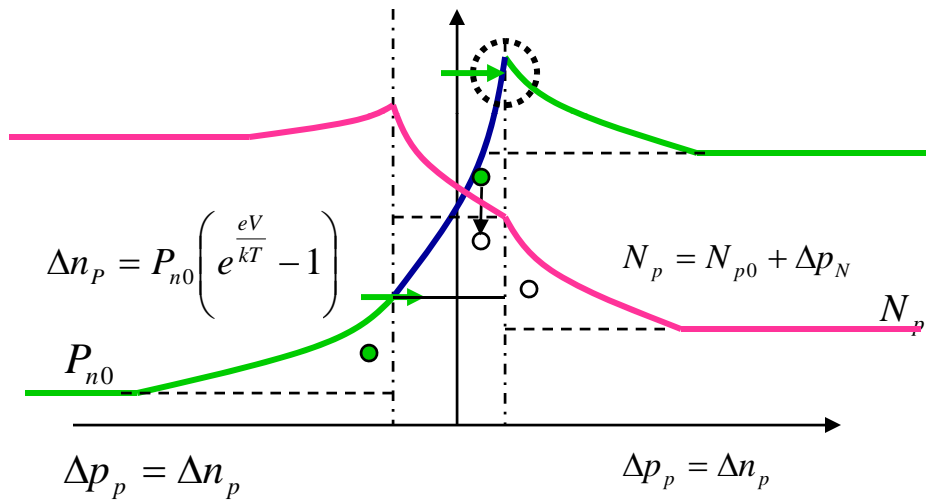


Example 3.3.1:

$$L_n = \sqrt{D_n \tau_n} \approx 13 \mu m$$

$$L_p = \sqrt{D_p \tau_p} \approx 3 \mu m$$

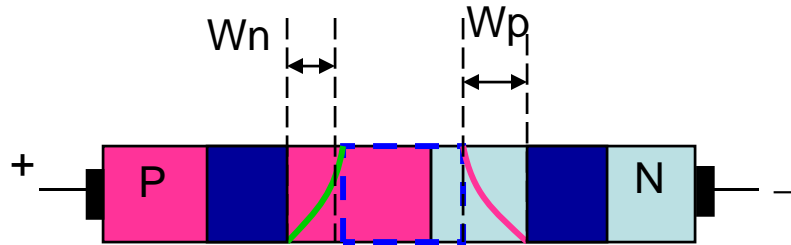
Minimum LED size? why?



$$\Delta p_N = N_{p0} \left( e^{\frac{eV}{kT}} - 1 \right) \quad J_p = \frac{qD_p}{L_p} \frac{n_i^2}{N_D} \left( e^{\frac{qV}{kT}} - 1 \right)$$

$$J_n = \frac{qD_n}{L_n} \frac{n_i^2}{N_A} \left( e^{\frac{qV}{kT}} - 1 \right)$$

$$J_r = \left( \frac{qn_i W}{2\tau_n} + \frac{qn_i W}{2\tau_p} \right) \left( e^{\frac{V}{\eta kT}} - 1 \right)$$



