

Calculus of Variations

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Calculus of Variations

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Background

Definition

- A function is a mapping of single values to single values.
- A functional is a mapping of function values to single or function values. It usually contains single or multiple variables and their derivatives.
- *Dirichlet Principle*: There exists one stationary ground state for energy.
- Euler's Equation defines the condition for finding the extrema of functionals. → An extremal is the maximum or minimum integral curves of Euler's equation of a functional.
- Calculus of Functionals: Determining the properties of functionals.
- Calculus of Variations: Finding the extremals of functionals.

Background

Single value calculus:

 Functions take extreme values on bounded domain. Necessary condition for extremum at x₀, if f is differentiable:

$$f'(x_0) = 0$$

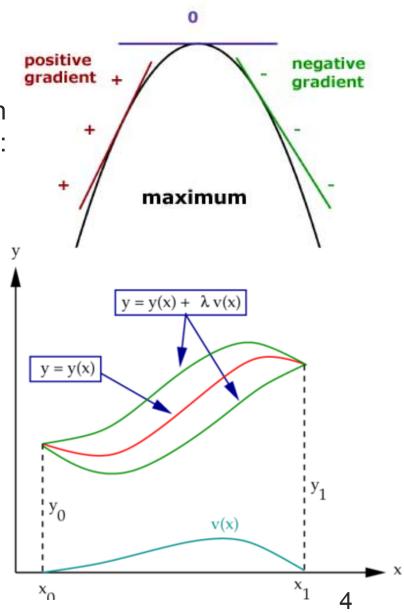
Calculus of variations

 Test function v(x), which vanishes at endpoints, used to find extremal:

$$w(x) = u(x) + \varepsilon v(x)$$
 $I[\varepsilon] = \int_a^b F(x, w, w_x) dx$

Necessary condition for extremal:

$$\frac{dI}{d\varepsilon} = 0$$



Maximum and Minimum of Functions

Maximum and minimum

(a) If f(x) is twice continuously differentiable on $[x_0, x_1]$ i.e.

Nec. condition for a max. (min.) of f(x) at $x \in [x_0, x_1]$ is that F'(x) = 0

Suff. condition for a max (min.) of f(x) at $x \in [x_0, x_1]$ are that F'(x) = 0 and $F''(x) \le 0$

(b) If f(x) over closed domain D. Then nec. and suff. condition for a max. (min.)

of f(x) at
$$x_0 \in D - \partial D$$
 are that $\frac{\partial f}{\partial x_i}\Big|_{x=x_0} = 0$ i = 1,2...n and also that $\frac{\partial^2 f}{\partial x_i \partial x_j}\Big|_{x=x_0}$ is a negative infinite .

Maximum and Minimum of Functions

(c) If f(x) on closed domain DIf we want to extremize f(x) subject to the constraints $g_i(x_1, K, x_n) = 0$ i=1,2,...k (k < n)

EX: Find the extrema of f(x,y) subject to g(x,y) = 0

i) Approach One: Direct differentiation of g(x, y)

$$dg = g_x dx + g_y dy = 0$$

$$\Rightarrow dy = -\frac{g_x}{g_y} dx$$

To extremize *f*

$$df = f_x dx + f_y dy = 0$$

$$(f_x - f_y \frac{g_x}{g_y}) dx = 0$$

Maximum and Minimum of Functions

We have

$$f_x g_y - f_y g_x = 0$$
 and $g = 0$

to find (x_0, y_0) which is to extremize f subject to g = 0

ii) Approach Two: Lagrange Multiplier

Let
$$v(x, y, \lambda) = f(x, y) + \lambda g(x, y)$$

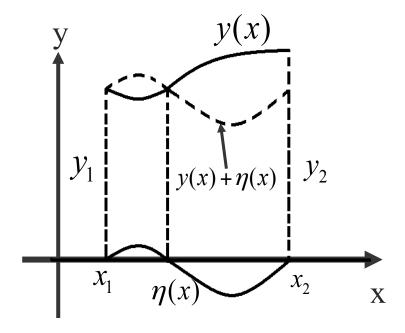
 \Rightarrow extrema of v without any constraint $\langle --- \rangle$ extrema of f subject to g = 0

To extremize
$$v \Rightarrow \begin{cases} \frac{\partial v}{\partial x} = f_x + \lambda g_x = 0 \\ \frac{\partial v}{\partial y} = f_y + \lambda g_y = 0 \end{cases} \Rightarrow \begin{cases} f_x g_y - f_y g_x = 0 \\ \frac{\partial v}{\partial \lambda} = g = 0 \end{cases}$$

We obtain the same equations by extremizing v. where λ is called the Lagrange Multiplier.

