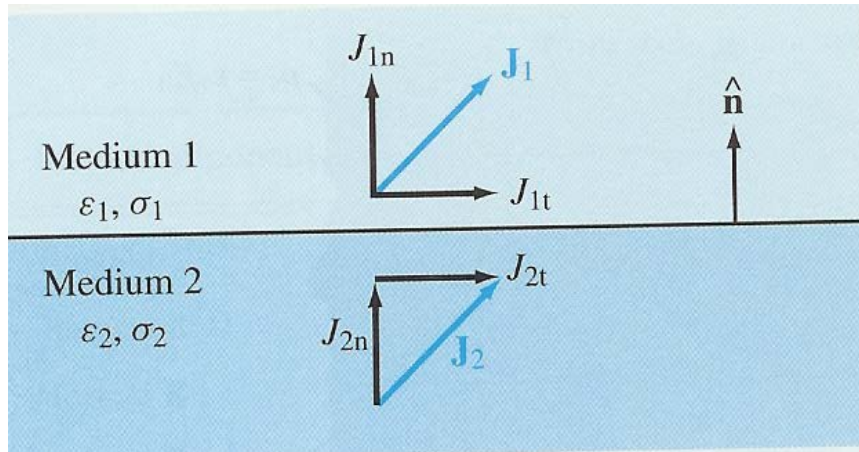


$$\vec{E}_{1t} = \vec{E}_{2t} = 0,$$

$$\vec{D}_{1n} = \rho_s,$$



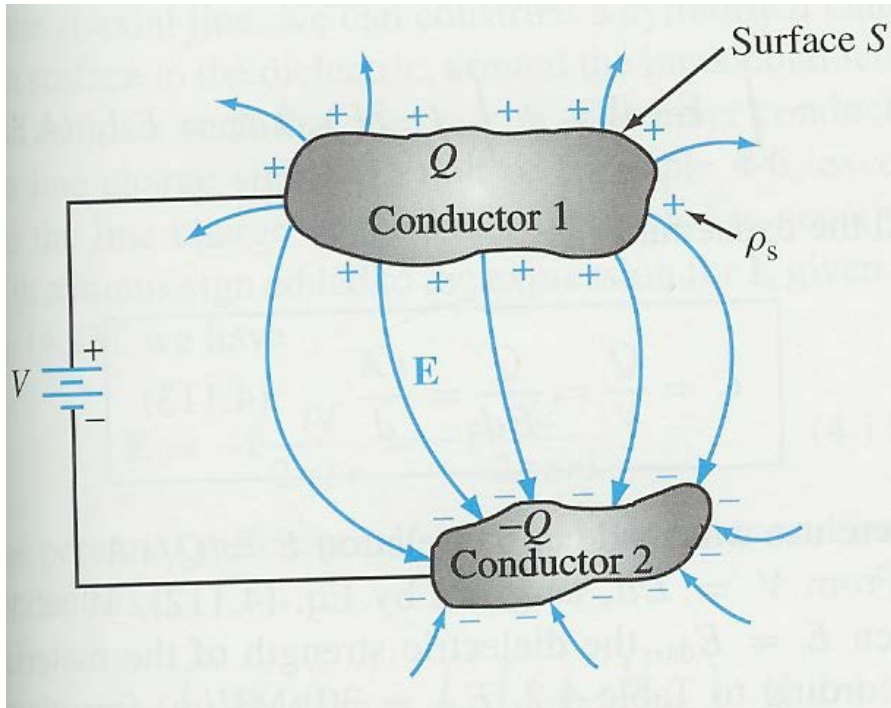
Conductor-Conductor boundary



$$\vec{E}_{1t} = \vec{E}_{2t}, \quad \vec{D}_{1n} - \vec{D}_{2n} = \epsilon_1 \vec{E}_{1n} - \epsilon_2 \vec{E}_{2n} = \rho_s,$$

$$\frac{\vec{J}_{1t}}{\sigma_1} = \frac{\vec{J}_{2t}}{\sigma_2}, \quad \epsilon_1 \frac{\vec{J}_{1n}}{\sigma_1} - \epsilon_2 \frac{\vec{J}_{2n}}{\sigma_2} = \rho_s,$$

$$\vec{J}_{1n} \left(\frac{\epsilon_1}{\sigma_1} - \frac{\epsilon_2}{\sigma_2} \right) = \rho_s,$$

Capacitance

$$C = \frac{Q}{V},$$

$$Q = \int_s \epsilon \vec{E} \cdot d\vec{s},$$

$$V = -\int_l \vec{E} \cdot d\vec{l}$$

$$C = \frac{\int_s \epsilon \vec{E} \cdot d\vec{s}}{-\int_l \vec{E} \cdot d\vec{l}},$$

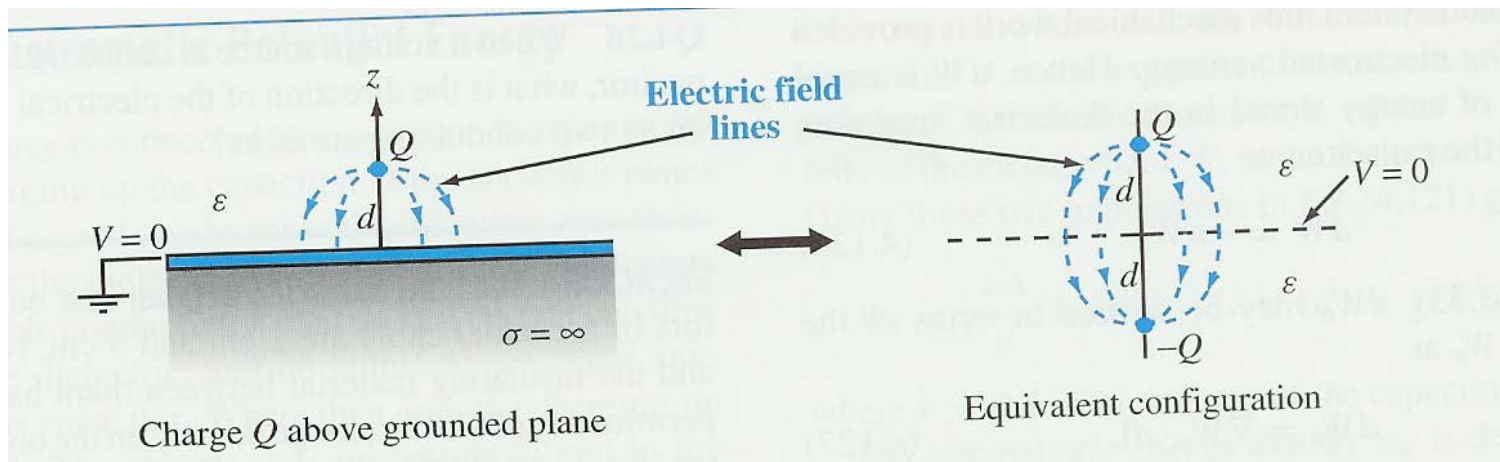
$$R = \frac{V}{I} = \frac{-\int_{x_2}^{x_1} \vec{E} \cdot d\vec{l}}{\int_A \vec{J} \cdot d\vec{s}} = \frac{-\int_{x_2}^{x_1} \vec{E} \cdot d\vec{l}}{\int_A \sigma \vec{E} \cdot d\vec{s}},$$

$$RC = \frac{\epsilon}{\sigma},$$

Electrostatic Potential Energy

$$W_e = \frac{1}{2} \vec{D} \cdot \vec{E}, \quad dW_e = -\vec{F} \cdot d\vec{l} = \nabla W_e \cdot d\vec{l},$$

$$\vec{F} = -\nabla W_e,$$

Image Method

Any given charge above an infinite, perfect conducting plane is electrically equivalent to the combination of the give charge and it's image with conducting plane removed.