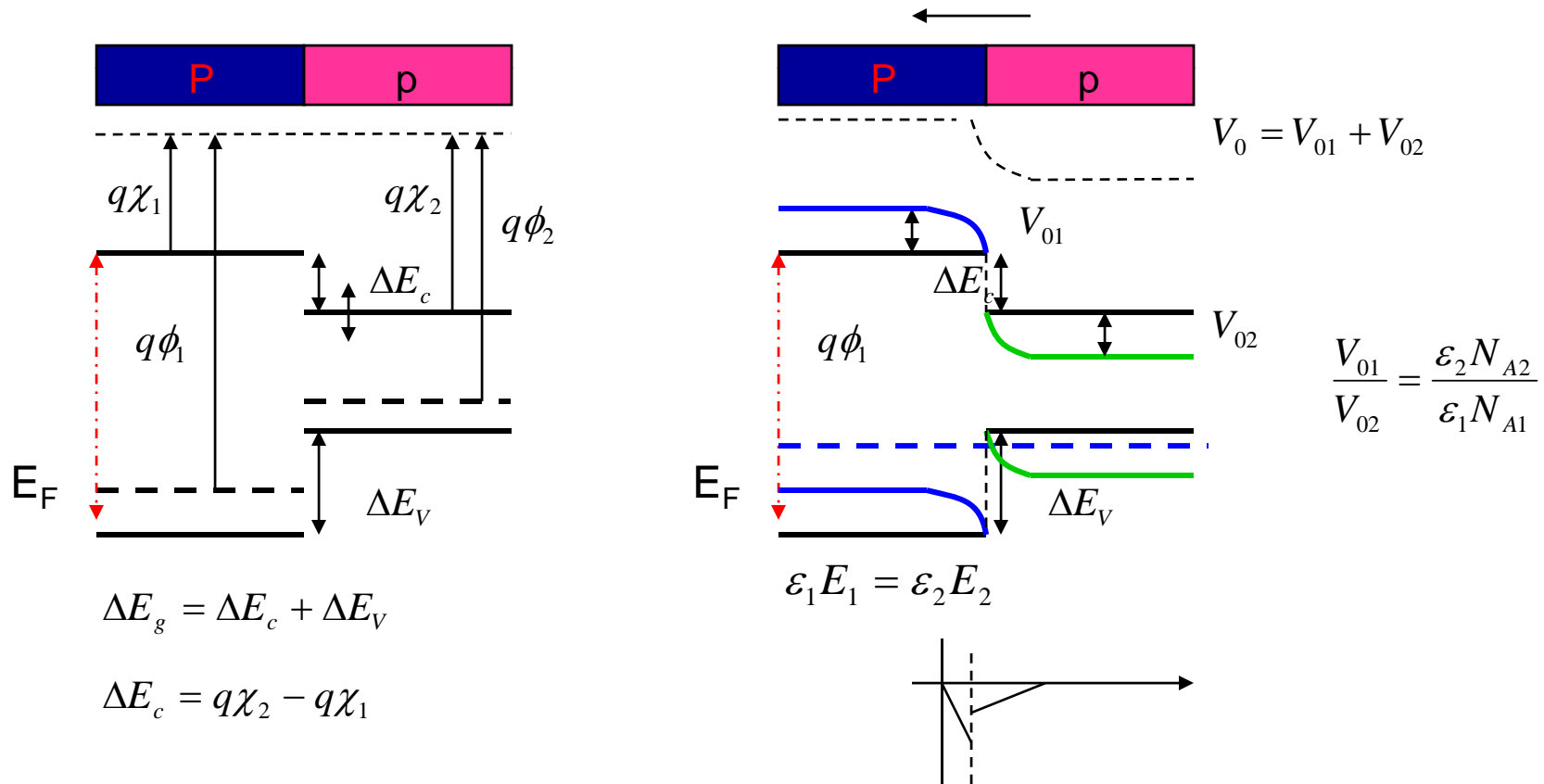
Carrier confinement

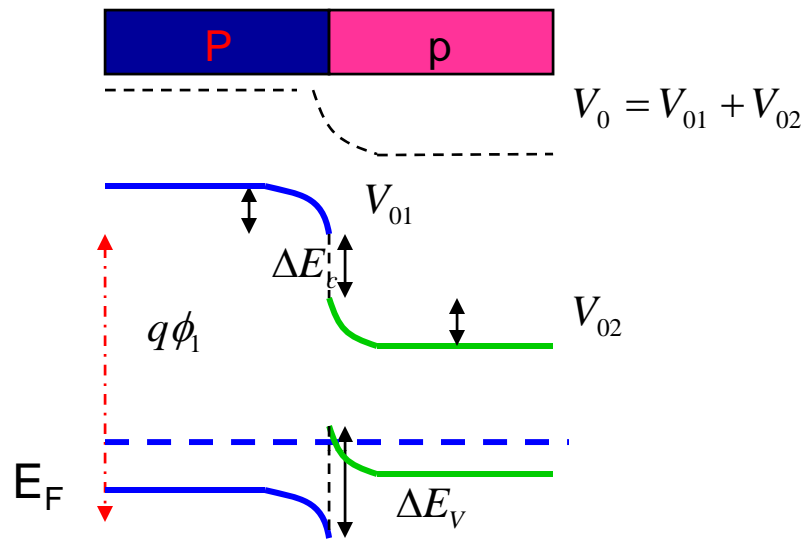
$$W_n \ll L_n = \sqrt{D_n \tau_n} \approx 13 \mu m$$

$$W_p \ll L_p = \sqrt{D_p \tau_p} \approx 3 \mu m$$

Better light emitting efficiency

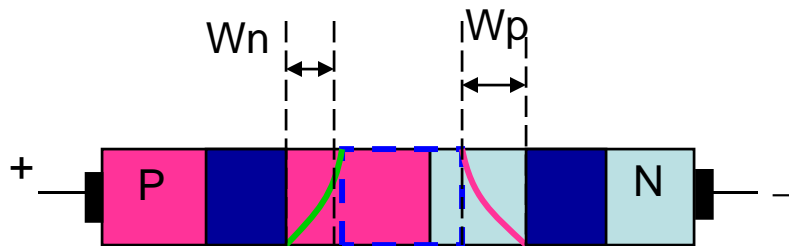
Band diagram

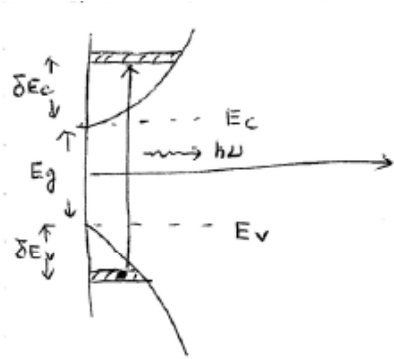




$$n_1 = n_2 e^{-\frac{q(V_0 + \Delta E_c)}{kT}}$$

Carrier confinement:





$$E_2 - E_c = \frac{\hbar^2 k_c^2}{2m_e^*}$$

$$E_v - E_1 = \frac{\hbar^2 k_v^2}{2m_h^*}$$

$k_c = k_v = k$  no-photon transitions.

