22.581 Module 6: The Unsteady Bernoulli Equation

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Outline



Unsteady Bernoulli Equation

• From Governing Equations

General equations of motion

• Equate change in momentum per unit volume with the forces per unit volume:

$$\rho \vec{a} = \rho \frac{D\vec{u}}{Dt} = -\nabla p + \rho \vec{g} + \frac{\dot{F}_{visc}}{Vol}$$

Expanded in three dimensions:

$$\begin{aligned} x - direction \ : \rho \left(\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right) &= -\frac{\partial p}{\partial x} + \rho g_x + \frac{F_{x_{visc}}}{Vol} \\ y - direction \ : \rho \left(\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} \right) &= -\frac{\partial p}{\partial y} + \rho g_y + \frac{F_{y_{visc}}}{Vol} \\ z - direction \ : \rho \left(\frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} \right) &= -\frac{\partial p}{\partial z} + \rho g_z + \frac{F_{z_{visc}}}{Vol} \end{aligned}$$

• IF $\frac{\vec{F}_{visc}}{Vol} = 0$, then the above equations are the Euler Equations for an incompressible fluid.

Assume that the flow is:

- Inviscid
- Incompressible

$$\rho \vec{a} = \rho \frac{D \vec{u}}{D t} = -\nabla p + \rho \vec{g}$$

$$\frac{\partial \vec{u}}{\partial t} + (\vec{u} \cdot \nabla) \, \vec{u} = \frac{\partial \vec{u}}{\partial t} + \nabla \left(\frac{|u|^2}{2}\right) - \vec{u} \times (\nabla \times \vec{u})$$

$$\rho \left[\frac{\partial \vec{u}}{\partial t} + \nabla \left(\frac{|u|^2}{2} \right) - \vec{u} \times (\nabla \times \vec{u}) \right] = -\nabla p + \rho \vec{g}$$

• We then integrate from point 1 to 2 along a streamline:

$$\int_{1}^{2} \left[\frac{\partial \vec{u}}{\partial t} + \nabla \left(\frac{u^{2}}{2} \right) - \vec{u} \times (\nabla \times \vec{u}) = -\frac{\nabla p}{\rho} - \nabla (gz) \right] d\vec{s}$$

• Here $d\vec{s}$ is a differential length of the streamline.

$$\int_{1}^{2} \left[\frac{\partial \vec{u}}{\partial t} + \nabla \left(\frac{u^{2}}{2} \right) - \vec{u} \times (\nabla \times \vec{u}) = -\frac{\nabla p}{\rho} - \nabla (gz) \right] d\vec{s}$$

• Let's expand and rearrange terms a little:

$$\int_{1}^{2} \vec{u} \times (\nabla \times \vec{u}) \cdot d\vec{s} = \int_{1}^{2} \frac{\partial \vec{u}}{\partial t} \cdot d\vec{s} + \int_{1}^{2} \underbrace{\nabla \left(\frac{u^{2}}{2}\right) \cdot d\vec{s}}_{d\left(\frac{v^{2}}{2}\right)} \\ + \int_{1}^{2} \underbrace{\frac{\nabla p}{\rho} \cdot d\vec{s}}_{\frac{dp}{\rho}} + \int_{1}^{2} \underbrace{\nabla gz \cdot d\vec{s}}_{d(gz)}$$

• Which simplifies to:

$$\int_{1}^{2} \underbrace{\vec{u} \times (\nabla \times \vec{u}) \cdot d\vec{s}}_{see \ comments \ below} = \int_{1}^{2} \frac{\partial \vec{u}}{\partial t} \cdot d\vec{s} + \int_{1}^{2} \frac{dp}{\rho} + \left[\frac{\vec{u}_{2}^{2}}{2} - \frac{\vec{u}_{1}^{2}}{2}\right] + \left[gz_{2} - gz_{1}\right]$$

• Let's examine the first term of the above equation. The integrand is zero along the streamline (the dot product of a vector with a orthonormal vector is zero, by definition):

$$\int_{1}^{2} \vec{u} \times (\nabla \times \vec{u}) \cdot d\vec{s} = 0 \quad along \ streamline$$

• If the flow is *irrotational* everywhere (ie. $\nabla \times \vec{u} = 0$) then the above restriction is not just along a streamline anymore.

• If we also set $\rho = const$. then:

$$\int_{1}^{2} \frac{\partial \vec{u}}{\partial t} \cdot d\vec{s} + \frac{p_2}{\rho} + \frac{u_2^2}{2} + gz_2 = \frac{p_1}{\rho} + \frac{u_1^2}{2} + gz_1$$

True...

between any two points along a streamline

2 everywhere if $\nabla \times \vec{u} = 0$

• If $\nabla \times \vec{u} = 0$ we can also write $\nabla \phi = \vec{u}$, and then $\int_{1}^{2} \frac{\partial \vec{u}}{\partial t} \cdot d\vec{s} = \frac{\partial \phi}{\partial t}$

- The assumptions are:
 - Flow is inviscid
 - Plow is incompressible