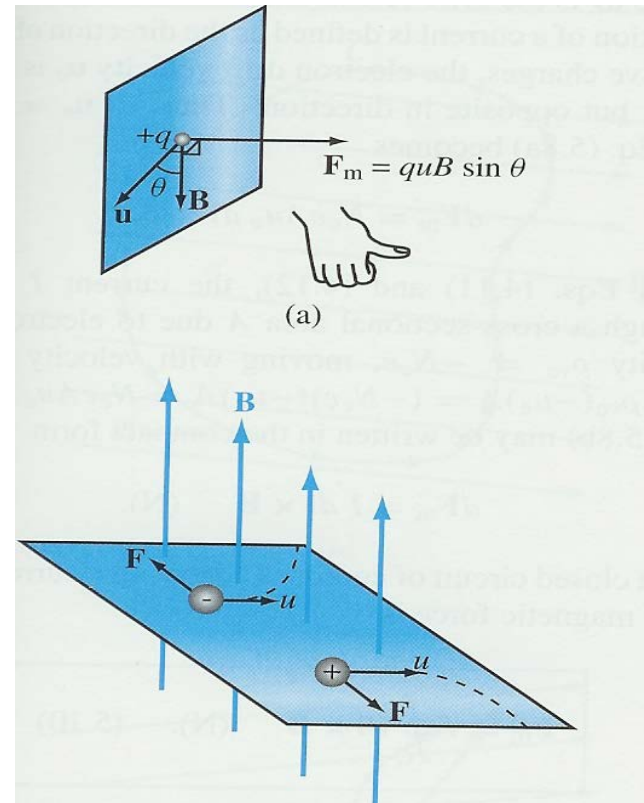
Magnetostatics

$$\begin{aligned} \nabla \cdot \vec{B} &= 0, & \vec{B} &= \mu \vec{H}, \\ \nabla \times \vec{H} &= \vec{J}, & \mu &= \mu_r \mu_0, \end{aligned}$$

Magnetic forces and Torques

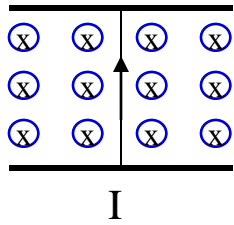
$$\vec{F}_m = q\vec{u} \times \vec{B},$$

$$\vec{F} = \vec{F}_e + \vec{F}_m = q\vec{E} + q\vec{u} \times \vec{B},$$



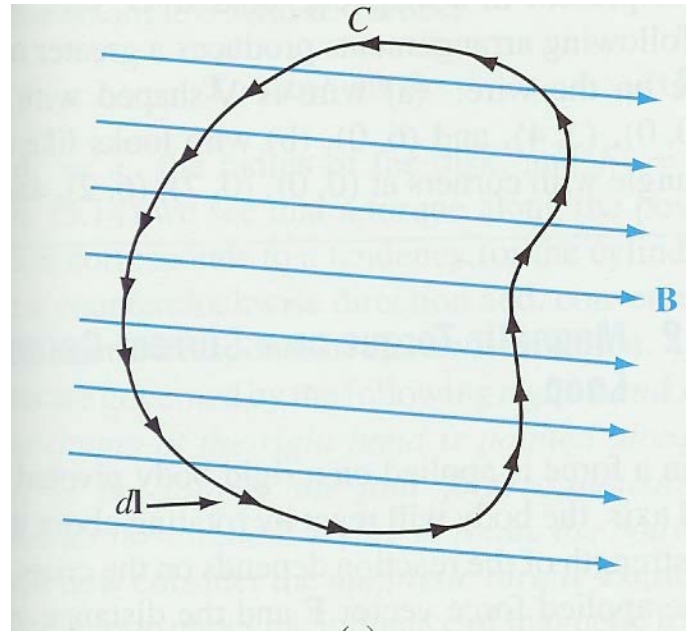
16.360 Lecture 21

Magnetic forces on a current-carrying conductor



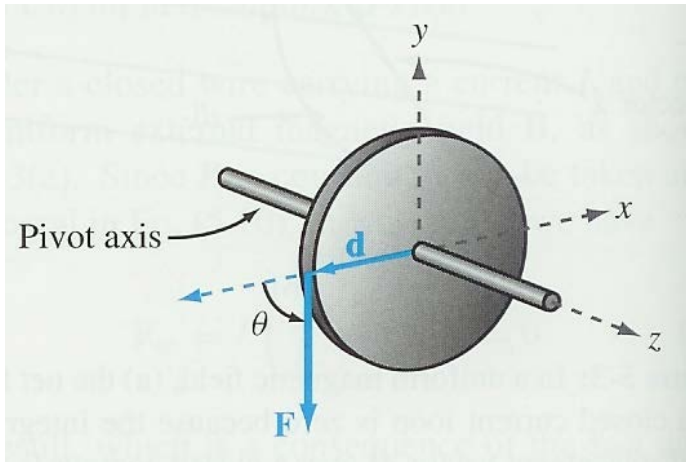
$$dF_m = Id\vec{l} \times \vec{B},$$

$$F_m = I \left(\oint_C d\vec{l} \times \vec{B} \right),$$



16.360 Lecture 21

Magnetic torques on a current-carrying conductor



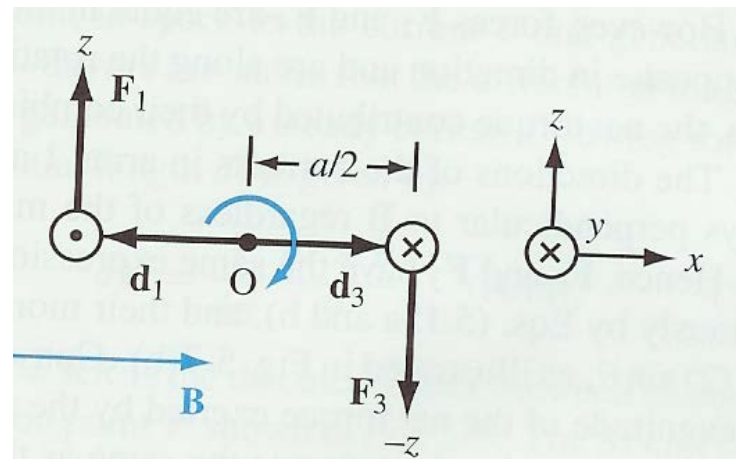
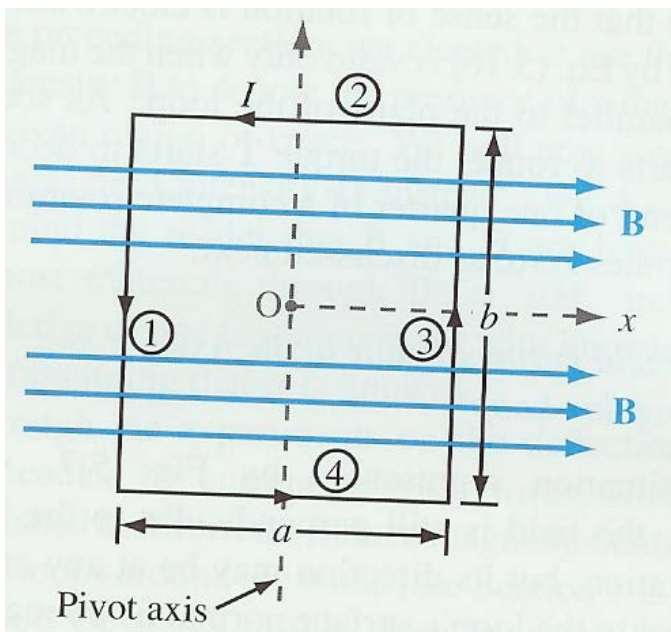
$$\vec{T} = \vec{d} \times \vec{F},$$

$$\vec{F}_1 = I(-\hat{y}b) \times (\hat{x}B_0) = \hat{z}IbB_0,$$

$$\vec{F}_3 = I(\hat{y}b) \times (\hat{x}B_0) = -\hat{z}IbB_0,$$

$$\vec{T} = \vec{d}_1 \times \vec{F}_1 + \vec{d}_3 \times \vec{F}_3 = \left(-\hat{x}\frac{a}{2}\right) \times (\hat{z}IbB_0)$$