$$\Rightarrow E_2 - E_c = \frac{m_h^*}{m_e^*} (E_v - E_1)$$

$$h\nu = E_2 - E_1$$

$$\begin{aligned} \Rightarrow E_2 - E_c &= E_2 - E_1 + E_1 - E_v + E_v - E_c \\ &= h\nu + (E_v - E_1) - E_g \\ &= h\nu - E_g - (E_v - E_1) \end{aligned}$$

$$\Rightarrow E_2 - E_c = \frac{1}{1 + \frac{m_e^*}{m_h^*}} (h\nu - E_g) = \frac{m_h^*}{m_e^* + m_h^*} (h\nu - E_g)$$

$$E_v - E_1 = \frac{1}{1 + \frac{m_h^*}{m_e^*}} (h\nu - E_g) = \frac{m_e^*}{m_e^* + m_h^*} (h\nu - E_g)$$

$$\therefore E_2 - E_c = \frac{m_h^*}{m_e^*} (E_v - E_1)$$

$$\Rightarrow dE_2 = -\frac{m_h^*}{m_e^*} dE_1$$

Reduced density of states

$$P(h\nu)dE = P(E_2 - E_1) d(E_2 - E_1) = P(h\nu)(dE_2 - dE_1)$$

$$dE = \hbar^2 \left( \frac{1}{m_e^*} + \frac{1}{m_h^*} \right) k dk$$

$$\Rightarrow \frac{dk}{dE} = \frac{1}{\hbar^2 \left[ \frac{1}{m_e^*} + \frac{1}{m_h^*} \right] k}$$

$$\Rightarrow P(h\nu) = \frac{4\pi k^2}{(2\pi)^3} \cdot \frac{dk}{dE} = \frac{1}{2\pi^2 k^2} \frac{m_e^* m_h^*}{m_e^* + m_h^*} k$$

$$kR = \left[ 2m_e^* (E_0 - E_g) \right]^{1/2}$$

$$\Rightarrow P(h\nu) = \frac{1}{4\pi^2 k^2} \left( \frac{2m_e^*}{k^2} \right)^{3/2} (h\nu - E_g)$$

$$M_r^* = \frac{m_e^* m_h^3}{m_e^* + m_h^*}$$

$P(h\nu)$  ← reduced density of states, and reflects the fact the conditions are placed on the recombination.

## Photon density of State (DOS)

Optical mode volume:

$$k_x = m_x \frac{\pi}{L_x}$$

$$V_{\text{mode}} = \frac{2\pi}{L_x} \frac{2\pi}{L_y} \frac{2\pi}{L_z}$$

Optical mode density from  $k$  to  $k+dk$ , # of optical modes / volume ( $\text{m}^3$ ) /  $dk$

$$\left(\frac{1}{2\pi}\right)^3 4\pi k^2 dk$$

$$E = c\hbar k \quad dk = \frac{E dE}{c\hbar}$$

Optical mode density from  $E$  to  $E+dE$ , #/ of optical modes / volume ( $\text{m}^3$ ) /  $dE$

$$\left(\frac{1}{2\pi}\right)^3 4 \left(\frac{E}{\hbar c}\right)^3 dE$$

## Photon density of State (DOS)

Optical mode density from E to E+dE, # of optical modes / volume (m<sup>3</sup>) /dE

$$\left(\frac{1}{2\pi}\right)^3 4\pi \left(\frac{E}{\hbar c}\right)^3 dE = \left(\frac{1}{2\pi}\right)^3 4\pi \left(\frac{\omega}{c}\right)^3 dE$$

# of photons per optical mode

$$\frac{1}{e^{\hbar\omega/kT} - 1}$$

# of photons from E to E+dE

$$n(E) = \left(\frac{1}{2\pi}\right)^3 4\pi \left(\frac{\omega}{c}\right)^3 \frac{1}{e^{\hbar\omega/kT} - 1} dE$$

$$n(E) = 4\pi \left(\frac{\nu}{c}\right)^3 \frac{h}{e^{h\nu/kT} - 1} d\nu = \frac{4\pi h \nu^3}{c^3} \frac{d\nu}{e^{h\nu/kT} - 1}$$