Maximum and Minimum of Functionals

Functionals are function's function.



• The basic problem in calculus of variations.

Determine $y(x) \in c^2[x_0, x_1]$ such that the functional :

$$I(y(x)) = \int_{x_0}^{x_1} F(x, y(x), y'(x)) dx$$
 as an extrema

where $F \in \mathbb{C}^2$ over its entire domain, subject to $y(x_0) = y_0$, $y(x_1) = y_1$ at the end points.

Maximum and Minimum of Functionals

Using integrating by parts of the 2nd term, it leads to

$$\Rightarrow [F_{y'}(x,y,y')\eta]_{x_0}^{x_1} - \int_{x_0}^{x_1} \left[\frac{d}{dx}F_{y'}(x,y,y') - F_y(x,y,y')\right] \eta dx = 0 \qquad -----(1)$$

Since $\eta(x_0) = \eta(x_1) = 0$ and $\eta(x)$ is arbitrary,

$$\Rightarrow \frac{d}{dx} [F_{y'}(x, y, y')] - F_y(x, y, y') = 0$$
 ------(2) (Euler's Equation)

Natural B.C's

$$\left[\frac{\partial F}{\partial y'}\right]_{x_0} = 0 \quad \text{or/and} \quad \left[\frac{\partial F}{\partial y'}\right]_{x=x_1} = 0$$

The above requirements are called national b.c's.

The Variational Notation

Variations

Imbed u(x) in a "parameter family" of function $\phi(x, \varepsilon) = u(x) + \varepsilon \eta(x)$ the variation of u(x) is defined as

$$\delta u = \varepsilon \eta(x)$$

The corresponding variation of F, δF to the order in ε is ,

since
$$\delta F = F(x + y + \varepsilon \eta, y' + \varepsilon \eta') - F(x, y, y')$$

$$= \frac{\partial F}{\partial y} \delta y + \frac{\partial F}{\partial y'} \delta y'$$
and
$$I(u + \varepsilon \eta) = \int_{x_0}^{x_1} F(x, u + \varepsilon \eta, u' + \varepsilon \eta') dx = G(\varepsilon)$$
Then
$$\delta I = \delta \int_{x_0}^{x_1} F(x, y, y') dx$$

$$= \int_{x_0}^{x_1} \delta F(x, y, y') dx$$

$$= \int_{x_0}^{x_1} \left(\frac{\partial F}{\partial y} \delta y + \frac{\partial F}{\partial y'} \delta y' \right) dx$$

The Variational Notation

$$= \int_{x_0}^{x_1} \left[\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) \right] \delta y dx + \left[\frac{\partial F}{\partial y'} \delta y \right]_{x_0}^{x_1}$$

Thus, a stationary function for a functional is one for which the first variation.

$$\frac{\partial F}{\partial y} = 0$$

For more general cases

(a) Several dependent variables

EX:
$$I = \int_{x_0}^{x_1} F(x, y, z; y', z') dx$$
Euler's Eq. $\Longrightarrow \frac{\partial F}{\partial y} - \frac{d}{dx} (\frac{\partial F}{\partial y'}) = 0$, $\frac{\partial F}{\partial z} - \frac{d}{dx} (\frac{\partial F}{\partial z'}) = 0$

(b) Several Independent variables

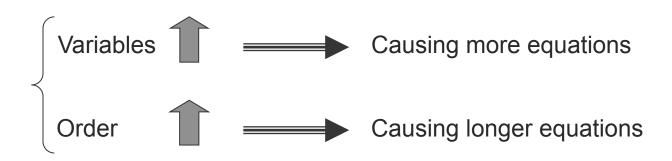
EX:
$$I = \iint_{R} F(x, y, u, u_{x}, u_{y}) dx dy$$
 Euler's Eq. $\Rightarrow \frac{\partial F}{\partial u} - \frac{\partial}{\partial x} (\frac{\partial F}{\partial u_{x}}) - \frac{\partial}{\partial y} (\frac{\partial F}{\partial u_{y}}) = 0$

The Variational Notation

(c) High Orders

EX:
$$I = \int_{x_0}^{x_1} F(x, y, y', y'') dx$$

Euler's Eq.
$$\Longrightarrow \frac{\partial F}{\partial y} - \frac{d}{dx} (\frac{\partial F}{\partial y'}) + \frac{d^2}{dx^2} (\frac{\partial F}{\partial y''}) = 0$$



Lagrange multiplier

Lagrange multiplier can be used to find the extreme value of a multivariate function f subjected to the constraints.