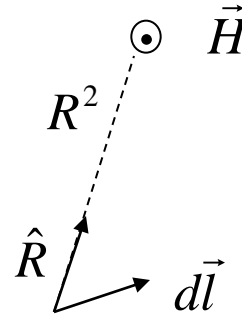
$$+ \left(\hat{x}\frac{a}{2}\right) \times (-\hat{z}IbB_0) = \hat{y}IabB_0 = \hat{y}IAB_0,$$



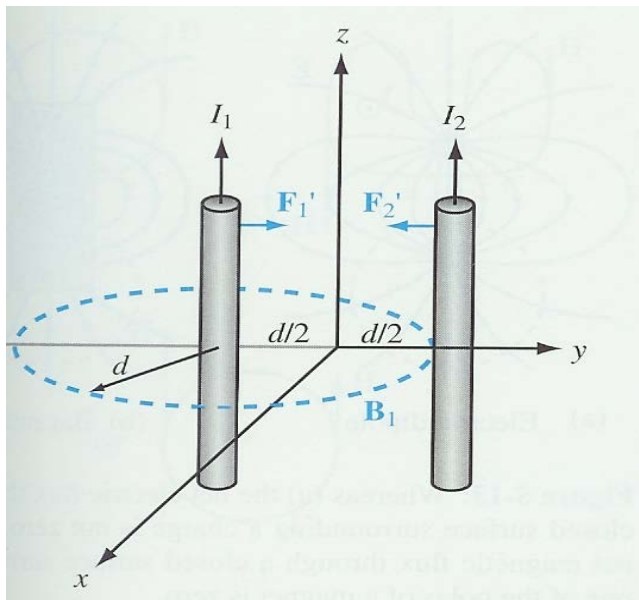
The Biot-Savart Law

$$d\vec{H} = \frac{I}{4\pi} \frac{d\vec{l} \times \hat{R}}{R^2},$$

$$\vec{H} = \frac{I}{4\pi} \oint_l \frac{d\vec{l} \times \hat{R}}{R^2},$$



Magnetic force between two parallel conductor



$$\vec{B}_1 = -\hat{x} \frac{\mu_0 I_1}{2\pi d},$$

$$\vec{F}_2 = I_2 l \hat{z} \times \vec{B}_1 = I_2 l \hat{z} \times (-\hat{x}) \frac{\mu_0 I_1}{2\pi d}$$

$$= -\hat{y} \frac{\mu_0 I_1 I_2 l}{2\pi d},$$

Gauss's Law for Magnetism

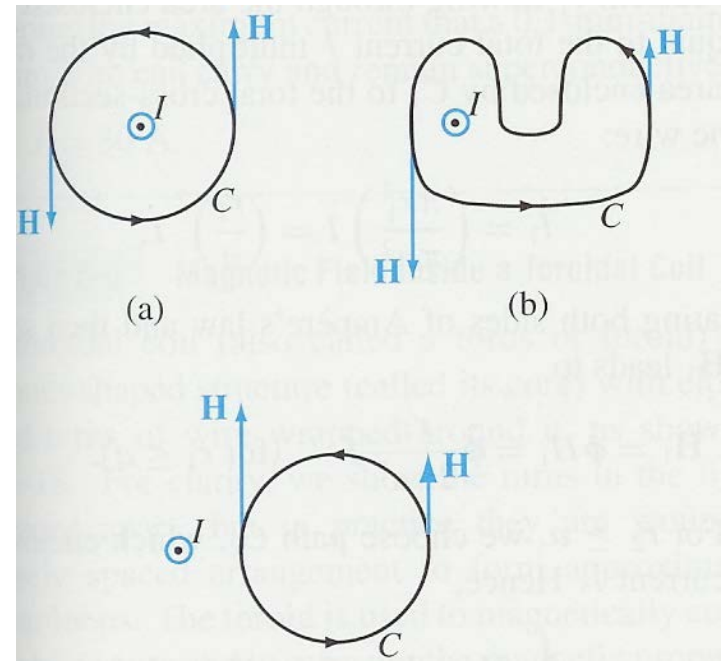
$$\nabla \cdot \vec{D} = \rho_v, \quad \oint_S \vec{D} \cdot d\vec{s} = Q_v,$$

$$\nabla \cdot \vec{B} = 0, \quad \oint_S \vec{B} \cdot d\vec{s} = 0,$$

Ampere's Law

$$\nabla \times \vec{H} = \vec{J}, \quad \int_S (\nabla \times \vec{H}) \cdot d\vec{s} = \int_S \vec{J} \cdot d\vec{s} = I,$$

$$\oint_C \vec{H} \cdot d\vec{l} = \int_S \vec{J} \cdot d\vec{s} = I,$$



Vector Magnet Potential

$$\begin{array}{ll} \nabla \cdot \vec{B} = 0, & \vec{B} = \mu \vec{H}, \\ \nabla \times \vec{H} = \vec{J}, & \mu = \mu_r \mu_0, \end{array}$$

For any vector A

$$\nabla \cdot (\nabla \times \vec{A}) = 0, \quad \nabla \cdot \vec{B} = 0, \quad \text{Define: } \vec{B} \equiv \nabla \times \vec{A}, \quad \text{with } \nabla \cdot \vec{A} = 0,$$

$$\nabla \times \vec{H} = \vec{J}, \quad \nabla \times (\nabla \times \vec{A}) = \mu \vec{J},$$

$$\nabla \times (\nabla \times \vec{A}) = \nabla(\nabla \cdot \vec{A}) - \nabla^2 \vec{A} = -\nabla^2 \vec{A} = \mu \vec{J},$$