$$\rho(h\nu) = h\nu \frac{4\pi h \nu^3}{c^3} \frac{1}{e^{h\nu/kT} - 1} d\nu$$

In J/m<sup>3</sup>/Hz: Photon energy density per Hz

$$\rho(h\nu) = h\nu \frac{8\pi h \nu^3}{c^3} \frac{d\nu}{e^{h\nu/kT} - 1}$$

2 polarizations

## Photon density of State (DOS)

Optical mode volume:

$$k_x = m_x \frac{\pi}{L_x}$$

$$V_{\text{mode}} = \frac{2\pi}{L_x} \frac{2\pi}{L_y} \frac{2\pi}{L_z}$$

Optical mode density per dk, #/ of optical modes / volume (m<sup>3</sup>) /dk

$$\left(\frac{1}{2\pi}\right)^3 4\pi k^2 dk$$

$\rho(h\nu)$

In m<sup>3</sup>/J/s/Hz,  
stimulated emission coefficient

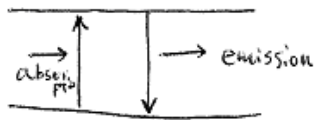
$$E = c\hbar k \quad dk = \frac{EdE}{c\hbar}$$

Optical mode density per dE, #/ of optical modes / volume (m<sup>3</sup>) /dE



## 1. Emission, Absorption &amp; Amplification

... two level system model



Note: this is a simplified model,

$$(2) E_2 - E_1 = h\nu$$

a). absorption rate:

$$R_{ab} = B_{12} N_1 \rho(h\nu)$$

$\rho(h\nu)$

In  $J/m^3/Hz$ : Photon energy density per Hz

b). emission:

① spontaneous

$$R_{sp} = A_{21} N_2$$

① random phase

② random polarization

$R_{ab}$  In  $\#/m^3 s$ , absorption rate per unit volume

② stimulate emission

$$R_{st} = B_{21} N_2 \rho(h\nu)$$

① fixed phase

② fixed polarization.

$A_{21}$  In  $/s$ , spontaneous emission rate

$$N_2 = N_1 \exp - (E_2 - E_1) / kT.$$

$B$

In  $m^3/J/s/Hz$ ,  
stimulated emission coefficient

At equilibrium, absorption = emission rate

$$B_{12} N_1 \rho(h\nu) = A_{21} N_2 + B_{21} N_2 \rho(h\nu) \dots \textcircled{1}$$

in thermal equilibrium,  $N_2 = N_1 e^{- (E_2 - E_1) / kT}$

$$\rho_{eq}(h\nu) = \frac{8\pi h\nu^3}{c^3 [\exp(\frac{h\nu}{kT}) - 1]}$$

← plank's  
black body  
distribution Law

$$\dots \left[ B_{12} N_1 - B_{21} N_2 e^{-(E_2 - E_1)/kT} \right] \cdot \frac{8\pi h\nu^3}{c^3 [e^{h\nu/kT} - 1]} = A_{21} N_2$$

$$\Rightarrow \boxed{B_{12} = B_{21}}, \quad \boxed{A_{21} = B_{21} \cdot \frac{8\pi h\nu^3}{c^3}}$$

$$A_{21} = B_{21} \times \# \text{ optical modes}$$

Einstein coefficients.  $\leftarrow$  valid in any condition.

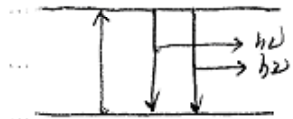
- stimulated emission rate to spontaneous emission rate

$$\dots \frac{R_{st}}{R_{sp}} = \frac{B_{21} N_2 \rho(h\nu)}{A_{21} N_2} = \frac{B_{21}}{A_{21}} \rho(h\nu) \cdot \frac{c^3}{8\pi h\nu^3} = N(h\nu)$$

- stimulated emission rate to absorption

$$\dots \frac{R_{st}}{R_{ab}} = \frac{B_{21} N_2 \rho(h\nu)}{B_{12} N_1 \rho(h\nu)} = \frac{N_2}{N_1}$$

... what stimulated emission rate  $R_{st} > R_{ab}$ .



- $\Rightarrow N_2 > N_1$ , population inversion.
- $N_2 = N_1$ , optical transparency.

Example: Calculate the spontaneous to stimulate emission ratio at frequency  $5 \times 10^{14} \text{ Hz}$  @ 1300 K.