EX:

(a) Find the extreme value of $I = \int_{x_0}^{x_1} F(x, u, v, u_x, v_x) dx$

where
$$u(x_1) = u_1$$
 $u(x_2) = u_2$
 $v(x_1) = v_1$ $v(x_2) = v_2$

and subject to the constraints

$$G(x,u,v) = 0 \qquad ----(3)$$

From
$$\delta I = \int_{x_1}^{x_2} \left\{ \left[\frac{\partial F}{\partial v} - \frac{d}{dx} \left(\frac{\partial F}{\partial u_x} \right) \right] \delta u + \left[\frac{\partial F}{\partial v} - \frac{d}{dx} \left(\frac{\partial F}{\partial v_x} \right) \right] \partial v \right\} dx = 0$$
 ----(4)

Because of the constraints, we don't get two Euler's equations.

From

$$\delta G = \frac{\partial G}{\partial u} \delta u + \frac{\partial G}{\partial v} \delta v = 0 \implies -\frac{G_v}{G_u} \delta v = \delta u$$

Therefore, Eq. (4) becomes

$$\Rightarrow \delta I = \int_{x_0}^{x_1} \left\{ -\frac{G_v}{G_u} \left[\frac{\partial F}{\partial u} - \frac{d}{dx} \left(\frac{\partial F}{\partial u_x} \right) \right] + \left[\frac{\partial F}{\partial v} - \frac{d}{dx} \left(\frac{\partial F}{\partial v_x} \right) \right] \right\} \delta v dx = 0 \qquad -----(5)$$

$$\Rightarrow \frac{\partial G}{\partial v} \left[\frac{\partial F}{\partial u} - \frac{d}{dx} \left(\frac{\partial F}{\partial u_x} \right) \right] - \frac{\partial G}{\partial u} \left[\frac{\partial F}{\partial v} - \frac{d}{dx} \left(\frac{\partial F}{\partial v_x} \right) \right] = 0 \qquad -----(6)$$

The above equations, together with Eq. (3), are used to solve for u, v.

(b) Simple Isoparametric Problem

To extremize $I = \int_{x_1}^{x_2} F(x, y, y') dx$, subject to the constraint:

- i) $J = \int_{x_1}^{x_2} G(x, y, y') dx = const.$
- ii) $y(x_1) = y_1, y(x_2) = y_2$

Take the variation of two-parameter family : $y + \delta y = y + \varepsilon_1 \eta_1(x) + \varepsilon_2 \eta_2(x)$ (where $\eta_1(x)$ and $\eta_2(x)$ are some equations which satisfy

$$\eta_1(x_1) = \eta_2(x_1) = \eta_1(x_2) = \eta_2(x_2) = 0$$

Then,
$$I(\varepsilon_1, \varepsilon_2) = \int_{x_1}^{x_2} F(x, y + \varepsilon_1 \eta_1 + \varepsilon_2 \eta_2, y' + \varepsilon_1 \eta_1' + \varepsilon_2 \eta_2') dx$$

$$J(\varepsilon_1, \varepsilon_2) = \int_{x_1}^{x_2} G(x, y + \varepsilon_1 \eta_1 + \varepsilon_2 \eta_2, y' + \varepsilon_1 \eta_1' + \varepsilon_2 \eta_2') dx$$

To base on the Lagrange Multiplier Method, we can get:

$$\frac{\partial}{\partial \varepsilon_{1}} (I + \lambda J) \Big|_{\varepsilon_{1} = \varepsilon_{2} = 0} = 0$$

$$\frac{\partial}{\partial \varepsilon_{2}} (I + \lambda J) \Big|_{\varepsilon_{1} = \varepsilon_{2} = 0} = 0$$

$$\Rightarrow \int_{x_{1}}^{x_{2}} \left\{ \left[\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) \right] + \lambda \left[\frac{\partial G}{\partial y} - \frac{d}{dx} \left(\frac{\partial G}{\partial y'} \right) \right] \right\} \eta_{i} dx = 0 \qquad i = 1, 2$$