$$\nabla^2 \vec{A} = -\mu \vec{J},$$

Vector Poisson's equation

$$\vec{A} = \frac{\mu}{4\pi} \int_{v'} \frac{\vec{J}}{R'} dv',$$

16.360 Lecture 22

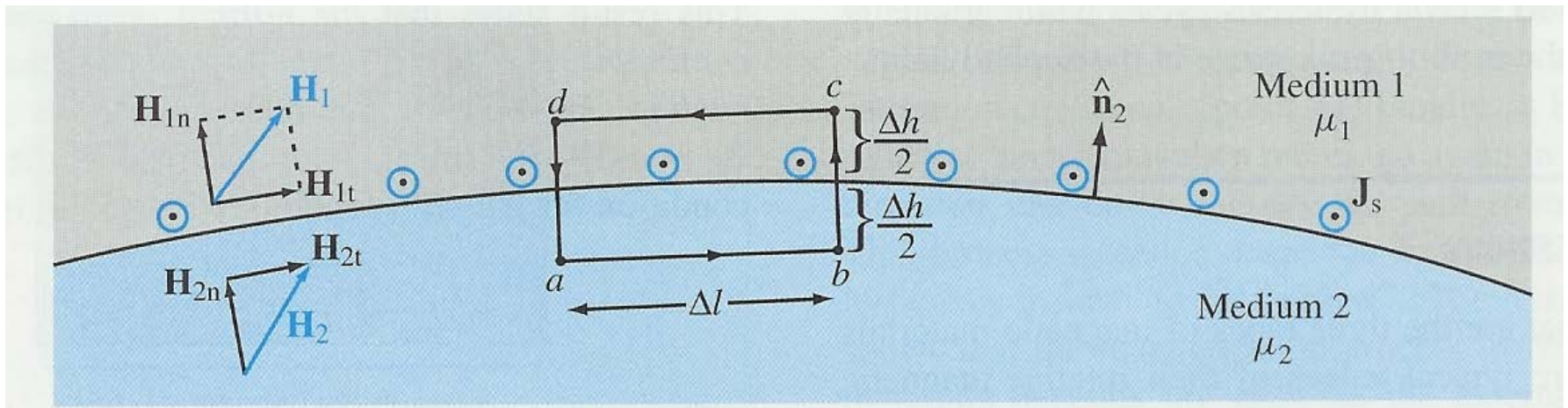
Magnetic boundary conditions

$$\oint_C \vec{D} \cdot d\vec{s} = \rho_s; \quad \vec{D}_{1n} - \vec{D}_{2n} = \rho_s,$$

$$\nabla \cdot \vec{B} = 0, \quad \oint_C \vec{B} \cdot d\vec{s} = 0; \quad \vec{B}_{1n} - \vec{B}_{2n} = 0,$$

$$\oint_C \vec{E} \cdot d\vec{l} = 0; \quad \vec{E}_{1t} = \vec{E}_{2t},$$

$$\nabla \times \vec{H} = \vec{J}, \quad \oint_C \vec{H} \cdot d\vec{l} = J; \quad \vec{H}_{1t} - \vec{H}_{2t} = J,$$

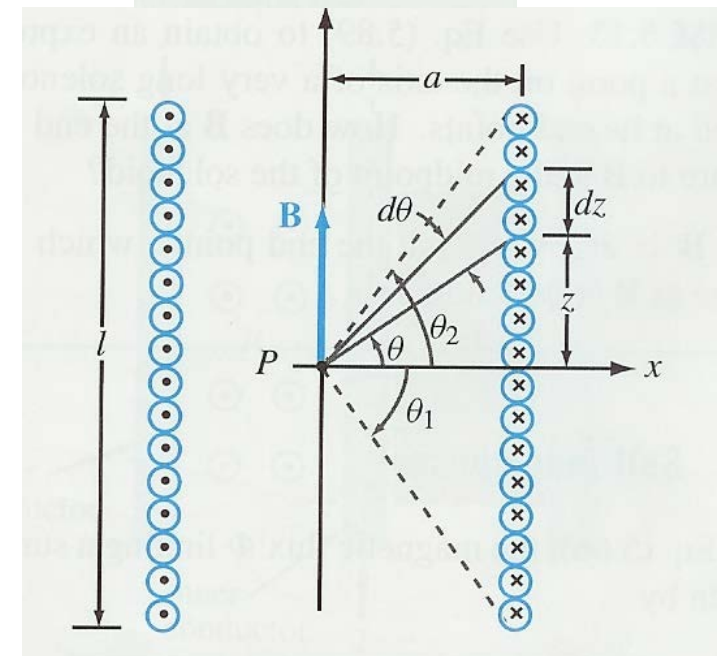
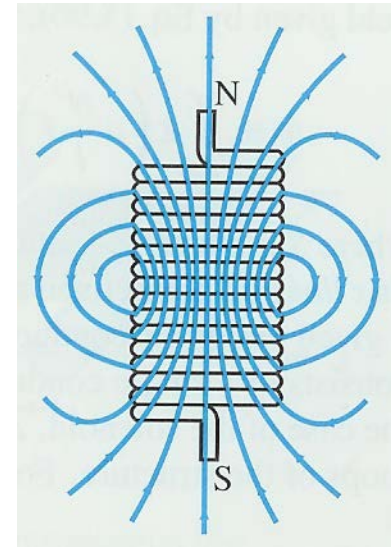


InductanceMagnetic field in a Solenoid

$$\vec{H} = \hat{z} \frac{I' a^2}{2(a^2 + z^2)^{3/2}},$$

$$d\vec{B} = \mu dH = \hat{z} \frac{\mu n I' a^2}{2(a^2 + z^2)^{3/2}} dz,$$

$$\vec{B} = \hat{z} \frac{\mu n I'}{2} \sin(\theta_2 - \theta_1),$$



Self Inductance

$$\Phi = \int_s \vec{B} \cdot d\vec{s} = \mu \frac{NIS}{l},$$

$$\Lambda = N\Phi = N \int_s \vec{B} \cdot d\vec{s} = \mu \frac{N^2 IS}{l},$$

$$L = \frac{\Lambda}{I} = \frac{N \int_s \vec{B} \cdot d\vec{s}}{I} = \mu \frac{N^2 S}{l},$$

Mutual Inductance

$$L_{12} = \frac{\Lambda_{12}}{I_1} = \frac{N_2 \int_s \vec{B}_1 \cdot d\vec{s}}{I_1}, \quad L_{21} = \frac{\Lambda_{21}}{I_2} = \frac{N_1 \int_s \vec{B}_2 \cdot d\vec{s}}{I_2},$$

Magnetic Energy

$$W_m = \frac{1}{2} \int_{v'} \vec{B} \cdot \vec{H} dv',$$

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Static field

$$\nabla \cdot \vec{D} = \rho_v, \quad \vec{E} = \int d\vec{E} = \frac{1}{4\pi\epsilon} \int_{v'} \hat{R}' \frac{\rho_v dv'}{R'^2},$$

$$\nabla \times \vec{E} = 0,$$

$$\nabla \cdot \vec{B} = 0, \quad \vec{H} = \frac{I}{4\pi} \oint_l \frac{d\vec{l} \times \hat{R}}{R^2},$$
$$\nabla \times \vec{H} = \vec{J},$$

Dynamic Field

$$\nabla \cdot \vec{D} = \rho_v, \quad \nabla \cdot \vec{B} = 0,$$
$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}, \quad \nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t},$$

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Table 6-1: Maxwell's equations.

Reference	Differential Form	Integral Form
Gauss's law	$\nabla \cdot \mathbf{D} = \rho_v$	$\oint_S \mathbf{D} \cdot d\mathbf{s} = Q$ (6.1)
Faraday's law	$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$	$\oint_C \mathbf{E} \cdot d\mathbf{l} = -\int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{s}$ (6.2)*
No magnetic charges (Gauss's law for magnetism)	$\nabla \cdot \mathbf{B} = 0$	$\oint_S \mathbf{B} \cdot d\mathbf{s} = 0$ (6.3)
Ampère's law	$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}$	$\oint_C \mathbf{H} \cdot d\mathbf{l} = \int_S \left(\mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \right) \cdot d\mathbf{s}$ (6.4)

*For a stationary surface S .