$$\frac{R_{sp}}{R_{st}} = \frac{A_{21} N_2}{B_{21} N_2 \rho(\nu)} = \frac{A_{21}}{B_{21} \rho(\nu)}$$

$$= \frac{\frac{8\pi h\nu^3}{c^3}}{\rho(\nu)}$$

$\rho(\nu)$  at thermal equilibrium

$$\rho(\nu) = \frac{8\pi h\nu^3}{c^3 (e^{\frac{h\nu}{kT}} - 1)}$$

$$\Rightarrow \frac{R_{sp}}{R_{st}} \approx e^{\frac{h\nu}{kT}} = e^{18.4}$$

Time-independent perturbation theory

$$H_0|n^{(0)}\rangle = E_n^{(0)}|n^{(0)}\rangle, \quad n = 1, 2, 3, \dots$$

$$(H_0 + \lambda V)|n\rangle = E_n|n\rangle.$$

$$E_n = E_n^{(0)} + \lambda E_n^{(1)} + \lambda^2 E_n^{(2)} + \dots$$

$$|n\rangle = |n^{(0)}\rangle + \lambda|n^{(1)}\rangle + \lambda^2|n^{(2)}\rangle + \dots$$

$$\begin{aligned} (H_0 + \lambda V)(|n^{(0)}\rangle + \lambda|n^{(1)}\rangle + \dots) \\ = (E_n^{(0)} + \lambda E_n^{(1)} + \lambda^2 E_n^{(2)} + \dots)(|n^{(0)}\rangle + \lambda|n^{(1)}\rangle + \dots) \end{aligned}$$

$$H_0|n^{(1)}\rangle + V|n^{(0)}\rangle = E_n^{(0)}|n^{(1)}\rangle + E_n^{(1)}|n^{(0)}\rangle$$

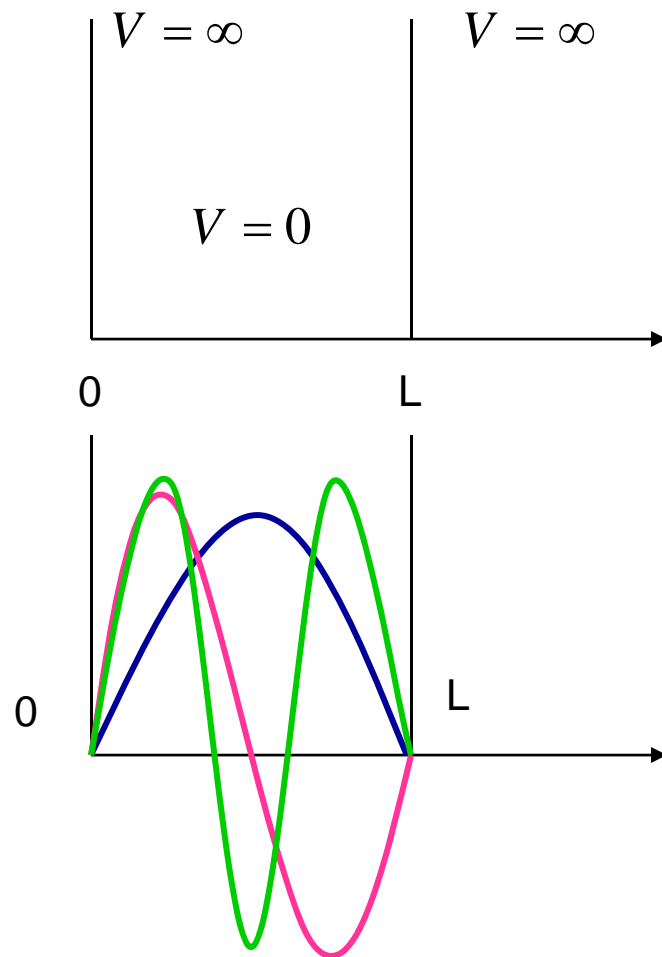
$$E_n^{(1)} = \langle n^{(0)}|V|n^{(0)}\rangle$$

$$(E_n^{(0)} - H_0)|n^{(1)}\rangle = \sum_{k \neq n} |k^{(0)}\rangle \langle k^{(0)}|V|n^{(0)}\rangle$$

$$|n^{(1)}\rangle = \sum_{k \neq n} \frac{\langle k^{(0)}|V|n^{(0)}\rangle}{E_n^{(0)} - E_k^{(0)}} |k^{(0)}\rangle$$

$$E_n = E_n^{(0)} + \langle n^{(0)}|V|n^{(0)}\rangle + \sum_{k \neq n} \frac{|\langle k^{(0)}|V|n^{(0)}\rangle|^2}{E_n^{(0)} - E_k^{(0)}} + \dots$$