The Euler's equation becomes

$$\frac{\partial}{\partial y} \big(F + \lambda G \big) - \frac{d}{dx} \bigg[\frac{\partial}{\partial y'} \big(F + \lambda G \big) \bigg] = 0$$
 when
$$\frac{\partial G}{\partial y} - \frac{d}{dx} \bigg(\frac{\partial G}{\partial y'} \bigg) = 0 \quad , \; \lambda \; \text{ is arbitrary numbers.}$$

 \Longrightarrow The constraint is trivial, we can ignore λ .

Helmholtz Equation

EX: Force vibration of a membrane

$$\frac{\partial^2 u}{\partial t^2} - c^2 \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) = f(x, y, t) \tag{7}$$

if the forcing function f is of the form

$$f(x, y, t) = P(x, y)\sin(\omega t + \alpha)$$

we may write the steady state displacement *u* in the form

$$u = v(x, y)\sin(\omega t + \alpha)$$

$$\Rightarrow c^{2}(\frac{\partial^{2} v}{\partial x^{2}} + \frac{\partial^{2} v}{\partial v^{2}}) + \omega^{2} v + p = 0$$
-----(8)

$$\int_{R} \left[c^{2} \left(\frac{\partial^{2} v}{\partial x^{2}} + \frac{\partial^{2} v}{\partial y^{2}} \right) + \omega^{2} v + p \right] \delta v dx dy = 0$$

Consider

$$c^{2} \int_{R} v_{xx} \delta v dx dy$$

$$= c^{2} \int_{R} [(v_{x} \delta v)_{x} - v_{x} \delta v_{x}] dx dy$$

Note that $(v_x \delta v)_x = v_{xx} \delta v + v_x \delta v_x$

$$V = V_{x}i + V_{y}j$$

$$V_{x} = v_{x}\delta v$$

$$V_{y} = v_{y}\delta v$$

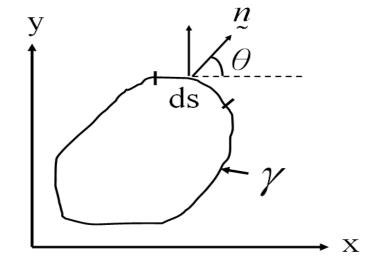
$$M = \cos\theta i + \sin\theta j$$

$$\nabla \bullet V = \frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} = \frac{\partial (v_x \delta v)}{\partial x} + \frac{\partial (v_y \delta v)}{\partial y}$$

$$c^{2} \int_{R} v_{yy} \delta v dx dy$$

$$= c^{2} \int_{R} [(v_{y} \delta v)_{y} - v_{y} \delta v_{y}] dx dy$$

$$(v_y \delta v)_y = v_{yy} \delta v + v_y \delta v_y$$



$$\int_{\Re} (\nabla \cdot V) da = \oint_{\gamma} V \cdot n \, ds = \oint_{\gamma} (v_x \, \delta v \cos \theta + v_y \, \delta v \sin \theta) \, ds$$

$$c^2 \int_{R} v_{xx} \, \delta v \, dx \, dy + \int_{R} c^2 v_{yy} \, \delta v \, dx \, dy$$

$$= c^2 \oint_{\gamma} v_x \, \delta v \cos \theta \, ds - \int_{R} \frac{1}{2} c^2 \delta(v_x)^2 \, dx \, dy$$

$$+ c^2 \oint_{\gamma} v_y \, \delta v \sin \theta \, ds - \int_{R} \frac{1}{2} c^2 \delta(v_y)^2 \, dx \, dy$$

$$\Rightarrow \oint_{\gamma} c^2 (v_x \cos \theta + v_y \sin \theta) \, \delta v \, ds - \int_{R} \frac{1}{2} c^2 \delta[(v_x)^2 + (v_y)^2] \, dx \, dy$$

$$+ \int_{R} \frac{1}{2} \omega_2 \delta(v^2) \, dx \, dy + \int_{R} P \delta v \, dx \, dy = 0$$

$$\Rightarrow \int_{\gamma} c^2 \frac{\partial v}{\partial n} \delta v \, ds - \delta \int_{R} \left[\frac{1}{2} c^2 (\nabla v)^2 - \frac{1}{2} \omega^2 v^2 - P v \right] \, dx \, dy = 0$$