Faraday's Law

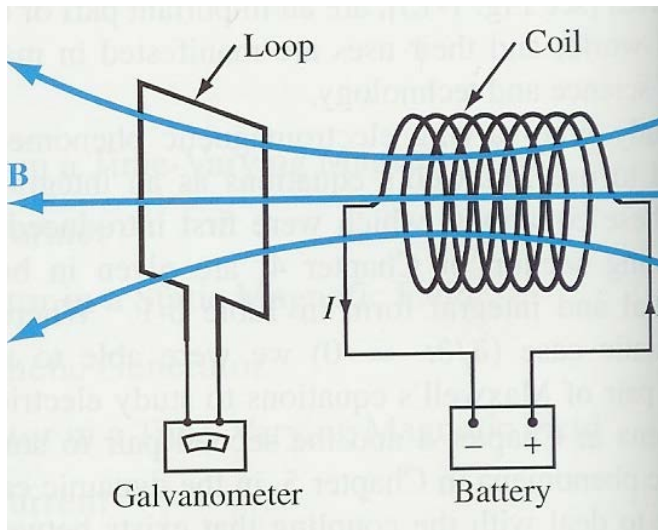
$$\oint_C \vec{E} \cdot d\vec{l} = -\oint_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s} \quad \Phi = \int_s \vec{B} \cdot d\vec{s},$$

$$= -\frac{\partial}{\partial t} \oint_s \vec{B} \cdot d\vec{s}, = -\frac{\partial \Phi}{\partial t},$$

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Electromotive force

$$V_{emf} = \oint_C \vec{E} \cdot d\vec{l} = - \oint_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s} = - \frac{\partial \Phi}{\partial t},$$



$$V_{emf} = V_{emf}^{tr} + V_{emf}^m,$$

Stationary Loop in a Time-varying Magnetic field

$$V_{emf}^{tr} = - \oint_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s}, \quad I = \frac{V_{emf}^{tr}}{R_i + R} = - \oint_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s}, \quad \underline{\text{Lenz's law}}$$

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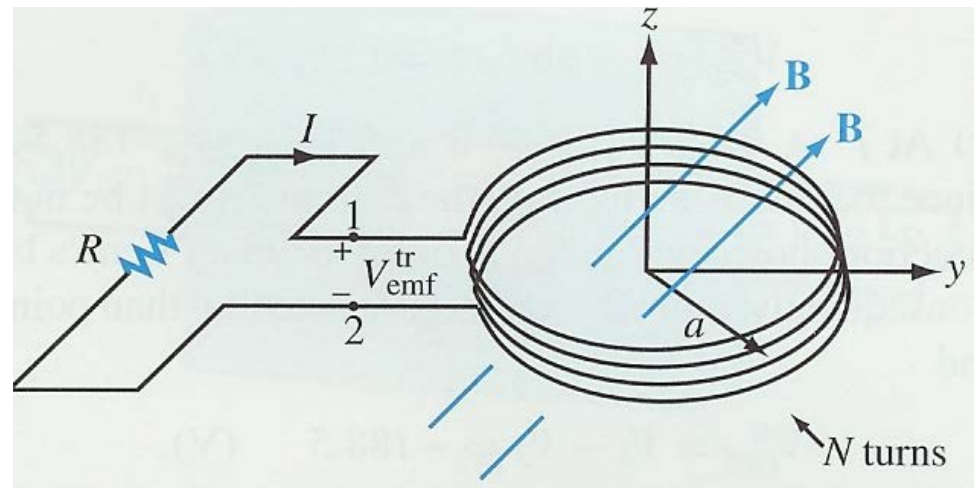
$$V_{emf} = \oint_C \vec{E} \cdot d\vec{l} = - \oint_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s} = \oint_s (\nabla \times \vec{E}) \cdot d\vec{s},$$

$$\nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t}, \quad \underline{\text{Faraday's law, differential form}}$$

An example:

$$\vec{B} = B_0 (\hat{y}2 + \hat{z}3) \sin \omega t,$$

- The magnetic flux link of a single turn of the inductor.
- The transformer emf,.
- The polarity of the emf.
- The induced current.

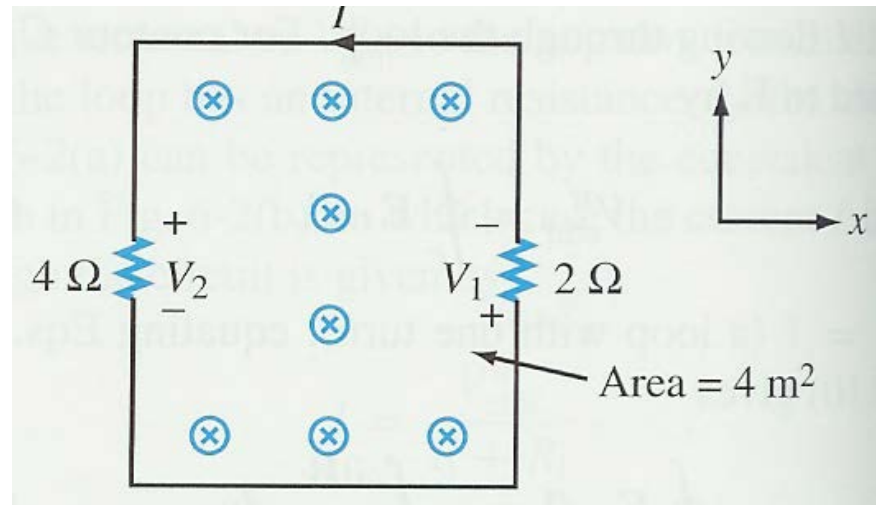


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Example II

$$\vec{B} = -\hat{z}0.3t,$$

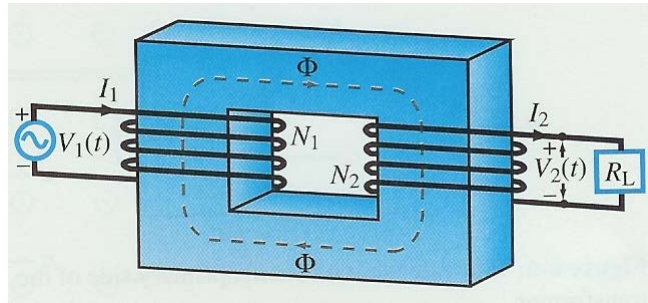
Determine the voltage drops across the two resistors



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The ideal Transformer properties:

- $\mu = \infty$
- $I = 0$ in the core.
- The magnetic flux is confined within the core



Questions:

- $I = ?$, with applied voltage of V_1 and with R_L
- V_2 , and $I_2 = ?$

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Voltage transformer:

$$V_1 = -N_1 \frac{d\Phi}{dt}, \quad V_2 = -N_2 \frac{d\Phi}{dt}, \quad \frac{V_2}{V_1} = \frac{N_2}{N_1},$$

Power relations:

$$P_1 = P_2, \quad \text{Why?}$$

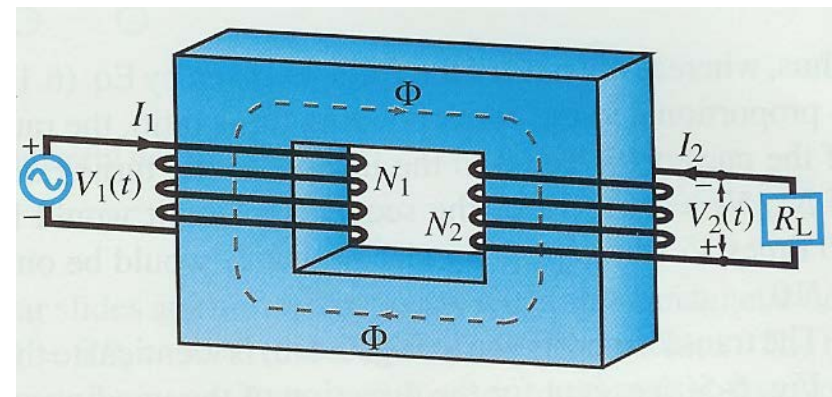
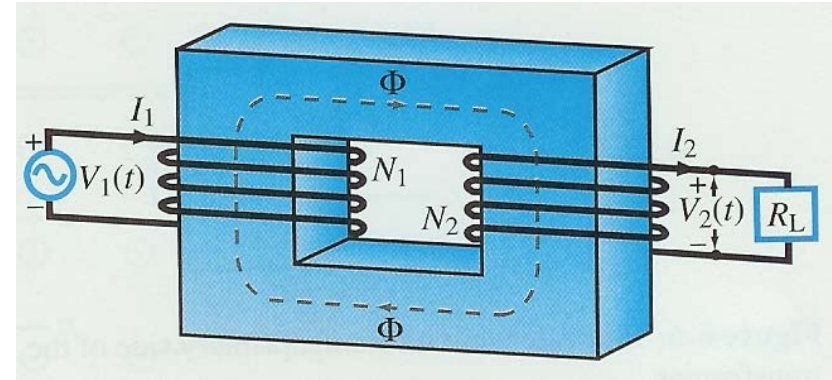
Current transformer:

$$P_1 = V_1 I_1, \quad P_2 = V_2 I_2, \quad \frac{I_2}{I_1} = \frac{N_1}{N_2},$$

Impedance transformer:

$$R_1 = V_1 / I_1, \quad R_2 = V_2 / I_2, \quad \frac{R_1}{R_2} = \left(\frac{N_1}{N_2}\right)^2,$$

$$R_{in} = \left(\frac{N_1}{N_2}\right)^2 R_L,$$



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Moving conductor in a static magnetic field:

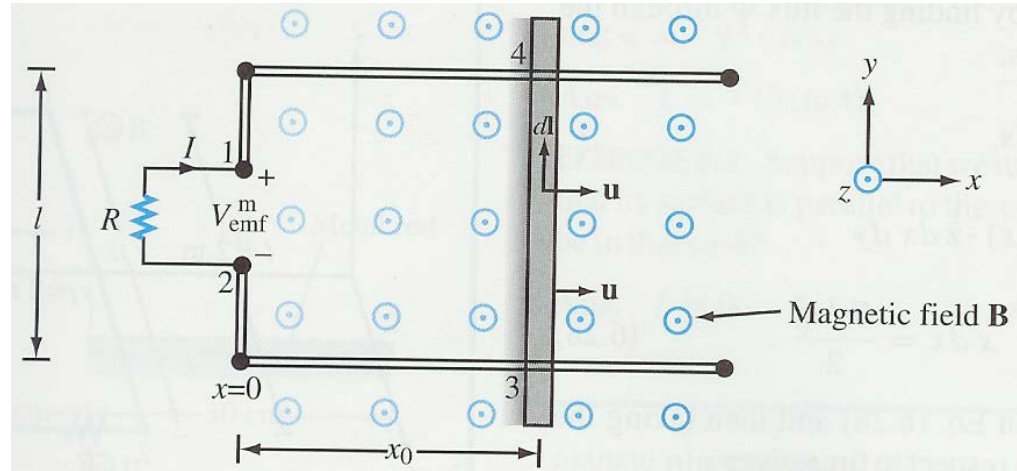
$$V_{emf} = \oint_C \vec{E} \cdot d\vec{l} = - \oint_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s} = - \frac{\partial \Phi}{\partial t},$$

$$V_{emf} = V_{emf}^{tr} + V_{emf}^m,$$

$$V_{emf}^m = - \frac{\partial \Phi}{\partial t} = - \frac{\partial (\vec{B} \cdot d\vec{s})}{\partial t} = (\vec{u} \times \vec{B}) \cdot d\vec{l},$$

$$d\vec{s} = \vec{w} \times d\vec{l},$$

$$\vec{A} \cdot (\vec{B} \times \vec{C}) = \vec{B} \cdot (\vec{A} \times \vec{C}) = \vec{C} \cdot (\vec{A} \times \vec{B}),$$

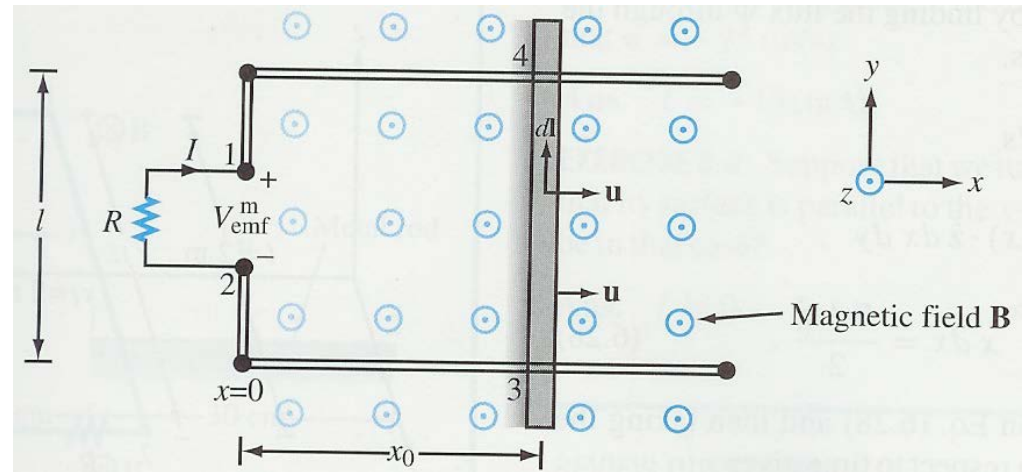


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Another way to look at it:

$$\mathbf{F}_m = q(\vec{u} \times \vec{B}), \quad \mathbf{E}_m = \frac{\mathbf{F}_m}{q},$$

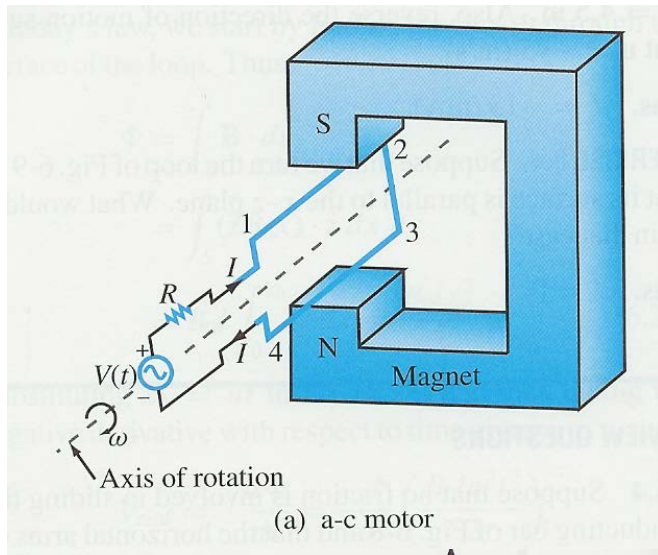
$$V_{emf}^m = \int_{l_1}^{l_2} \mathbf{E}_m \cdot d\vec{l} = \int_{l_1}^{l_2} (\vec{u} \times \vec{B}) \cdot d\vec{l},$$



Next lecture:

- The electromagnetic generator
- Moving conductor in a time varying magnetic field