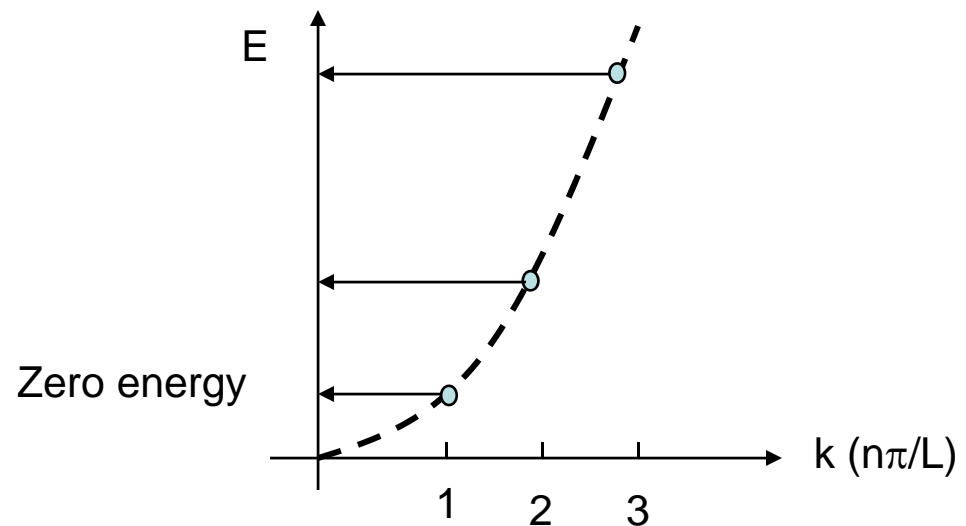
Example



$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) + V(x)\psi(x) = E\psi(x)$$

$$\psi(x) = A \sin kx, \quad \psi(x) = A \sin\left(\frac{n\pi}{L}x\right),$$

$$kL = \frac{\sqrt{2mE}}{\hbar} L = n\pi, \quad n > 0 \quad E = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$



Example

$$E_n^0 = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$

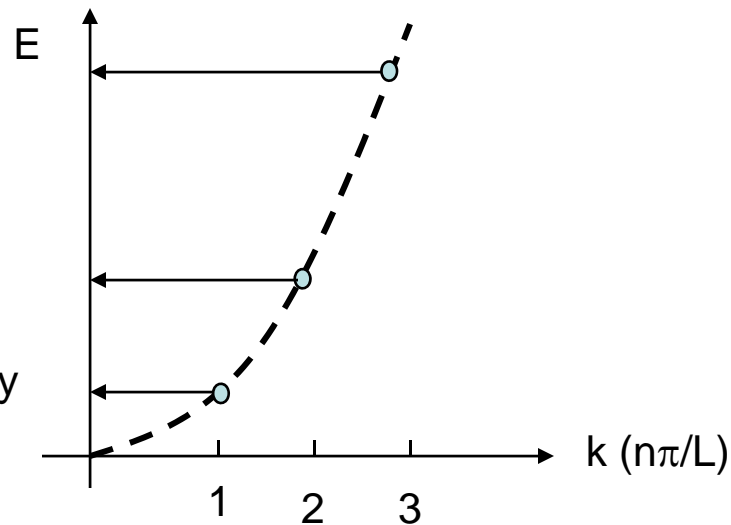
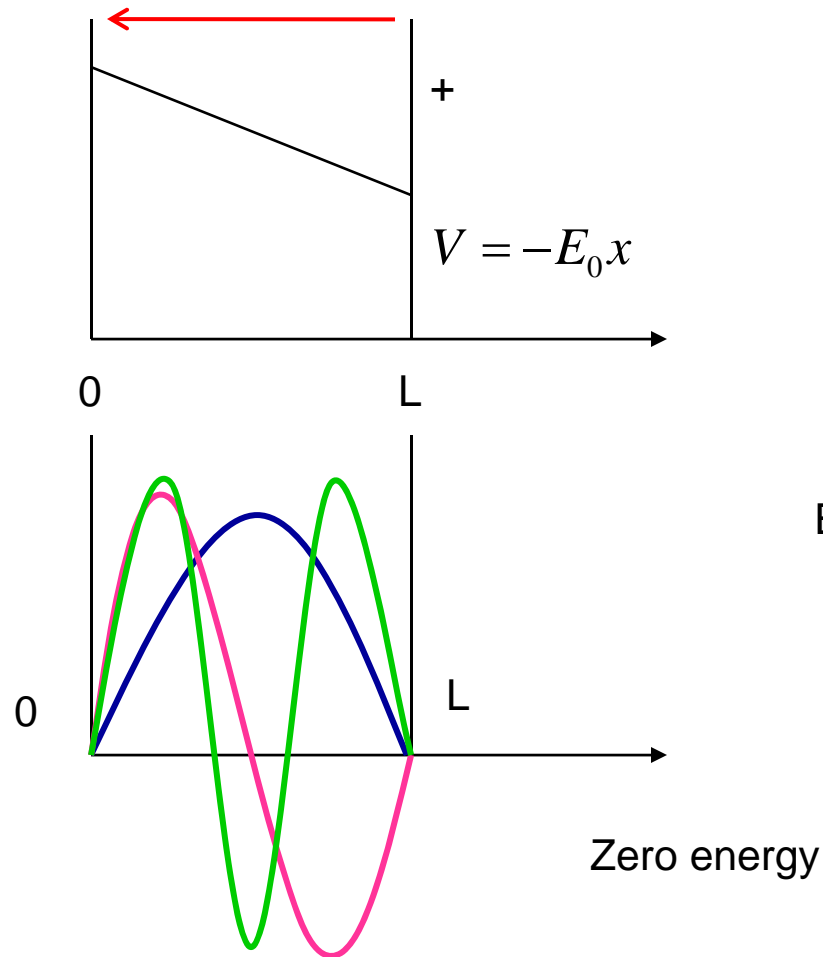
$$E_n^1 = \langle \psi_n^0 | V_0 | \psi_n^0 \rangle$$

$$= \int_0^L V_0 \sin^2 \frac{n\pi x}{L} dx$$

$$= \int_0^L V_0 \frac{1 - \cos \frac{2n\pi x}{L}}{2} dx$$

$$= V_0 \frac{L}{2}$$

Wrong, why?