Hence,

i) if
$$v = f(x, y)$$
 is given on γ i.e. $\delta v = 0$ on γ

then the variational problem

$$\Rightarrow \delta \int_{R} \left[\frac{c^{2}}{2} (\nabla v)^{2} - \frac{1}{2} \omega^{2} v^{2} - p v \right] dx dy = 0$$
 -----(10)

ii) if
$$\frac{\partial v}{\partial n} = 0$$
 is given on γ

the variation problem is same as Eq. (10)

iii) if
$$\frac{\partial v}{\partial n} = \psi(s)$$
 is given on γ

$$\Rightarrow \delta \left[\int_{R} \left\{ \frac{1}{2} c^{2} (\nabla v)^{2} - \frac{1}{2} \omega^{2} v^{2} - p v \right\} dx dy - \int_{\tau} c^{2} \psi v dx \right] = 0$$
 ----(11)

Diffusion Equation

EX: Steady State Heat Condition

$$\nabla \cdot (k\nabla T) = f(x,T) \text{ in } D$$

B.C's:

$$T = T_1 \quad \text{on} \quad B_1$$

$$-kn \cdot \nabla T = q_2 \quad \text{on} \quad B_2$$

$$-kn \cdot \nabla T = h(T - T_3) \quad \text{on} \quad B_3$$

Multiply the equation by δT , and integrate over the domain D. After integrating by parts, we find the variational problem as follow.

$$\delta \left[\int_{D} \left\{ \frac{1}{2} k (\nabla T)^{2} + \int_{T_{0}}^{T} f(x, T') dT' \right\} d\tau + \int_{B_{2}} q_{2} T d\sigma + \frac{1}{2} \int_{B_{3}} h (T - T_{3})^{2} d\sigma \right] = 0$$

with $T = T_I$ on B_{I} .

Poisson's Equation

EX: Torsion of a Prismatic Bar

$$\nabla^2 \psi = -2 \qquad \text{in} \quad \mathbb{R}$$

$$\psi = 0 \qquad \text{on} \quad \gamma$$

where ψ is the Prandtl stress function and

$$\sigma_{\tau z} = G_{\alpha} \frac{\partial \psi}{\partial y}$$
, $\sigma_{zy} = \sigma_{\alpha} \frac{\partial \psi}{\partial x}$

The variation problem becomes

$$\delta \left\{ \int_{R} \left[(D\psi)^{2} - 4\psi \right] dx dy \right\} = 0$$

with $\psi = 0$ on γ .

I) Method of Weighted Residuals (MWR)

$$L[u] = 0$$
 in D

with homogeneous b.c's in B. Assume an approximate solution.

$$u = u_n = \sum_{i=1}^n C_i \phi_i$$

where each trial function ϕ_i satisfies the b.c's. The residual is

$$R_n = L[u_n]$$

In this method (MWR), C_i are chosen such that R_n is forced to be zero in an average sense.

i.e.
$$< w_j$$
, $R_n > = 0$, $j = 1,2,...,n$

where w_i are the weighting functions..

II) Galerkin Method

 w_j are chosen to be the trial functions ϕ_j hence the trial functions is chosen as members of a complete set of functions.

Galerkin method force the residual to be zero w.r.t. an orthogonal complete set.

EX: Torsion of a Square Shaft

$$\nabla^2 \psi = -2$$

$$\psi = 0 \quad on \quad x = \pm a, \quad y = \pm a$$

i) One – term approximation

$$\psi_1 = c_1(x^2 - a^2)(y^2 - a^2)$$

$$R_i = \nabla^2 \psi_1 + 2 = 2c_1[(x - a)^2 + (y - a)^2] + 2$$

$$\phi_1 = (x^2 - a^2)(y^2 - a^2)$$

From
$$\int_{-a}^{a} \int_{-a}^{a} R_1 \phi_1 dx dy = 0$$

$$\Rightarrow c_1 = \frac{5}{8} \frac{1}{a^2}$$

Therefore,

$$\psi_1 = \frac{5}{8a^2} (x^2 - a^2)(y^2 - a^2)$$

The torsional rigidity is determined by

$$D_1 = 2G \int_R \psi dx dy = 0.1388G(2a)^4$$

The exact value of D is

$$D_a = 0.1406G(2a)^4$$

The approximation error is -1.2%.