The electromagnetic generator

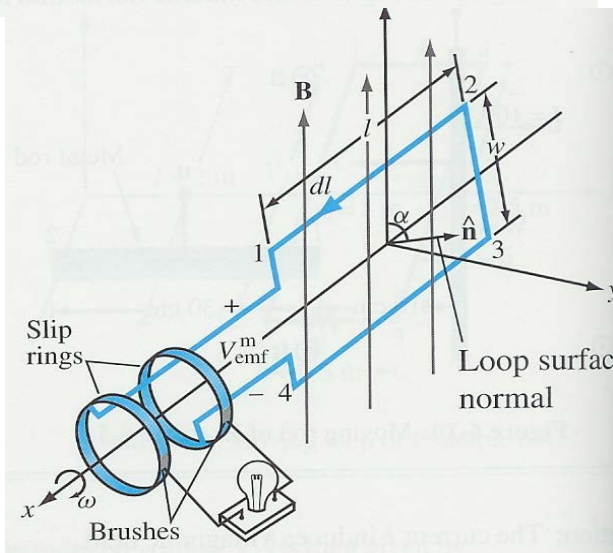


$$\Phi = \int_S \vec{B} \cdot d\vec{s} = B_0 A \cos(\alpha)$$

$$= B_0 A \cos(\omega t + C_0),$$

$$V_{emf} = -\frac{d\Phi}{dt} = -\frac{d}{dt} B_0 A \cos(\omega t + C_0)$$

$$= A \omega B_0 \sin(\omega t + C_0),$$

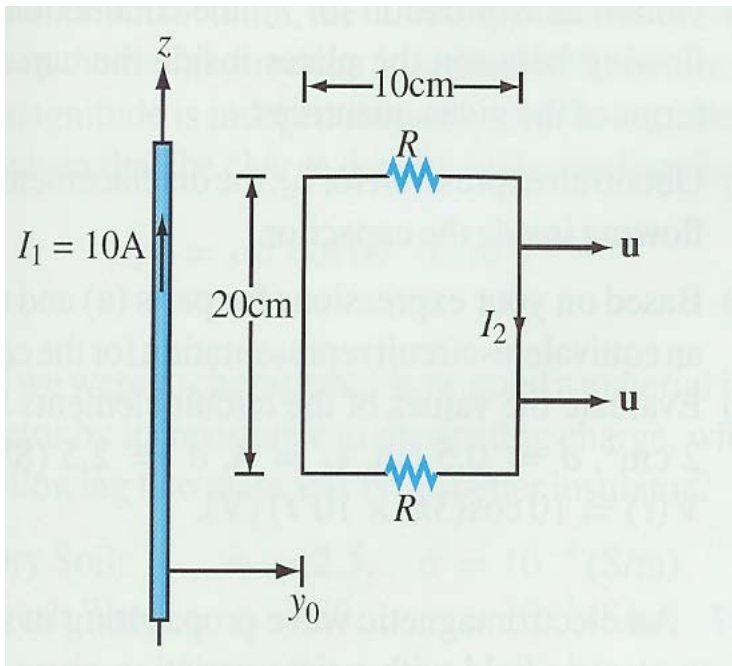


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Moving conductor in a time-varying magnetic field

$$\oint_C \vec{E} \cdot d\vec{l} = V_{emf} = V_{emf}^{tr} + V_{emf}^m = -\int_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s} + \oint_C (\vec{u} \times \vec{B}) \cdot d\vec{l},$$

Example:



$$I_1 = 10A, \quad \vec{u} = 5\hat{y}, \quad R = 10\Omega$$

$$I_2 = ?$$

Displacement current

- Ampere's law in static electric field

$$\nabla \times \vec{H} = \vec{J},$$

- Ampere's law in time-varying electric field

$$\nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t},$$

- proof of Ampere's law:

$$\nabla \cdot \vec{D} = \rho_v, \quad Q = \int_v (\nabla \cdot \vec{D}) dv = \oint_s \vec{D} \cdot d\vec{s},$$

$$I = \oint_s \vec{J}' \cdot d\vec{s} = \frac{\partial}{\partial t} Q = \frac{\partial}{\partial t} \int_v (\nabla \cdot \vec{D}) dv = \frac{\partial}{\partial t} \oint_s \vec{D} \cdot d\vec{s} = \oint_s \left(\frac{\partial}{\partial t} \vec{D} \right) \cdot d\vec{s},$$

$$J' = \frac{\partial \vec{D}}{\partial t}$$

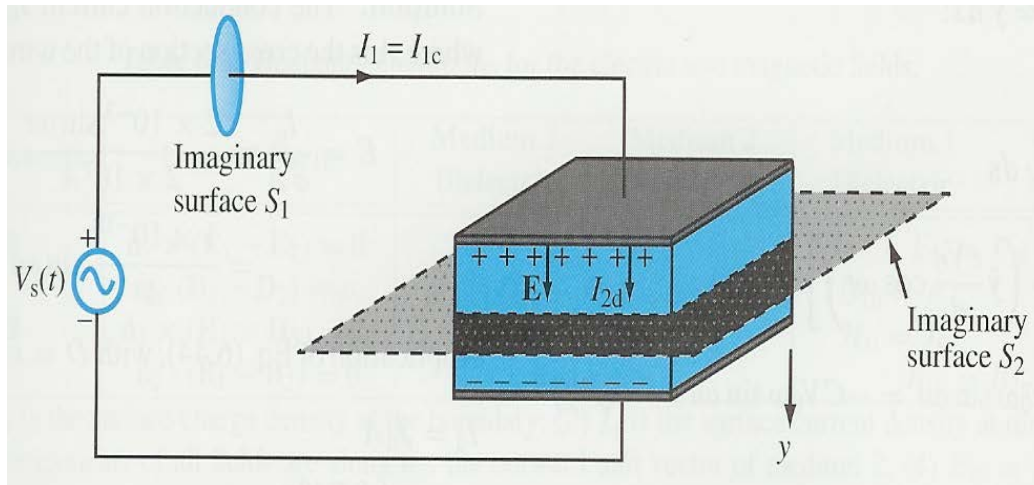
Displacement current density

Displacement current

- Ampere's law in time-varying electric field

$$\nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t},$$

Example:



$$V_s = V_0 \cos \omega t,$$

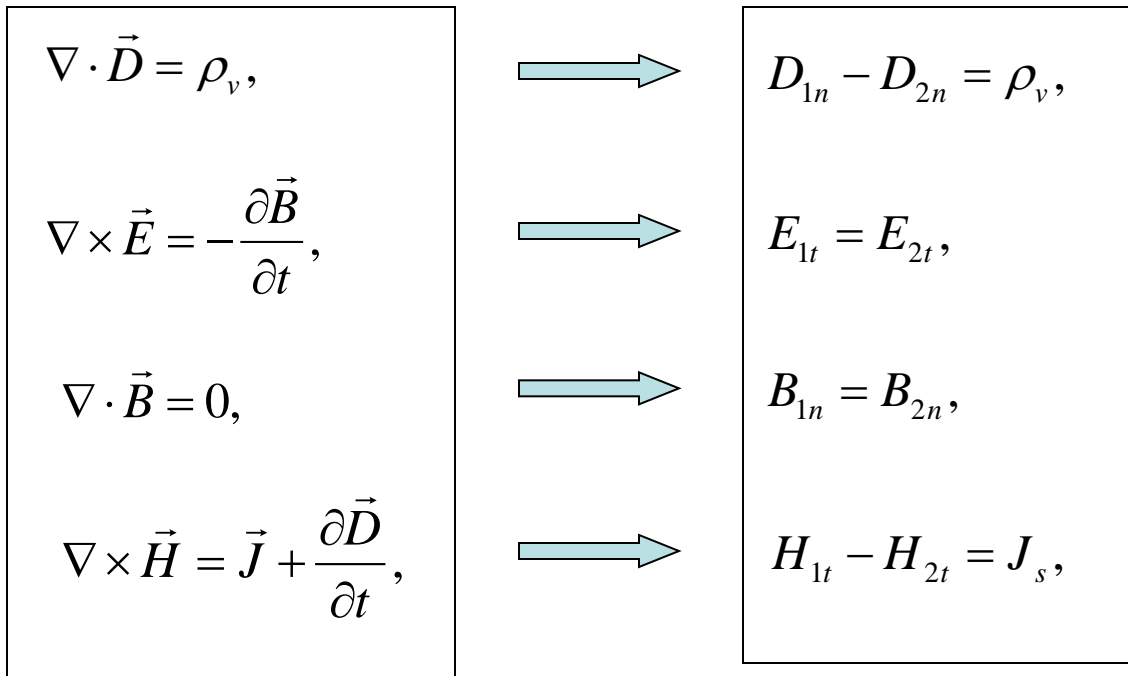
$$I_{1c} = C \frac{dV_c}{dt} = -CV_0 \omega \sin \omega t,$$

$$\vec{E} = \hat{y} \frac{V_c}{d} = \hat{y} \frac{V_0}{d} \cos \omega t,$$

$$I_{2d} = \int_s \frac{\partial \vec{D}}{\partial t} \cdot d\vec{s} = -CV_0 \omega \sin \omega t,$$