Time-dependent perturbation theory

$$i\hbar \frac{\partial}{\partial t} \psi = H \psi \quad i\hbar \frac{\partial}{\partial t} \psi_n^0 = H^0 \psi_n^0 \quad i\hbar \frac{\partial}{\partial t} \psi_n^0 = E_n^0 \psi_n^0 \quad \psi_n^0 = C_n^0 e^{-\frac{iE_n t}{\hbar}}$$

Initial state

$$\psi_n^0 = C_n^0 e^{-\frac{iE_n t}{\hbar}} \left| \psi_n^0 \right\rangle$$

$$i\hbar \frac{\partial}{\partial t} \psi = (H^0 + H^1) \psi$$

Final state

$$\psi^1 = \sum_k C_k^1 e^{-\frac{iE_k t}{\hbar}}$$

$$i\hbar \frac{\partial}{\partial t} \sum_k C_k^1 e^{-\frac{iE_k t}{\hbar}} \psi_k^0 = (H^0 + H^1) \sum_k C_k^1 e^{-\frac{iE_k t}{\hbar}} \psi_k^0$$

$$i\hbar \frac{\partial}{\partial t} \sum_k C_k^1 e^{-\frac{iE_k t}{\hbar}} \psi_k^0 = (H^0 + H^1) \sum_k C_k^1 e^{-\frac{iE_k t}{\hbar}} \psi_k^0$$

$$i\hbar \sum_k C_k^1 \frac{\partial}{\partial t} e^{-\frac{iE_k t}{\hbar}} \psi_k^0 + i\hbar \sum_k \dot{C}_k^1 e^{-\frac{iE_k t}{\hbar}} \psi_k^0 = H^0 \sum_k C_k^1 e^{-\frac{iE_k t}{\hbar}} \psi_k^0 + H^1 \sum_k C_k^1 e^{-\frac{iE_k t}{\hbar}} \psi_k^0$$

$$i\hbar \sum_k \dot{C}_k^1 e^{-\frac{iE_k t}{\hbar}} \psi_k^0 = H^1 \sum_k C_k^1 e^{-\frac{iE_k t}{\hbar}} \psi_k^0$$

$$i\hbar \sum_k \dot{C}_k^1 e^{-\frac{iE_k t}{\hbar}} |\psi_k^0\rangle = H^1 \sum_k C_k^1 e^{-\frac{iE_k t}{\hbar}} |\psi_k^0\rangle$$

$$i\hbar \dot{C}_m^1 = \sum_k C_k^1 e^{-\frac{i(E_k - E_m)t}{\hbar}} \langle \psi_m^0 | H^1 | \psi_k^0 \rangle$$

$$\dot{C}_m^1 = -\frac{i}{\hbar} \sum_k C_k^1 e^{-\frac{i(E_k - E_m)t}{\hbar}} \langle \psi_m^0 | H^1 | \psi_k^0 \rangle$$



$$\dot{C}_m^1 = -\frac{i}{\hbar} \sum_k C_k^1 e^{-\frac{i(E_k - E_m)t}{\hbar}} \langle \psi_m^0 | H' | \psi_k^0 \rangle \quad \text{Initial state } |\psi_n^0\rangle$$

$$\dot{C}_m^1 = \frac{1}{i\hbar} e^{-\frac{i(E_k - E_m)t}{\hbar}} \langle \psi_m^0 | H' | \psi_n^0 \rangle$$

Example 1:  $H' = V_0$

$$\begin{aligned} C_m^1 &= -\frac{i}{\hbar} \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \int_0^t e^{-\frac{i(E_k - E_m)t}{\hbar}} dt \\ &= \frac{i}{\hbar} \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \frac{\hbar}{i(E_k - E_m)} e^{-\frac{i(E_k - E_m)t}{\hbar}} \Big|_0^t \\ &= \frac{\langle \psi_m^0 | V_0 | \psi_n^0 \rangle}{(E_k - E_m)} \left( 1 - e^{-\frac{i(E_k - E_m)t}{\hbar}} \right) \\ &= \frac{\langle \psi_m^0 | V_0 | \psi_n^0 \rangle}{(E_k - E_m)} e^{-\frac{i(E_k - E_m)t}{2\hbar}} \left( 2i \sin \frac{(E_k - E_m)t}{2\hbar} \right) \end{aligned}$$