ii) Two – term approximation

$$\psi_{2} = (x^{2} - a^{2})(y^{2} - a^{2})[c_{1} + c_{2}(x^{2} + y^{2})]$$

$$\Longrightarrow R_{2} = \nabla \psi_{2} + 2$$

$$\phi_{1} = (x^{2} - a^{2})(y^{2} - a^{2})$$

$$= (x^{2} - a^{2})(y^{2} - a^{2})$$

$$\Longrightarrow R_{2}\phi_{1}dxdy = 0$$

$$\phi_{2} = (x^{2} - a^{2})(y^{2} - a^{2})(x^{2} + y^{2})$$
and
$$\int_{R} R_{2}\phi_{2}dxdy = 0$$

We obtain
$$c_1 = \frac{1295}{2216} \frac{1}{a^2}$$
, $c_2 = \frac{525}{4432} \frac{1}{a^2}$

Therefore

$$D_2 = 2G \int_R \psi_2 dx dy = 0.1404 G(2a)^4$$

 \rightarrow The error is -0.14%.

I) Kantorovich Method [Kantorovich (1948)]

Assuming the approximate solution as : $u = \sum_{i=1}^{n} C_i(x_n)U_i$

where U_i is a known function decided by b.c. condition.

 C_i is a unknown function decided by minimal "I". \Longrightarrow Euler Equation of C_i

EX: The Torsional Problem with a Functional "I".

$$I(u) = \int_{-a}^{a} \int_{-b}^{b} \left[\left(\frac{\partial u}{\partial x} \right)^{2} + \left(\frac{\partial u}{\partial y} \right)^{2} - 4u \right] dx dy$$

Assuming the one-term approximate solution as:

$$u(x, y) = (b^2 - y^2)C(x)$$

Then,

$$I(C) = \int_{-a}^{a} \int_{-b}^{b} \{(b^2 - y^2)^2 [C'(x)]^2 + 4y^2 C^2(x) - 4(b^2 - y^2) C(x)\} dxdy$$

Integrate by *y*

$$I(C) = \int_{-a}^{a} \left[\frac{16}{15} b^{5} C'^{2} + \frac{8}{3} b^{3} C^{2} - \frac{16}{3} b^{3} C \right] dx$$

Euler's equation is

$$C''(x) - \frac{5}{2b^2}C(x) = -\frac{5}{2b^2}$$
 where b.c. condition is $C(\pm a) = 0$

General solution is

$$C(x) = A_1 \cosh(\sqrt{\frac{5}{2}} \frac{x}{b}) + A_2 \sinh(\sqrt{\frac{5}{2}} \frac{x}{b}) + 1$$

where
$$A_1 = -\frac{1}{\cosh(\sqrt{\frac{5}{2}}\frac{a}{b})}$$
, $A_2 = 0$

and
$$C(x) = \left\{ 1 - \frac{\cosh(\sqrt{\frac{5}{2}} \frac{x}{b})}{\cosh(\sqrt{\frac{5}{2}} \frac{a}{b})} \right\}$$

Therefore, the one-term approximate solution is

$$\mathbf{u} = \left\{ 1 - \frac{\cosh(\sqrt{\frac{5}{2}} \frac{x}{b})}{\cosh(\sqrt{\frac{5}{2}} \frac{a}{b})} \right\} (b^2 - y^2)$$

II) Rayleigh-Ritz Method

This is used when the exact solution is impossible or difficult to obtain.