16.360 Lecture 28

- Boundary conditions for Electromagnetic



Maxwell equations

boundary conditions

- Charge-Current continuity Relation

$$I = -\frac{\partial}{\partial t} Q = -\frac{\partial}{\partial t} \int_v \rho_v dv,$$

$$I = \oint_s \vec{J} \cdot d\vec{s} = -\frac{\partial}{\partial t} \int_v \rho_v dv,$$

$$\oint_s \vec{J} \cdot d\vec{s} = \int_v \nabla \cdot \vec{J} = -\frac{\partial}{\partial t} \int_v \rho_v dv,$$

$$\nabla \cdot \vec{J} = -\frac{\partial}{\partial t} \rho_v,$$

charge current continuity equation

$$\oint_s \vec{J} \cdot d\vec{s} = 0,$$

steady state integral form

$$\sum_i I_i = 0,$$

Kirchhoff's current law

- Free-charge dissipation in a conductor

$$\nabla \cdot \vec{J} = -\frac{\partial}{\partial t} \rho_v, \quad \vec{J} = \sigma E,$$

$$\frac{\sigma}{\epsilon} \nabla \cdot \vec{D} = \nabla \cdot \vec{J} = -\frac{\partial}{\partial t} \rho_v,$$

$$\frac{\sigma}{\epsilon} \rho_v = -\frac{\partial}{\partial t} \rho_v,$$

$$\rho_v = \rho_{v0} e^{-t/\tau_r}, \quad \tau_r = \frac{\epsilon}{\sigma},$$

- Electromagnetic Potentials

Electrostatics: $\nabla \times \vec{E} = 0, \quad \vec{E} = -\nabla V,$

$$\nabla \cdot \vec{B} = 0, \quad \vec{B} = \nabla \times \vec{A},$$

Dynamic case:

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}, \quad \nabla \times \vec{E} = -\frac{\partial}{\partial t}(\nabla \times \vec{A}),$$

$$\vec{E} = -\nabla V - \frac{\partial \vec{A}}{\partial t},$$

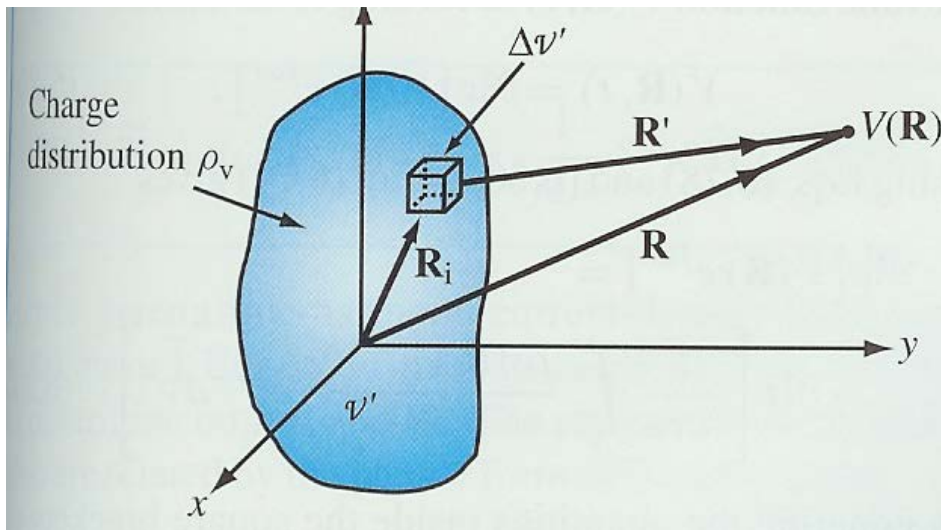
$$\vec{B} = \nabla \times \vec{A},$$

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• Retard Potentials

Electrostatics:
$$V(\vec{R}) = \frac{1}{4\pi\epsilon} \int_{v'} \frac{\rho_{v'}(R_i)}{\vec{R}'} dv',$$

Dynamic case:



$$V(\vec{R}, t) \neq \frac{1}{4\pi\epsilon} \int_{v'} \frac{\rho_{v'}(R_i, t)}{\vec{R}'} dv',$$

$$V(\vec{R}, t) = \frac{1}{4\pi\epsilon} \int_{v'} \frac{\rho_{v'}(R_i, t - R'/u_p)}{\vec{R}'} dv',$$

$$A(\vec{R}, t) = \frac{\mu}{4\pi} \int_{v'} \frac{\vec{J}(R_i, t - R'/u_p)}{\vec{R}'} dv',$$

- Time-Harmonic Potentials

$$\rho_{v'}(R_i, t) = \operatorname{Re} \left[\tilde{\rho}_{v'}(R_i) e^{j\omega t} \right],$$

$$\rho_{v'}(R_i, t - R'/u_p) = \operatorname{Re} \left[\tilde{\rho}_{v'}(R_i) e^{j\omega t - j\omega R'/u_p} \right] = \operatorname{Re} \left[\tilde{\rho}_{v'}(R_i) e^{j\omega t - jkR'} \right],$$

$$V(\vec{R}, t) = \operatorname{Re} \left[\frac{1}{4\pi\epsilon} \int_{v'} \frac{\rho_{v'}(R_i, t - R'/u_p)}{\vec{R}'} dv' \right] = \operatorname{Re} \left[\frac{1}{4\pi\epsilon} \int_{v'} \frac{\tilde{\rho}_{v'}(R_i) e^{j\omega t} e^{-jkR'}}{\vec{R}'} dv' \right],$$

$$\tilde{V}(\vec{R}) = \frac{1}{4\pi\epsilon} \int_{v'} \frac{\tilde{\rho}_{v'}(R_i) e^{-jkR'}}{\vec{R}'} dv',$$

$$\tilde{A}(\vec{R}) = \frac{1}{4\pi\epsilon} \int_{v'} \frac{\tilde{J}_{v'}(R_i) e^{-jkR'}}{\vec{R}'} dv',$$

- Time-Harmonic Potentials

$$\nabla \times \tilde{\vec{H}} = j\omega\epsilon \tilde{\vec{E}}, \quad \nabla \times \tilde{\vec{E}} = -j\omega\mu \tilde{\vec{H}},$$

$$\nabla \times (\nabla \times \tilde{\vec{E}}) = \nabla(\nabla \cdot \tilde{\vec{E}}) - \nabla^2 \tilde{\vec{E}} = -\omega^2 \mu\epsilon \tilde{\vec{E}},$$

$\nabla \cdot \tilde{\vec{E}} = 0$, if no free charge, trans-wave, why?

$$\nabla^2 \tilde{\vec{E}} + \omega^2 \mu\epsilon \tilde{\vec{E}} = 0, \quad \tilde{\vec{E}}(\vec{R}') = \tilde{\vec{E}}_0 e^{-jkR'}, \quad k^2 = \omega^2 \mu\epsilon,$$

- example

$$\vec{E}(z,t) = \hat{x}10 \sin(10^{10}t - kz), \quad \text{find } k?$$