$$C_m^1 = \frac{i}{\hbar} \langle \psi_m^0 | V_0 | \psi_n^0 \rangle e^{-\frac{i(E_k - E_m)t}{2\hbar}} \frac{\sin \frac{(E_k - E_m)t}{2\hbar}}{(E_k - E_m) / 2\hbar}$$

$$C_m^1 = \frac{i\pi}{\hbar} \langle \psi_m^0 | V_0 | \psi_n^0 \rangle e^{-\frac{i(E_k - E_m)t}{2\hbar}} \frac{\sin \frac{(E_k - E_m)t}{2\hbar}}{\pi(E_k - E_m) / 2\hbar}$$

$$|C_m^1|^2 = \frac{\pi^2}{\hbar^2} \left| \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \right|^2 \delta(E_k - E_m) \quad \text{Scattering}$$

$$\delta(x) = \lim_{L \rightarrow \infty} \frac{\sin(Lx)}{\pi x}$$

$$\dot{C}_m^1 = \frac{1}{i\hbar} e^{-\frac{i(E_k - E_m)t}{\hbar}} \langle \psi_m^0 | H' | \psi_n^0 \rangle$$

Example 2:  $H' = V_0 \cos \omega_{op} t$

$$\begin{aligned} C_m^1 &= -\frac{i}{\hbar} \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \int_0^t e^{-\frac{i(E_k - E_m)t}{\hbar}} \cos \omega t dt \\ &= -\frac{i}{\hbar} \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \int_0^t \left[ e^{-\frac{i(E_k - E_m - \hbar\omega)t}{\hbar}} + e^{-\frac{i(E_k - E_m + \hbar\omega)t}{\hbar}} \right] dt \\ &= \frac{i}{\hbar} \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \frac{\hbar}{i(E_k - E_m)} e^{-\frac{i(E_k - E_m - \hbar\omega)t}{\hbar}} \Big|_0^t \\ &= \frac{\langle \psi_m^0 | V_0 | \psi_n^0 \rangle}{(E_k - E_m)} \left( 1 - e^{-\frac{i(E_k - E_m)t}{\hbar}} \right) \\ &= \frac{\langle \psi_m^0 | V_0 | \psi_n^0 \rangle}{(E_k - E_m - \hbar\omega)} e^{-\frac{i(E_k - E_m - \hbar\omega)t}{2\hbar}} \left( 2i \sin \frac{(E_k - E_m - \hbar\omega)t}{2\hbar} \right) \end{aligned}$$

$$C_m^1 = \frac{it}{\hbar} \langle \psi_m^0 | V_0 | \psi_n^0 \rangle e^{-\frac{i(E_k - E_m - \hbar\omega)t}{2\hbar}} \left( \frac{\sin \frac{(E_k - E_m - \hbar\omega)t}{2\hbar}}{(E_k - E_m - \hbar\omega)t / 2\hbar} \right)$$

$$|C_m^1|^2 = \frac{t^2}{\hbar^2} \left| \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \right|^2 \left( \frac{\sin \frac{(E_k - E_m - \hbar\omega)t}{2\hbar}}{(E_k - E_m - \hbar\omega)t / 2\hbar} \right)^2$$

$$|C_{total}^1|^2 = \frac{t^2}{\hbar^2} \left| \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \right|^2 \int \left( \frac{\sin \frac{(\omega_k - \omega_m - \omega)t}{2}}{(\omega_k - \omega_m - \omega)t / 2} \right)^2 \rho(\hbar\omega) d\hbar\omega$$

$$|C_{total}^1|^2 = \frac{t^2}{\hbar^2} \left| \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \right|^2 \frac{2\hbar}{t} \int \left( \frac{\sin \frac{(\omega_k - \omega_m - \omega)t}{2}}{(\omega_k - \omega_m - \omega)t / 2} \right)^2 \rho(\hbar\omega) d(\omega_k - \omega_m - \omega)t / 2$$

$$|C_{total}^1|^2 = \frac{t^2}{\hbar^2} \left| \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \right|^2 \frac{2\hbar\pi}{t}$$

$$|C_{total}^1|^2 = \frac{2\pi t}{\hbar} \left| \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \right|^2$$

$$W = \frac{d}{dt} |C_{total}^1|^2 = \frac{2\pi}{\hbar} \left| \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \right|^2 \rho(\hbar\omega)$$

$$W = \frac{2\pi}{\hbar} \left| \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \right|^2 \rho(\hbar\omega)$$

Fermi-Golden rule

$$W_{nm} = \frac{2\pi}{\hbar} \left| \langle \psi_m^0 | V_0 | \psi_n^0 \rangle \right|^2 \delta(E_m - E_n - \hbar\omega)$$

Fermi-Golden rule