First, we assume the approximate solution as : $u = \sum_{i=1}^{n} C_i U_i$

where U_i are some approximate function which satisfy the b.c's. Then, we can calculate extreme I.

$$I = I(c_1, K, c_n)$$
 Choose $c_1 \sim c_n$ i.e. $\frac{\partial I}{\partial c_1} = K = \frac{\partial I}{\partial c_n} = 0$

EX:
$$y'' + xy = -x$$
 $y(0) = y(1) = 0$

Its solution can be obtained from

$$\int_0^1 (y'' + xy + x) \delta y dx = 0 \implies I = \int_0^1 \left[\frac{1}{2} (y')^2 - \frac{1}{2} xy^2 - xy \right] dx$$

Assuming that

$$y = x(1-x)(c_1 + c_2x + c_3x^2 K)$$

i) One-term approximation

$$y = c_1 x (1 - x) = c_1 (x - x^2) \quad y' = c_1 (1 - 2x)$$
Then, $I(c_1) = \int_0^1 \left[\frac{1}{2} c_1^2 (1 - 4x + 4x^2) - \frac{x}{2} c_1^2 (x^2 - 2x^3 + x^4) - c_1 x (x - x^2) \right] dx$

$$= \frac{c_1^2}{2} \left(1 - 2 + \frac{4}{3} \right) - \frac{c_1^2}{2} \left(\frac{1}{4} - \frac{2}{5} + \frac{1}{6} \right) - c_1 \left(\frac{1}{3} - \frac{1}{4} \right) = \frac{19}{120} c_1^2 - \frac{c_1}{12}$$

$$\frac{\partial I}{\partial c_1} = 0 \implies \frac{19}{60} c_1 - \frac{1}{12} = 0 \implies c_1 = 0.263 \implies y(1) = 0.263 x (1 - x)$$

ii) Two-term approximation

$$y = x(1-x)(c_1 + c_2x) = c_1(x-x^2) + c_2(x^2 - x^3)$$

Then
$$y' = c_1(1-2x) + c_2(2x-3x^2)$$

$$I(c_1, c_2) = \int_0^1 \left[\frac{1}{2} \left\{ c_1^2 \left(1 - 4x + 4x^2 \right) + 2c_1c_2 \left(2x - 7x^2 + 6x^3 \right) + c_3^2 \left(4x^2 - 12x^3 + 9x^4 \right) \right\} - \frac{1}{2} \left\{ c_1^2 \left(x^3 - 2x^4 + x^5 \right) + 2c_1c_2 \left(x^4 - 2x^5 + x^6 \right) + c_2^2 \left(x^5 - 2x^6 + x^7 \right) \right\} - \left\{ c_1 \left(x^2 - x^3 \right) + c_2 \left(x^3 - x^4 \right) \right\} dx$$

$$= \frac{c_1^2}{2} \left(1 - 2 + \frac{4}{3} - \frac{1}{4} + \frac{2}{5} - \frac{1}{6} \right) + c_1c_2 \left(1 - \frac{7}{3} + \frac{3}{2} - \frac{1}{5} + \frac{1}{3} - \frac{1}{7} \right)$$

$$- \frac{c_1^2}{2} \left(\frac{4}{3} - 3 + \frac{9}{5} - \frac{1}{6} + \frac{2}{7} - \frac{9}{8} \right) - \frac{c_1}{12} - \frac{c_2}{20}$$

$$= \frac{19}{120} c_1^2 + \frac{11}{70} c_1c_2 + \frac{107}{1680} c_2^2 - \frac{c_1}{12} - \frac{c_2}{20}$$

$$\frac{\partial I}{\partial c_1} = 0 \implies \frac{19}{60}c_1 + \frac{11}{70}c_2 = \frac{1}{12}$$

$$\frac{\partial I}{\partial c_2} = 0$$
 $\implies \frac{11}{70}c_1 + \frac{109}{840}c_2 = \frac{1}{20}$

$$0.317 c_1 + 0.127 c_2 = 0.05$$
 \implies $c_1 = 0.177 , c_2 = 0.173$

$$\Rightarrow$$
 $y(2) = (0.177x - 0.173x^2)(1-x)$

Examples

I) The Brachistochrone (fastest descent) Curve Problem – Proposed by Johann Bernoulli (1696)

II) Structural Dynamics – Formulation of Governing Equation of Motion

Summary

- Many governing equations in physics and chemistry can be formulated by functionals. Finding the extremals of functionals can lead to the solutions in many problems in science and engineering.
- Euler's Equations provide the necessary condition for evaluating problems involving functionals. However, analyzing Euler's equations sometimes can be difficult.
- Calculus of Variations determines the extremals of functionals, even though the solution is only approximated.
- Calculus of Variations provides the theoretical basis for many methods in engineering, such as the Principle of Virtual Displacement (PVD) and the Finite Element Method (FEM).

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